



# In situ diffraction studies of phase-structural transformations in hydrogen and energy storage materials: An overview



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## ABSTRACT

The paper presents an overview of advanced *in situ* diffraction studies as a highly valuable tool to probe the structure and reacting mechanisms of hydrogen and energy storage materials. These studies offer benefits from the use of a high flux diffraction beam in combination with high resolution measurements, and allow, even when using very small samples, establishing the mechanism of the phase-structural transformations and their kinetics based on a rapid data collection for the various charge-discharge states at variable test conditions. The applied conditions include a broad range of hydrogen/deuterium pressures, from vacuum to high pressures reaching 1000 bar H<sub>2</sub>/D<sub>2</sub>, and temperatures, from cryocooling (2 K) to as high as 1273 K (1000 °C). Simultaneously, various state-of-charge and discharge are probed when studying metal-H/D systems for hydrogen/deuterium gas storage and as anode electrodes of metal hydride batteries. The range of the studied systems includes but is not limited to the AB<sub>5</sub>, AB<sub>2</sub> Laves type, AB, equiatomic ternary ABC intermetallics and Mg-containing layered AB<sub>3</sub> structures and composites. Prospects on Li-ion full cells are considered as well. Interrelations between the structure and hydrogen storage performance, including maximum and reversible H storage capacity, hysteresis of hydrogen absorption and desorption, H<sub>2</sub> absorption and desorption in dynamic equilibrium conditions, are considered and related to the ultimate goal of optimisation of the H storage behaviours of the advanced H storage materials. Numerous contributions of Dr. Michel Latroche to the field are highlighted, particularly in *in situ* studies during electrochemical transformations. The paper summarises a long-standing collaboration of the co-authors in the field, which also included 29 joint topical publications with Dr. Michel Latroche, and references 163 original publications.

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## 1. Introduction

Hydrogen is an energy vector that has a huge impact on a carbon-free energy economy, reducing greenhouse gas emissions as well as dependence on fossil fuel production. Hydrogen can be readily obtained from water using electricity from renewable sources. The energy is regained using a fuel cell to generate electricity (and water), or by the combustion of hydrogen. While further efficiencies may be reached in the production and use of hydrogen, the main stumbling block to large-scale use of hydrogen is storage and transport.

Although commercial hydrogen fuel cell electric vehicles (HFCEV) currently use high pressure (300–700 bar) pressurized gas tanks to store hydrogen on-board, the low density of gaseous

hydrogen at ambient conditions and the requirement to compress the gas resulting in a need for complex and expensive infrastructure remain a challenge. There is a need for alternative hydrogen storage systems that can reach high hydrogen densities at an affordable cost. An avenue that has been extensively studied is the use of various types of solid-state materials that are capable of reversibly storing hydrogen and providing it as a fuel to the fuel cell [1].

From a practical point of view, the time and ease of loading the material with hydrogen, the amount of hydrogen that can be usefully stored, the rate at which it can be retrieved, and the temperature and pressure conditions involved are critical requirements for determining the potential of a material. These practical use requirements are defined by the kinetics of hydrogen absorption and desorption, the thermodynamics of metal-hydrogen interactions (enthalpy of sorption/desorption), the diffusion of hydrogen through the material, the chemical bonding of hydrogen related to its location in the material's structure, chemical environment, and metal-

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hydrogen distances in the material as well as by material properties affecting its H storage performance.

In order to study these material properties and the features of the metal-hydrogen interactions, various techniques can be utilized [2], including gas absorption and desorption measurements to determine the H storage capacity, thermodynamics and kinetics, scanning and tunnelling electron microscopy to assess the particle sizes and elemental composition, Raman and infrared spectroscopy to investigate the chemical bonds, vibration modes and gas-solid interfaces, and neutron and X-ray scattering techniques that enable exploration of the material structure and the location and movement of hydrogen within that structure. For detailed examination of the physics and chemistry of the absorbed hydrogen interaction with the solid, neutron scattering is particularly well-suited as the neutrons penetrate the material and the hydrogen atoms when used as deuterium isotope have a large coherent neutron scattering length enabling their location in the crystal structures together with metal atoms [3].

Three main methods of neutron scattering are exploited to investigate the hydrogen-solid interactions and interface: neutron powder diffraction (NPD), inelastic neutron scattering (INS) and quasielastic neutron scattering (QENS) [4,5].

In NPD techniques, the neutrons are elastically scattered off the nuclei of the hydrogen and material atoms. For crystalline materials, the pattern of coherently scattered neutrons displays Bragg peaks which intensities depend on the types and positions of atoms in the crystal lattice, while incoherently scattered neutrons form a background in the powder diffraction pattern. Hydrogen has a large incoherent cross section, but in contrast its isotope deuterium has a large coherent cross section and a much smaller incoherent cross section. For this reason, most NPD studies are performed with deuterium wholly or partially substituted for hydrogen.

INS involves detecting the neutrons inelastically scattered while passing through the material, either by time-of-flight measurements or by filtering specific energies. INS measurements determine lattice vibrations and rotational transitions in amorphous and crystalline materials. This technique uses the incoherently scattered neutrons and can be employed for both hydrogen and deuterium.

QENS also uses the incoherent neutron scattering cross section but monitors energy changes in the inelastically scattered neutrons enabling the study of the hydrogen diffusion dynamics in the material.

While neutron scattering experiments are routinely performed successfully on *ex situ* samples, *in situ* studies have the advantages of examining the sample at variable testing conditions, when the temperature, gas pressure and hydrogen capacity are changed [6,7]. This enables the material to be studied while undergoing hydrogen ab- or desorption to investigate structural and chemical changes in the hydrogen-material composite to determine intermediate reaction products, phase transformations and reaction rates.

X-ray diffraction (XRD) can also be employed to study material structure [8,9], however, as the intensity of elastically scattered photons is proportional to the electron density, hydrogen is difficult to directly observe. XRD synchrotron facilities have the advantage of high photon flux compared to neutron sources, which enables the use of small sample sizes, for example milligrams, and rapid data collection times of a few seconds, depending on the sample and the instrument's beam energy and flux. In comparison, NPD requires sample of the order of grams, and collections times that are normally longer.

Previous reviews of the use of neutron diffraction for characterisation of potential hydrogen storage materials have covered the techniques as well as applications of neutron scattering. Klein et al. [4] discuss in detail the theory and practice of the different neutron scattering techniques introduced above. In addition, a number of types of solid-state hydrogen storage materials are described,

including metal hydrides, complex hydrides and porous materials, and the application of the scattering techniques to these categories.

The use of inelastic neutron scattering applied to high-pressure hydrides was also reviewed [10], focussing on INS studies of high-pressure monohydrides of *d*-metals of the VI–VIII groups and concluding that INS was effective at determining the H-H distances and therefore which materials do not obey the Switendick criterion, which postulates a minimum H-H distance of hydrogen atoms in the metal lattice of 2 Å.

Other neutron techniques have also been reviewed, including the use of SANS to investigate polymers, hydrogel, and bio-materials [11], neutron radiography and tomography to study the distribution of hydrogen in a storage bed [12], a comparison of neutron radiography to carrier gas hot extraction for the study of hydrogen mass transport in iron [13], neutron reflectivity to examine hydrogen concentration in multi-layer Fe-Nb thin films and a combination of *in situ* neutron radiography and IR-thermography to determine the driving forces for hydrogenation of sodium alanate [14].

An overview of the early works on the structural chemistry of intermetallic hydrides was presented in [15]. In a recent comprehensive review of AB<sub>2</sub> Laves type intermetallic compounds for use as storage materials, Yartys et al. [16], detailed many classes of Laves phase intermetallics and their structures as determined by neutron diffraction. Bououdina et al. [17] have also discussed Zr-based AB<sub>2</sub> Laves phase materials as part of a review of hydrogen storage materials. A review of the achievements of the European Cooperation in Science and Technology (COST) Action MP1103 [18] included use of *in situ* neutron diffraction for characterisation of nanostructured materials for solid-state hydrogen storage as well as neutron radiography to quantify hydrogen absorption by metal storage tanks. Reviews of more specific materials and the use of *in situ* neutron diffraction include those of lithium borohydride [19], magnesium borohydride [20], light metals [21] and porous materials [22].

In this article, the available neutron facilities are described and compared and the types of sample cells and sample environments necessary for *in situ* neutron studies are explored. In addition, a range of *in situ* neutron experiments yielding important results for hydrogen gas and electrochemical loading in a variety of potential materials for hydrogen storage applications is described. Among them, this manuscript contains 45 citations of the research articles by Dr. Michel Latroche, highlighting a great impact of his contribution to the field of hydrogen and energy storage materials, in particular when applying operando electrochemical studies related to the rechargeable batteries R&D.

## 2. Available neutron sources and their comparison

There is a large number of neutron facilities around the world. These can be divided into spallation sources and reactor sources, based on how the neutrons are generated. In a spallation source, a proton beam is accelerated into a heavy metal target, such as tantalum, tungsten, lead or mercury. Neutrons are ejected from the target metal atoms and slowed by moderators to provide a variety of energies to the instruments. In contrast to those produced in a nuclear fission reactor, the spallation source neutron beams can be pulsed by modulating the incident proton beam. While spallation sources have higher flux than reactor sources and use less energy, they are more expensive, and the targets require regular replacement. The main large neutron facilities currently in operation are listed below.

ISIS-Rutherford-Appleton Laboratories, United Kingdom has two spallation sources and generates a pulsed 800 MeV proton beam which passes through a carbon target (to produce muons) and then strikes a tantalum neutron target. Water, liquid methane and liquid hydrogen are used as moderators to slow the neutrons. Over 30 instruments use the neutrons generated, providing facilities for high

resolution powder diffraction, inelastic and quasielastic neutron spectroscopy as well as other techniques.

Institut Laue-Langevin, Grenoble, France uses a high-flux 58.3 MW nuclear reactor to provide a continuous total neutron flux of  $1.5 \times 10^{15}$  neutrons per second per  $\text{cm}^2$ . The neutrons are moderated by the heavy water used to cool the reactor as well as liquid deuterium. There are a total of 47 instruments, including 6 for powder diffraction, 6 for single crystal diffraction and another 6 for 3-axis spectrometry. The most suitable for the *in situ* NPD studies is a medium/high resolution two-axis diffractometer D20, which is capable of producing a neutron flux of  $10^8 \text{ s}^{-1} \text{ cm}^{-2}$ . D20 shows the fastest counting rate of any reactor-based neutron diffractometer and offers  $\Delta d/d \sim 2 \times 10^{-3}$ . Advanced electronics together with a variety of sample environments allows *in situ* time-resolved experiments at the timescale of a few tens of milliseconds [23]. Furthermore, at ILL, the D1B diffraction line has been widely used for real-time diffraction experiments as will be recognized several times in this review. Although the neutron flux at D1B ( $7 \times 10^6 \text{ n cm}^{-2} \text{ s}^{-1}$ ) is lower than that for the D20 instrument, high-quality diffraction patterns can be collected every 10 min thanks to a position-sensitive detector (PSD).

The NIST Center for Neutron Research (NCNR) reactor in Maryland, USA produces cold neutrons servicing 22 instruments for neutron scattering research. These instruments include small angle scattering (SANS) instruments, ultra-small (USANS) and very small angle scattering (VSANS), a high-resolution powder diffractometer and residual stress diffractometer as well as triple and multi-axis spectrometers.

Oak Ridge Neutron Facilities (SNS/HFIR) in Tennessee, USA, encompasses both reactor and spallation sources. The High Flux Isotope Reactor (HFIR) produces cold and thermal neutrons for research as well as isotope production and studies on gamma irradiation. The reactor supports 13 instruments for research including a neutron powder diffractometer, small angle, wide angle, magnetic and stress diffractometers as well as neutron imaging and spectrometers.

Housed in the same facility is the spallation neutron source (SNS). Negatively charged hydrogen ions are initially linearly accelerated before the electrons are stripped off and the resulting proton beam impacts a mercury spallation target. Pulsed at 60 Hz, the resulting neutrons are slowed by a moderator before use at 20 instruments, including powder and single crystal diffractometers, USANS and spectrometers.

Three powder diffractometers, General-Purpose Powder Diffractometer (POWGEN; SNS, ORNL), High-Resolution Powder Diffractometer (BT-1; NIST), and Neutron Powder Diffractometer (POWDER, HB-2A; HFIR, ORNL) are generally characterized by a narrow instrumental resolution over a large Q-range. These instruments are flux-limited, leading to data collection times of hours. However, when a fast data collection is targeted, instruments have been designed in a way that either the Q-range can be reduced, or the instrument resolution can be coarsened, increasing the flux on the sample by several orders of magnitude. When coupled with a large 2-dimensional detector, it is possible to record diffraction patterns in the sub-second regime, allowing for kinetic processes to be monitored. This provides alternatives to D20 in fast collection of the diffraction data.

The Los Alamos Neutron Science Center (LANSCE) is another spallation facility in New Mexico, USA. The Lujan center provides for scientific research producing cold and ultra-cold neutrons by cooling the neutrons produced in a tungsten target with water or solid deuterium moderators [24]. Instruments include diffractometers for stress and strain of polycrystalline materials and time-resolved studies, as well as neutron imaging/tomography and SANS.

The Swiss Spallation Neutron Source (SINQ) is a continuous spallation facility in the Paul Scherrer Institut (PSI), Switzerland.

Neutrons emitted from the interaction of a proton beam with a lead spallation target are cooled using a liquid deuterium moderator. 5 diffraction instruments use either thermal or cold neutrons for diffraction studies, with another two SANS instruments as well as reflectometers and spectrometers. The High Resolution Powder Diffractometer HRPT allows a data set collection which can be successfully processed by Rietveld refinements in just 3 min, and this has been utilised during the studies of several metal hydride materials described in this review.

FRM-II is a research reactor-based neutron facility in Munich, Germany. 25 scientific instruments as well as irradiation facilities use neutrons from the reactor and cooled through a liquid deuterium moderator. There are 11 instruments for structure and stress studies as well as SANS instruments and imaging and spectroscopy instruments.

The Institute for Solid State Physics, Tokyo University, Japan (ISSP) Neutron Scattering Laboratory in Tokai, is a reactor driven neutron facility with high resolution powder diffraction, single crystal diffraction, SANS, VSANS and USANS instruments. Imaging and spectrometry capabilities are also available.

J-PARC (Japan Proton Accelerator Research Complex) is a proton accelerator facility in Tokai, Japan. Three accelerators support facilities dedicated to materials and life sciences as well as particle physics and nuclear transmutation. A 3 GeV Rapid Cycling Synchrotron provides a pulsed proton beam to a mercury spallation target with liquid hydrogen moderators to provide neutrons and muons for the Materials and Life Science Experimental Facility (MLF). Instruments using elastic neutron scattering in MLF include single crystal diffraction, powder diffraction, small-angle and total scattering, and reflectivity instruments, while for inelastic scattering there are instruments for time-of-flight and spin echo spectrometry, as well as imaging and elemental analysis. The Super High Resolution Powder Diffractometer accommodated at the Materials and Life Science Experimental Facility has achieved a high resolution of  $\Delta d/d = 0.0365\%$  and allows access to an expanded Q-range during the data collection [25,26].

The Australian Centre for Neutron Scattering in Lucas Heights, Australia uses the OPAL (Open Pool Australian Lightwater) reactor to generate neutrons for use in 15 instruments, including two powder diffraction instruments (high resolution and high intensity) as well as strain, SANS, USANS, reflectometry and spectrometry instruments.

### 3. Experimental cells

*In situ* neutron scattering experiments investigating the interaction of hydrogen with a material demand a sample cell capable of maintaining the experiment conditions, such as pressure and temperature as well as interfering as little as possible with the passage of the neutron beam. In addition, the cell must be constructed of suitable material as hydrogen diffuses through many metals at higher temperatures and embrittles some metals, making the cell susceptible to catastrophic failure [27,28]. Hydrogen embrittlement is apparent in iron and most iron-based alloys such as martensitic steels [29], due to hydrogen diffusion along grain boundaries and concentration in crack tips [30]. Hydrogen also reacts to form hydrides with metals such as zirconium, tantalum, vanadium and palladium, which have much lower yield strengths and therefore are unable to withstand higher pressures of hydrogen gas. Thus, the walls of the sample cells aiming for a decreased contribution of the cell to the diffraction pattern should eliminate a direct contact with hydrogen gas if the mentioned metals are utilized. As many stainless steels such as 316 L allow only a negligible hydrogen diffusion at moderate temperatures ( $< 350 \text{ }^\circ\text{C}$ ) and pressures ( $< 100 \text{ bar}$ ), such type of steel is most frequently used as a material of choice to build the samples cells.

Because the diffusion of hydrogen through a SS 316 L metal cell wall becomes significant over 470 °C, it is possible to hydrogenate/dehydrogenate sample material enclosed in a crimped iron envelope without exposing the sample to oxygen [31]. However, generally diffusion invalidates the experiment because the quantity of hydrogen in the cell becomes no longer known, particularly with increasing H<sub>2</sub> pressures, and thus a special H-impermeable material selection should be made to eliminate this problem. Furthermore, hydrogen embrittlement should be properly addressed when making the material selection to safely build the sample cells and to eliminate any possibility of a catastrophic failure of the equipment during the H<sub>2</sub> loading tests.

The material of the cell wall is also exposed to the neutron radiation, which can add an unwanted signal to the scattered neutron data. While this can sometimes be filtered out of the data, it often complicates the data analysis or obscures sample data. Metals with very low coherent scattering cross sections, such as vanadium, contribute very little to the signal from the sample container (this metal however reacts with hydrogen and therefore should not be exposed to it). Null-matrix alloys made from two metals with coherent scattering lengths of opposite sign, such as titanium (-3.438 fm) and zirconium (7.16 fm), also contribute negligibly to the signal. In addition, many materials become radioactive after exposure to the neutron beam, for example, in stainless steel after several days of thermal neutron bombardment, isotopes of cobalt require long on-shelf storage for the radioactivity to reduce to a safe level.

The pressure and temperature conditions of the experimental work, as well as possible reactions with the sample, place limitations on the materials available for cell construction. For example, while aluminium does not embrittle or allows hydrogen to diffuse through, it is only suitable for low pressure experiments due to the moderate yield strength. Stainless steels are only slightly permeable to hydrogen at 200 °C, so they be used at up to 350 °C depending on the wall thickness, but H<sub>2</sub> diffusion is quite strong at the temperatures exceeding 500 °C. Vanadium cans, which are extensively used as sample holders at neutron facilities due to transparency to neutrons across a wide temperature range, cannot be used together with hydrogen due to their embrittlement. Ti<sub>68</sub>Zr<sub>32</sub> is a strong null-matrix alloy but also embrittles after exposure to hydrogen. Techniques such as plating the inside of the sample cell with a material impervious to hydrogen, or using a sleeve of appropriate material, may be used to overcome these problems. The following cells have been developed for specific sets of experimental conditions for *in situ* experiments with hydrogen.

### 3.1. Stainless steel cell to 100 bar

A stainless steel cell has been used successfully to 100 bar of hydrogen at moderate temperatures for *in situ* studies on potential hydrogen storage materials at ISIS and SINQ [32,33]. Manufactured from 316 L stainless steel, with an internal diameter of 10 mm, and a wall thickness of 0.75 mm, this cell was used at ISIS for studies on palladium as well as LaNi<sub>5</sub> [6], with and without a collimator. Fig. 1 shows the geometries required to use a collimator of a neutron absorbing material such as BN in order to restrict both the input beam and the detected neutrons to the centre of the sample, excluding as much as possible of the sample cell.

A similar cell (outer diameter 9.5 mm, inner diameter 7.8 mm) made of stainless steel was designed J.-M. Joubert [34] and used at ILL for deuterium pressures up to 100 bar and 300 °C [35–37].

IFE has been using SS 316 L cylindrical autoclaves (limiting H<sub>2</sub>/D<sub>2</sub> pressure to 50 bar) with external diameter 6 mm and ¼" with decreased thickness of the wall at the sample location, allowing them to fit into the standard fitting to the cryostat and the heating resistance furnace at HRTP, SINQ, PSI.

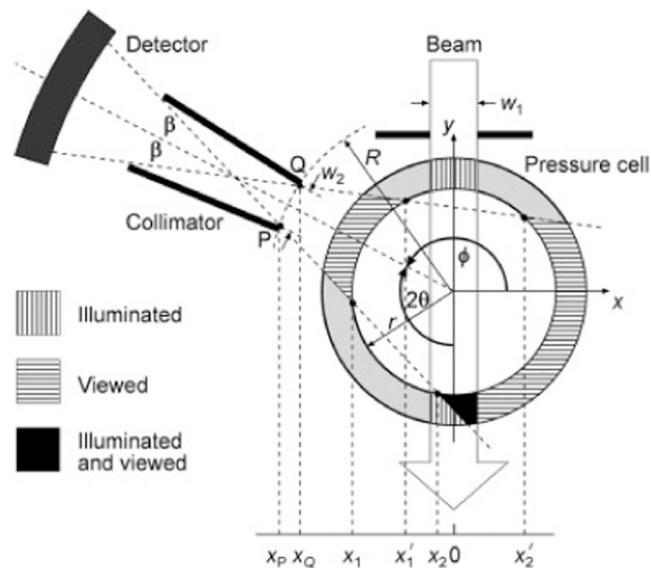


Fig. 1. Schematic of a collimated steel sample cell. (Reprinted with permission from [6]. Copyright (2012) Elsevier).

### 3.2. Zero matrix HP cell

A high pressure sample cell manufactured from Ti<sub>68</sub>Zr<sub>32</sub> has been used for *in situ* experiments at SINQ [38] and ILL [39,40]. The zero matrix alloy does not scatter neutrons coherently and does not contribute significantly to the diffraction data pattern. Because TiZr embrittles, a thin (0.1 mm) stainless steel sleeve was fitted inside the TiZr to prevent hydrogen from interacting with the cell body. Although the steel contributes to a signal, this has a small effect because of the low thickness, with the sufficient tensile strength of the TiZr allowing to contain the hydrogen pressure. Up to 1000 bar D<sub>2</sub> has been applied at SINQ [38] and ILL [39,40] (Fig. 2).

It is worth remarking that a zero matrix TiZr cell has been successfully implemented for the study of electrochemical reactions in Li-ion batteries [41].

### 3.3. Aluminium cell used on Engine-X

An aluminium cell with an ID of 18 mm and a wall thickness of 1 mm (Fig. 3) was constructed to perform a spatial analysis of the hydrogen adsorption by LaNi<sub>5</sub>. For this experiment, use of Al cell allowed a good thermal management as the heat of reaction was rapidly distributed through the water-cooled copper plates at each end of the cell. Engine-X, a stress/strain instrument at ISIS, was used to investigate spatial inhomogeneities in the uptake of deuterium following a rapid absorption step [42].

### 3.4. Sapphire cell

As a single crystal, sapphire has a high yield strength making it suitable for high pressure work. Sapphire is frequently used in X-ray diffraction *in situ* studies [8,43] as it is relatively easy to mask the peaks in the scattering patterns caused by the sapphire, when its single crystal is used. Ikeda, Ohshita, Otomo, Sakaki, Kim, Nakamura, Machida and Von Dreele [44] have used a commercial single crystal sapphire container (Fig. 4a) with a 5.8 mm ID, 2.5 mm wall thickness and a 5.0 mm base thickness for *in situ* neutron diffraction measurements on LaNi<sub>4.5</sub>Al<sub>0.5</sub> with the total scattering instrument, NOVA at JPARC in Japan.

Results were compared with an *ex situ* sample in a vanadium can, shown in Fig. 4b.



Fig. 2. Zero matrix high pressure TiZr cell with thin stainless steel liner.

### 3.5. Electrochemical cell

The electrochemical three-electrode cell developed at CNRS, France, for the study of electrode materials in Ni-MH batteries is shown in Fig. 5 [45]. It is based on an the developed earlier by Chabre et al. cell for the study at ILL of proton insertion in  $\gamma$ -MnO<sub>2</sub> [46,47]. The body of the cell is made of amorphous silica leading to a

wavy diffraction background. For the study of metal hydride materials, the working electrode is a composite electrode made from intermetallic powder ( $< 100 \mu\text{m}$ ) as active material mixed with carbon black and PTFE. The mixture is mechanically compressed over a nickel grid and rolled up on itself to form a cylinder of about 5 cm height and 10 mm diameter ( $\emptyset$ ). This working electrode is then sandwiched between inner ( $\emptyset$  8 mm) and outer ( $\emptyset$  12 mm) counter-

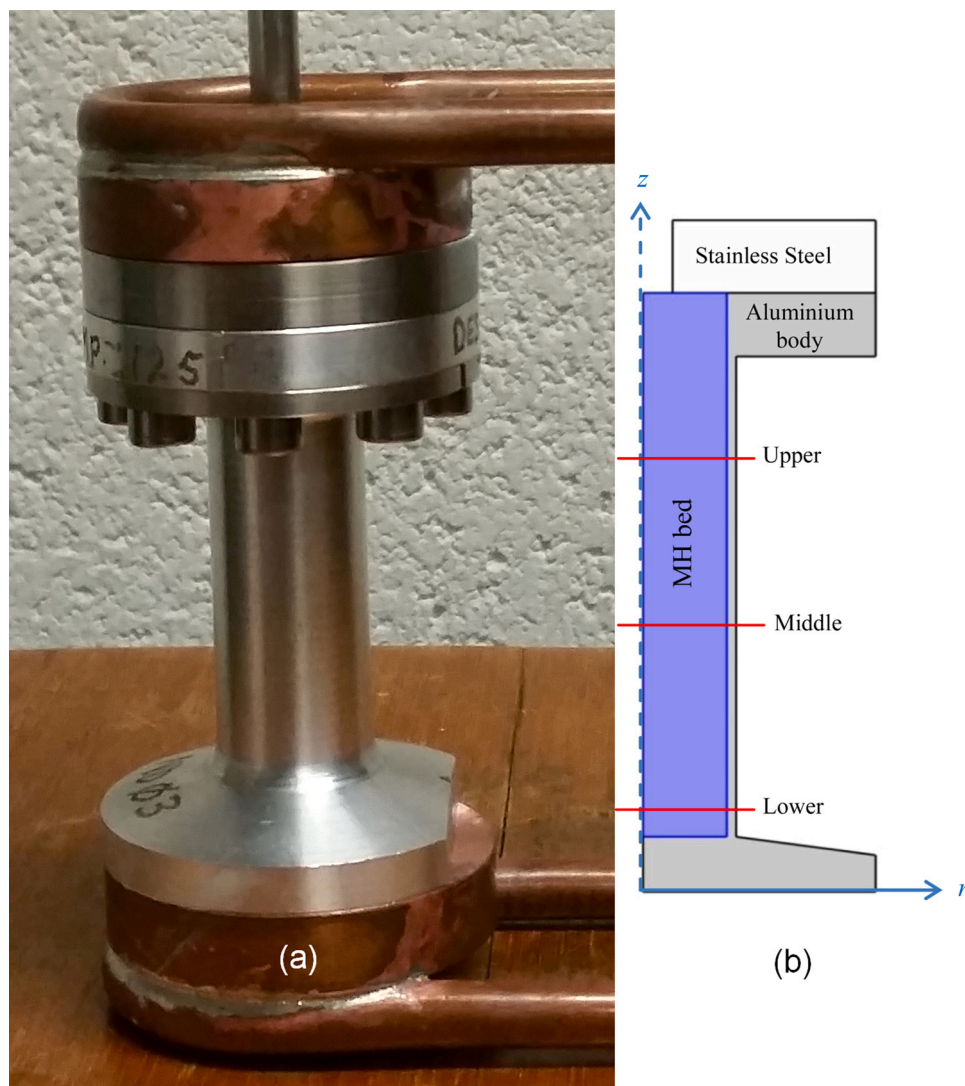
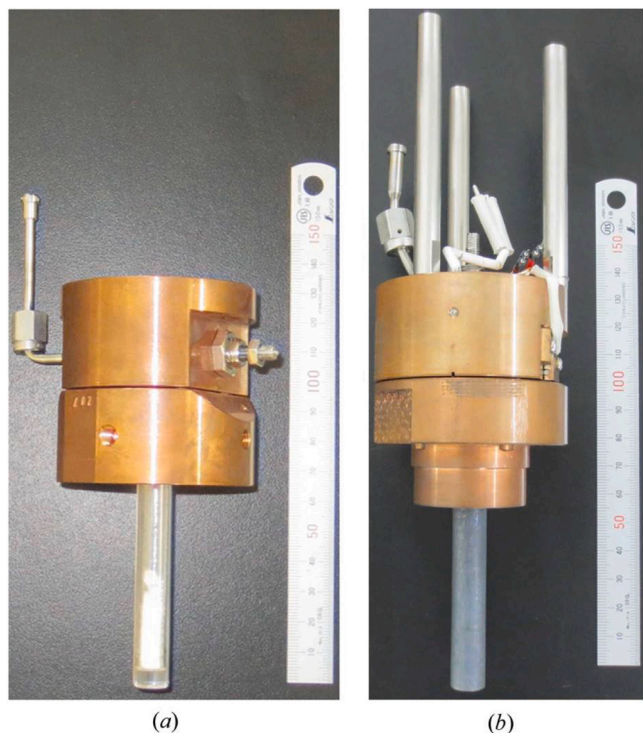


Fig. 3. Aluminium *in situ* cell employed at HRPD, ISIS for analysis of spatial inhomogeneities in deuterium absorption. (a) Layout; (b) Scheme. (b) reprinted with permission from [42]. (Copyright (2016) Elsevier).



**Fig. 4.** Sapphire (a) and vanadium (b) *in situ* cells [44]. Reprinted from an open access article distributed under the Creative Commons Attribution license (<http://creativecommons.org/licenses/by/4.0/>).

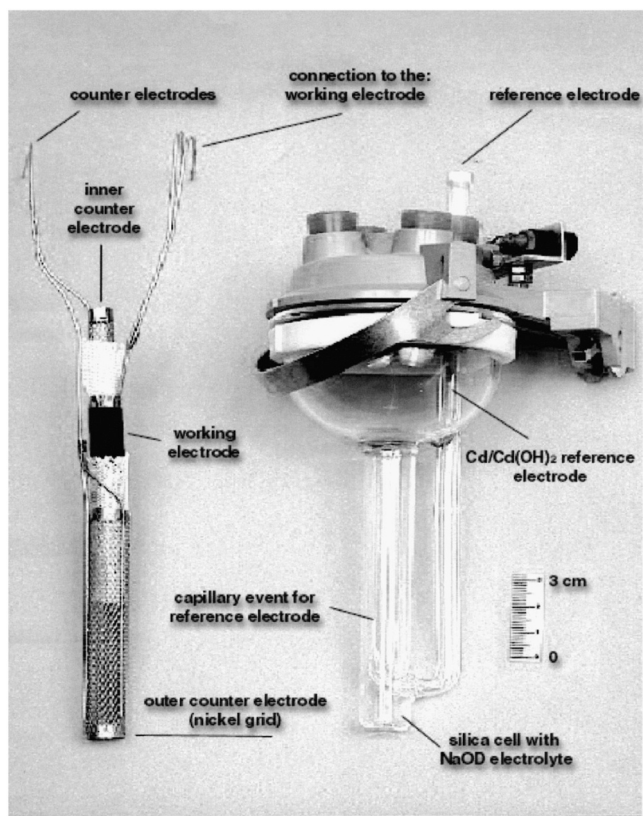
electrode cylinders made of nickel grid, with silica sheaths as separators on each side of the working electrode. The electrochemical couple  $\text{Cd}/\text{Cd}(\text{OH})_2$  is used as reference electrode. It is located in a side tube entering the main tube with a capillary going against the bottom side of the working electrode. The three electrodes are immersed in deuterated alkaline electrolyte (KOD [45,48] or NaOD [49] in  $\text{D}_2\text{O}$ ) to take advantage of the high coherent scattering section of deuterium. This cell has been widely used for *in situ* structural studies of both negative [45, 50–52] and positive [53] electrode materials for Ni-MH batteries. Its geometry is very close to that of commercial-type batteries and well-adapted to the Debye-Scherrer powder neutron diffraction beamlines. Recently, a coin-type cell has also been developed at the ISIS facility [54].

### 3.6. Coupling to ancillary equipment

During *in situ* neutron diffraction studies, the diffraction data is monitored as a function of different variables such as deuterium pressure, temperature or electrochemical potential. Changes in these variables modify the deuterium content and phase composition in the alloy. The refinement by the Rietveld method of deuterium site occupancies and phase amounts allows to determine the overall deuterium content in the sample during the experiment. However, these results require confirmation by complementary techniques to check and reinforce the validity of structural models. Thus, to control and measure the overall deuterium content in the alloys, the neutron diffraction cells are connected to different ancillary equipment such as Sieverts' apparatuses, pressure gauges or galvanostats. Moreover, most of these devices not only monitor changes on deuterium content but also control the time-evolution of variables during the experiments. Most popular devices are now described.

#### a) Manometric Sieverts' apparatus for Pressure-Composition-Temperature (PCT) measurements

Manometric Sieverts' apparatuses consist of calibrated cell volumes equipped with pressure transducers and valves for connection to different lines: deuterium inlet, deuterium venting and vacuum line. Fig. 6 shows a manual custom Sieverts' apparatus developed at CNRS and used in many experimental *in situ* diffraction studies [35–37,55,56]. It works at room temperature and deuterium



**Fig. 5.** View of the various parts of the *in situ* neutron diffraction electrochemical cell assembly. Reprinted with permission from [45]. (Copyright (2002) Elsevier).



**Fig. 6.** Manual homemade Sieverts' apparatus developed at CNRS for *in situ* neutron diffraction studies. At the top: inlets for  $\text{H}_2(\text{D}_2)$ , vacuum ("Vide") and venting ("atm"). They are connected through copper capillaries and yellow, green and green valves, respectively, to the main body of the apparatus. Two pressure gauges (1 and 100 bar) are located at the top-right. The three black valves connect the sample colder (PE outlet) to two large reservoirs (R1 and R2).

pressure between 100 Pa and 10 MPa. A long thin copper capillary tube connects the sample holder, placed in the beam, to the Sieverts' apparatus, where the user is located several meters away from the beam to avoid radioactive exposure. Using this system, deuterium uptake by the sample is measured by first determining by manometry the number of moles of deuterium in a closed calibrating volume, allowing the gas to expand towards the sample and, after completion of the solid-gas reaction, measuring again by manometry the amount of residual gas in the whole system. The difference between initial and final amount of deuterium corresponds to the sample deuterium uptake. The measurement is done by stepwise increasing (absorption) and decreasing (desorption) the pressure to register Pressure-Composition-Isotherms at constant temperature. Accurate measurements can be made when the appropriate care and methods are employed [57] and the real gas equations are used at higher pressures (> 1 MPa) [58].

#### a) Vacuum system and pressure gauges for Thermal Desorption (TD) measurements

Simultaneous measurement of neutron diffraction and thermal desorption data has been widely reported [36,56,59,60]. Both signals are synchronized in time and temperature which are correlated through a linear heating ramp. With a typical heating rate of 0.5 K/min and acquisition time of 5 min per diffractogram, the temperature resolution of diffraction patterns is 2.5 K [60]. Stainless steel or silica sample holders have been used. The choice is mainly based on the preference for diffraction patterns to present either additional Fe-lines [36] (Fig. 7) or a wavy background [61], respectively.

To monitor deuterium desorption, the sample holder is connected to a vacuum line equipped with a pump and a pressure sensor. A rotary pump is often used (residual pressure,  $P_{res} \sim 1$  Pa) since it easily evacuates the high amount of deuterium stored in large neutron samples (mass from 5 to 10 g) and ensures a constant pumping speed in the operational pressure range (typically between 1 and 1000 Pa). Assuming fast pumping rates and small sample-holder volume, the deuterium desorption rate ( $dx/dt$ ) from the sample is proportional to the released deuterium pressure [62]:

$$-\frac{dx(t)}{dt} = \frac{S}{RT} P(t)_{H_2} \sim AP(t)_{H_2} \quad (1)$$

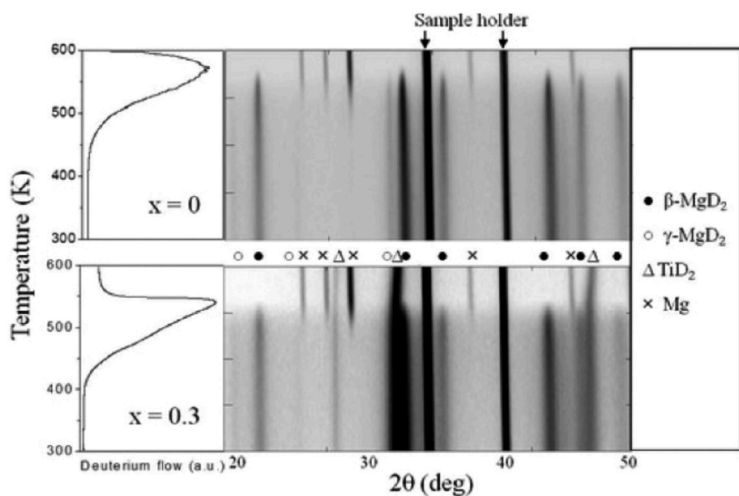
where  $x$  is the number of deuterium atoms per alloy formula unit,  $S$  is the pumping speed,  $R$  the gas constant,  $T$  the temperature of the vacuum line and  $P_{H_2}$  the gas pressure. In absence of a quadrupole mass spectrometer, the released gas is assumed to only consist of deuterium. The deuterium pressure is measured by a Pirani [60] or a capacitance manometer [36]. If the pumping speed  $S$  is constant during the whole desorption process, the monitored pressure scales to the deuterium desorption rate by the proportionality factor  $A$ :

$$A = \frac{x_{t=0}}{\int P dt} \quad (2)$$

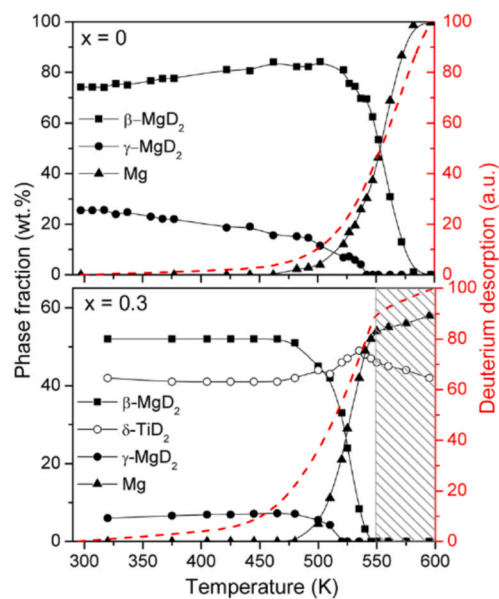
where  $x_{t=0}$  is the initial deuterium content in the sample, *i.e.* before starting the thermal desorption. In the end, as shown in Fig. 7, the deuterium desorption is synchronized with the neutron diffraction patterns (a) and the phase amounts, determined from Rietveld refinements, are correlated to the deuterium desorption signal (b).

#### a) Galvanostat device for electrochemical cycling

Coupling between electrochemical data and neutron acquisition is facilitated by the fact that both methods can be computer-monitored and their base time can be easily synchronized. Often, the electrochemical process is in galvanostatic mode and driven with a microprocessor-controlled galvanostat [45,49,51,63]. Typically, an applied charge current is imposed for a time exceeding the expected electrode capacity to take into account the competing Hydrogen Evolution Reaction (HER) through Heyrovsky or Tafel mechanisms [64]. On discharge, a constant current is applied until a given threshold potential is reached, typically 0.5 V against Cd/Cd(OH)<sub>2</sub> reference electrode. Then, to avoid electrode oxidation, the current is automatically switched-off and the electrode enters relaxation. Since one electron is provided by each hydrogen atom released from the electrode to the solution, the coulombic discharge,  $Q_{dis}$ , relates to the released hydrogen amount as:

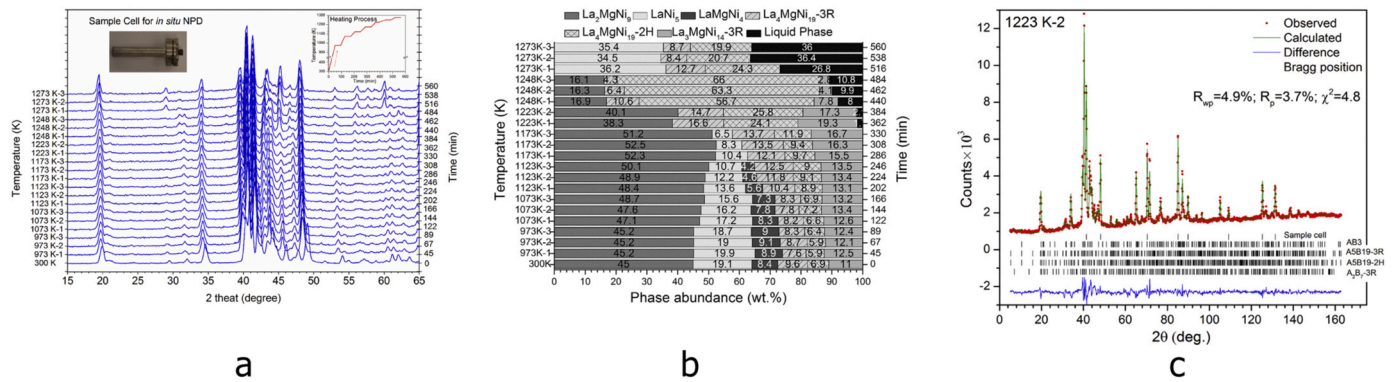


a



b

**Fig. 7.** (a) Simultaneous acquisition of neutron-diffraction (ND) patterns and thermal desorption (TD) spectra of  $(1-x)\text{MgD}_2 - x\text{TiD}_2$  nanocomposites for  $x = 0$  and 0.3 samples. Heating rate 0.5 K/min. (b) TD spectra and 2D evolution of ND patterns as a function of temperature. Evolution of the refined phase amounts with temperature. Reprinted with permission from [36]. (Copyright (2013) American Chemical Society).



**Fig. 8.** (a) Evolution of the *in situ* NPD pattern of  $\text{La}_2\text{MgNi}_9$  compounds during heating (K-1, K-2 and K-3 represent sequence of hold times; 0–22 min; 22–44 min; and 44–66 min, respectively). Inset: SS sample cell used to accommodate the alloy. Inset: temperature profile during the annealing. (b) Phase abundances as a function of temperature and heating time during the annealing of the as-cast  $\text{La}_2\text{MgNi}_9$  alloy from Rietveld refinements of the NPD data. (c) *In situ* powder neutron diffraction pattern of as-cast sample collected at 1223 K. Vertical bars show positions of the Bragg peaks for the phase constituents (from top to bottom): Fe (sample cell),  $\text{AB}_3$ ,  $\text{A}_5\text{B}_{19-3\text{R}}$ ,  $\text{A}_5\text{B}_{19-2\text{H}}$  and  $\text{A}_2\text{B}_7-3\text{R}$  phases. (Reprinted with permission from [65]. Copyright (2016) Elsevier).

$$\Delta x = Q_{\text{dis}}/nF \quad (3)$$

where  $\Delta x$  is the number of hydrogen atoms released by alloy formula unit,  $F$  is the faraday constant ( $26.8 \text{ Ah mol}^{-1}$ ) and  $n$  the number of alloy mols. Thus, the amount of electrochemical released deuterium can be compared with the evolution of deuterium content in the electrode as determined by refinement of neutron diffraction data.

#### 4. Temperature dependent transformations

##### 4.1. Phase-structural transformations in the metal hydride anode

###### 4.1.1. La-Mg-Ni alloys on their annealing between 300 K and 1273 K [65]

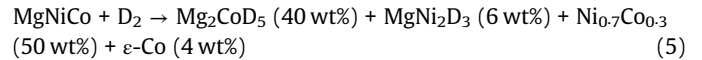
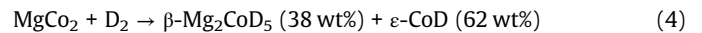
Studies of the temperature-induced phase-structural transformations in the as-cast  $\text{La}_2\text{MgNi}_9$  metal hydride battery electrode alloy were performed by *in situ* neutron powder diffraction at HRPT, SINQ/PSI at temperatures ranging from 300 K to 1273 K. Initial alloy is multi-phase structured, containing six different intermetallics with five stoichiometric compositions. These include the targeted rhombohedral  $\text{La}_2\text{MgNi}_9$  as the main phase constituent, together with three electrochemically active intermetallics,  $\text{La}_3\text{MgNi}_{14}$  (3 R),  $\text{La}_4\text{MgNi}_{19}$  (2 H) and  $\text{La}_4\text{MgNi}_{19}$  (3 R). Furthermore, the alloy contained two “ballast” intermetallics,  $\text{LaNi}_5$  and  $\text{LaMgNi}_4$ . Various transformations take place on heating, leading to the significant changes in the contents of the constituent intermetallics, first to the disappearance of  $\text{LaNi}_5$  and  $\text{LaMgNi}_4$ . Thermal volume expansions of the studied intermetallics were well-fitted by the linear dependences and resulted in similar values of 4.3–5.5 vol% at 1223 K as compared to 300 K. The data presenting results of the studies are given in Fig. 8, a, b and c. The temperature-dependent mechanism of transformations is closely related to the content of magnesium in  $\text{La}_{3-x}\text{Mg}_x\text{Ni}_9$  ( $x = 1.0, 1.1$  and  $1.2$ ) alloys [66].

###### 4.2. HDDR process in $\text{MgCo}_2\text{-D}_2$ system resulting in recombination of $\text{MgCo}_2$ and a different mechanism of interaction for $\text{MgCoNi}$ alloy [67]

$\text{MgCo}_2$  and  $\text{MgNiCo}$  crystallize with hexagonal Laves type intermetallic structures of the C14 type and do not form hydrides at ambient hydrogen pressures. However, applying high hydrogen pressures in the GPa range forces the hydrogen absorption and leads to the formation of multi-phase compositions, which contain 2.5 at. H/f.u. of  $\text{MgCo}_2$  or  $\text{MgNiCo}$  and remain thermally stable under

normal conditions. The hydrogenation of  $\text{MgCo}_2$  resulted in its decomposition to a ternary  $\text{Mg}_2\text{CoD}_5$  deuteride and metallic cobalt. The interaction of  $\text{MgNiCo}$  with deuterium under the synthesis conditions of 2.8 GPa and 200 °C proceeded in a more complex way. A very stable ternary deuteride  $\text{MgNi}_2\text{D}_3$  was leached away while Co was separated in the form of  $\text{Mg}_2\text{CoD}_5$  and the remaining nickel formed a solid solution with Co with the approximate composition  $\text{Ni}_{0.7}\text{Co}_{0.3}$ .

The studied transformations can be described by the following equations:



Time-resolved *in situ* neutron diffraction and Thermal Desorption Spectroscopy were employed to probe temperature-dependent phase-structural transformations in the  $\text{MgCo}_2\text{-D}$  system. Some experimental data of the studied of transformations in the  $\text{MgCo}_2\text{-D}_2$  system are shown in Fig. 9, a, b and c. The following mechanisms of decomposition of the  $\text{Mg}_2\text{CoD}_5$  deuteride and the formation of Mg-Co compounds in the Co-rich part of the Mg-Co system were proposed and it was concluded that:

(a) The  $\text{Mg}_2\text{CoD}_5$  deuteride decomposes into elementary Mg and Co at approximately 300 °C.

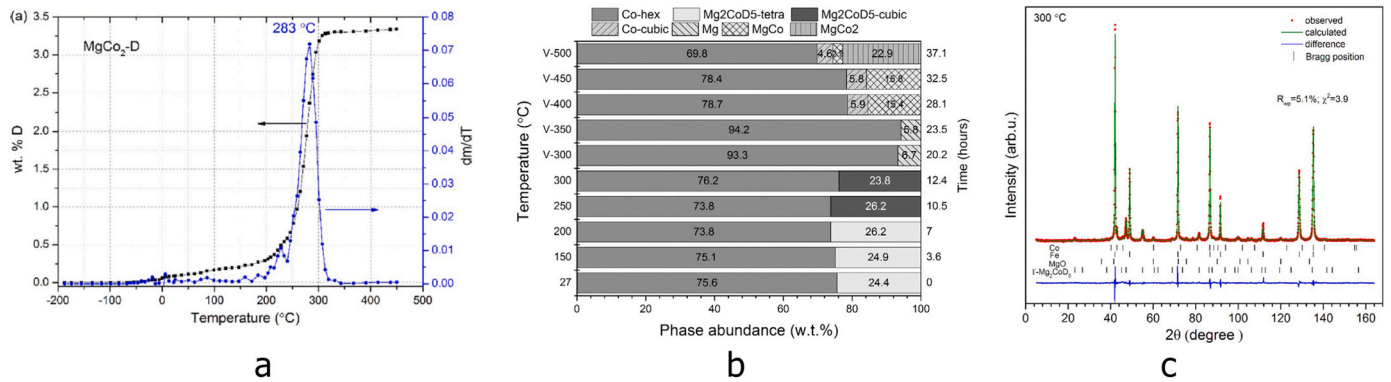
(b) A polymorphic transformation from the LT tetragonal to the HT cubic modification of  $\text{Mg}_2\text{CoD}_5$  takes place at around 250 °C.

(c)  $\text{MgCo}$  and  $\text{MgCo}_2$  compounds are formed in the Mg-Co matrix by solid-state reactions at 400 and 500 °C, respectively.

(d) A formation of  $\text{MgCo}_2$  after completion of the deuterium desorption shows that this Mg-Co intermetallic is a new example of the hydrogen storage alloy, for which a complete HDDR cycle was successfully accomplished.

The mechanism of interaction of the alloys with hydrogen becomes significantly affected by the transition metals, Co and Ni. *In situ* neutron powder diffraction of phase-structural transformations accompanying the vacuum desorption of deuterium in the temperature range of 27–500 °C showed a complete recovery of the initial  $\text{MgCo}_2$  intermetallic via a Hydrogenation-Disproportionation-Desorption-Recombination process.  $\text{MgCo}_2$  is an example of the hydrogen storage alloy, in which a successful HDDR processing results in the reversible formation of the initial intermetallic at much lower temperatures than in the equilibrium phase diagram of the Mg-Co system.





**Fig. 9.** (a) Thermal Desorption Spectroscopy data showing one event of deuterium vacuum desorption peaking at 283 °C; (b) Temperature-driven evolution of the phase-structural composition of the disproportionated  $\text{MgCo}_2\text{D}_{2.5}$  resulting in a recombination of  $\text{MgCo}_2$  intermetallic compound; (c) Neutron powder diffraction pattern collected at 300 C and showing a formation of the HT cubic modification of  $\text{Mg}_2\text{CoD}_5$  deuteride and cobalt.  $\text{Mg}_2\text{CoD}_5$  decomposes to the metallic Mg and Co in the cause of the desorption process before forming MgCo and, eventually,  $\text{MgCo}_2$ .

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### 4.3. Low temperature transformations at subzero temperatures

#### 4.3.1. Order-disorder transitions in Laves type $\text{ErMn}_2\text{D}_2$ hydride

$\text{ErMn}_2$  intermetallic compound crystallizes with hexagonal C14  $\text{MgZn}_2$  type of crystal structure and orders magnetically with magnetic moments beared by both Er and Mn.  $\text{ErMn}_2$  forms several individual hydrides with a maximum H storage capacity reaching 4.3 at.H/f.u. and temperature-dependent magnetic properties being sensitive to the hydrogen content, lattice expansion and symmetry and Mn-Mn distances which when exceeding a critical value of 2.7 Å (See Fig. 10 (a)) cause a change of the sign of magnetic interaction between the Er and Mn magnetic moments and a ferromagnetic or antiferromagnetic coupling [68].

$\text{ErMn}_2\text{D}_2$  deuteride has been studied by high-resolution synchrotron X-ray powder diffraction at temperatures between 150 and 298 K. Below 210 K a transformation from the hexagonal C14 Laves phase type structure (space group  $P6_3/mmc$ ;  $a = 5.55357(3)$ ;  $c = 9.05368(7)$  Å;  $T = 298$  K) into the monoclinic one crystallising in the space group  $C2/m$  ( $a = 9.61247(9)$ ;  $b = 5.57558(4)$ ;  $c = 9.07102(8)$  Å;  $\beta = 90.5451(5)^\circ$ ;  $T = 150$  K) takes place (see Fig. 10(b)) [69].

A rather large two-phase region of coexistence of both structural modifications was observed between 185 and 205 K (Fig. 11, (a) and

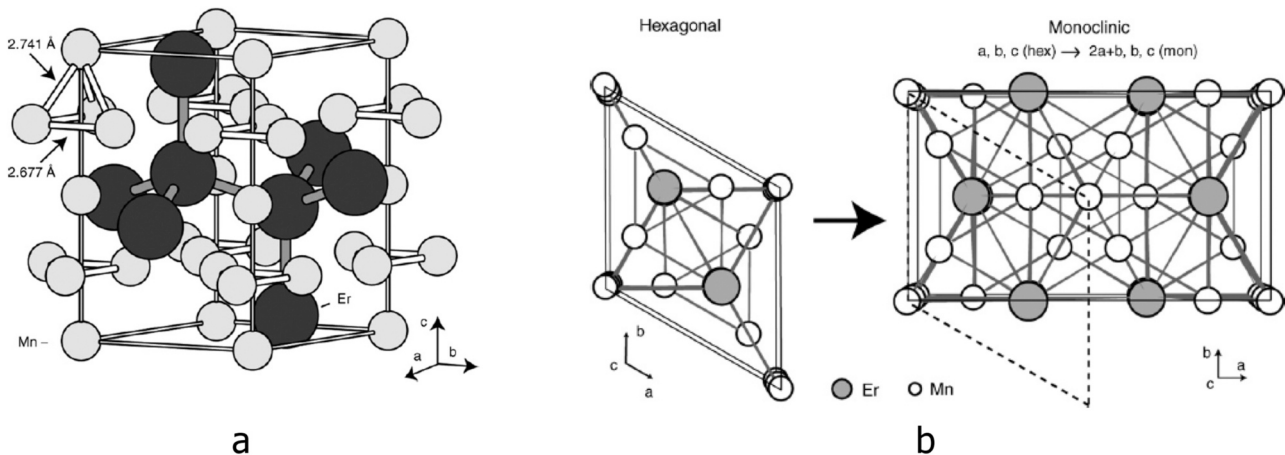
(b)). Interestingly, formation of the LT monoclinic structure on cooling down is accompanied by a volume expansion.

#### 4.3.2. Equiatomic ternary $\text{ABCH}_{1-2}$ hydrides

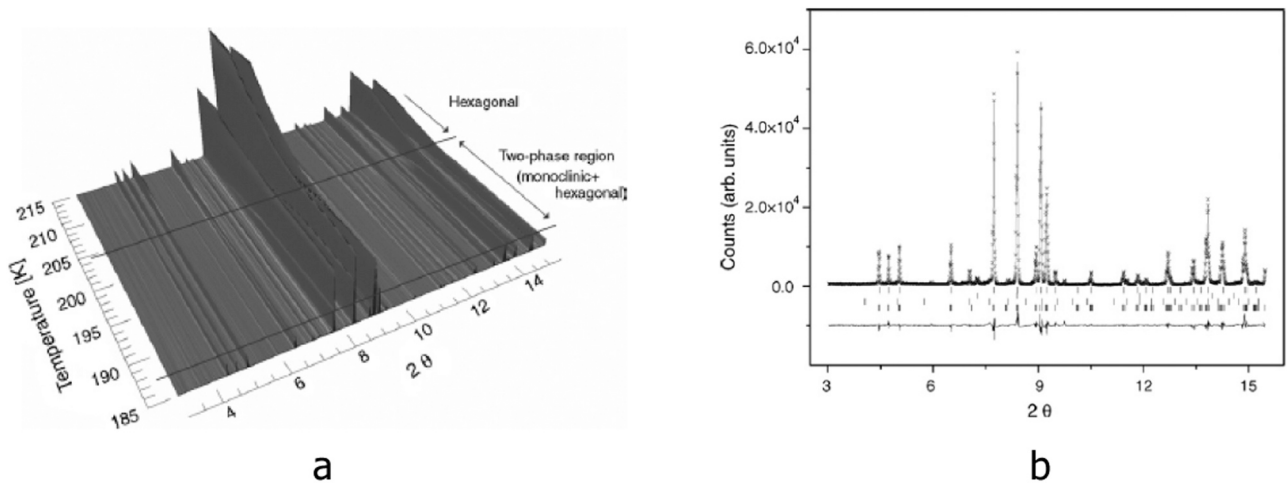
$\text{RTX}$  intermetallic compounds with  $R$ =Rare Earth Metal;  $T$ =Transition Metal, Fe, Co, Ni;  $X$ =Non transition element, Al, Si, Ge, Ga, Sn, In, crystallize with a variety of closely related structure types which can be obtained from the hexagonal  $\text{AlB}_2$  type through the ordering of B and C components and *via* internal and external deformation of the structures (see Fig. 12A). When the atoms of transition metal T are substituted by atoms X in an ordered way, this results in the structures of ZrBeSi, LiGaGe and TiNiSi types (see Fig. 12B).

Analysis of the crystal structure of the  $\text{RTX}$  hydride with TiNiSi type structure shows that it contains 4 types of interstitial sites shown in the Fig. 13a, which are suitable for the accommodation of hydrogen atoms. Neutron diffraction study of the dihydride  $\text{TbNiSiD}_{1.78}$  showed that D atoms completely fill the available  $\text{Tb}_3\text{Ni}/R_3\text{T}$  tetrahedra which form a spatial network *via* connecting to each other by vertexes and edges of the tetrahedra (see. Fig. 13b).

Table 1 gives an overview of the crystallographic and magnetic properties of several ternary metal hydrides.



**Fig. 10.** (a) The metal sublattice in the hexagonal crystal structure of  $\text{ErMn}_2\text{D}_2$  at room temperature. The  $\text{Mn}_4$  tetrahedra are slightly elongated along the  $c$ -axis (Mn-Mn interatomic distances in the  $(a, b)$  plane are 2.677 Å; distances out of the plane, along  $[001]$  are 2.741 Å). (b) Scheme of the transformation of the unit cell of the hexagonal phase of  $\text{ErMn}_2\text{D}_2$  existing at room temperature (left) into the monoclinic low temperature structure (right); the original hexagonal cell is shown with the dashed lines). During this transformation, the Er position splits into two positions (Er $\rightarrow$ Er1 and Er2), and the two Mn positions both split into two (Mn1 $\rightarrow$ Mn1 and Mn2; Mn2 $\rightarrow$ Mn3 and Mn4). (Reprinted with permission from [69]. Copyright (2006) Elsevier).



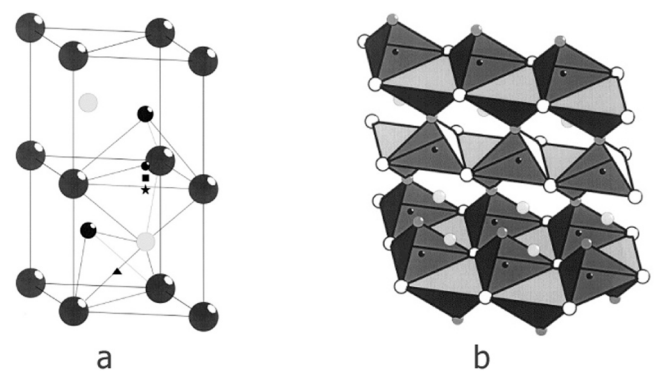
**Fig. 11.** (a) Evolution of the *in situ* SR-XRD patterns of the  $\text{ErMn}_2\text{D}_2$  as a function of temperature ( $T = 185\text{--}215$  K in steps of  $\Delta T = 5$  K;  $\lambda = 0.3750$  Å). (b) Example of the refinements of the diffraction pattern of  $\text{ErMn}_2\text{D}_2$  collected in the two-phase region. This plot shows the data obtained at  $T = 200$  K showing observed (crosses), calculated (upper line) and difference (bottom line) plots. The positions of the Bragg peaks are shown as ticks (from top to bottom: hexagonal  $\text{ErMn}_2\text{D}_2$ ,  $\text{ErD}_{2\pm x}$ ,  $\text{Er}_2\text{O}_3$ , monoclinic  $\text{ErMn}_2\text{D}_2$ ). (Reprinted with permission from [69]. Copyright (2006) Elsevier).

The following conclusions can be made concerning influence of hydrogen on the crystallographic and magnetic properties of the intermetallics.

The crystal structures of the intermetallic alloys of the orthorhombic  $\text{TiNiSi}$  undergo various transformations on the hydrogenation, including a) formation of internally deformed orthorhombic  $\text{CeNiSnD}$  and  $\text{TbNiSnD}$ ; b)  $\text{TiNiSi} \rightarrow \text{ZrBeSi}$  type rebuilding of the orthorhombic structure into a hexagonal one; c)  $\text{TiNiSi} \rightarrow \text{ZrNiAl}$  type rebuilding accompanied by a contraction of the unit cell volume. For the monohydrides  $\text{CeNiSnD}$  and  $\text{TbNiSnD}$  volume expansion is small, just 1.5% and 2.8%, which corresponds to a rather small value of 1.0–1.3 Å<sup>3</sup>/at.D. For  $\text{TbNiSiD}_{1.78}$  lattice expansion is much higher, 9.9%, while specific expansion reaches a typical for the metal hydride value 2.8 Å<sup>3</sup>/at.H close to 3 Å<sup>3</sup>/at.H.

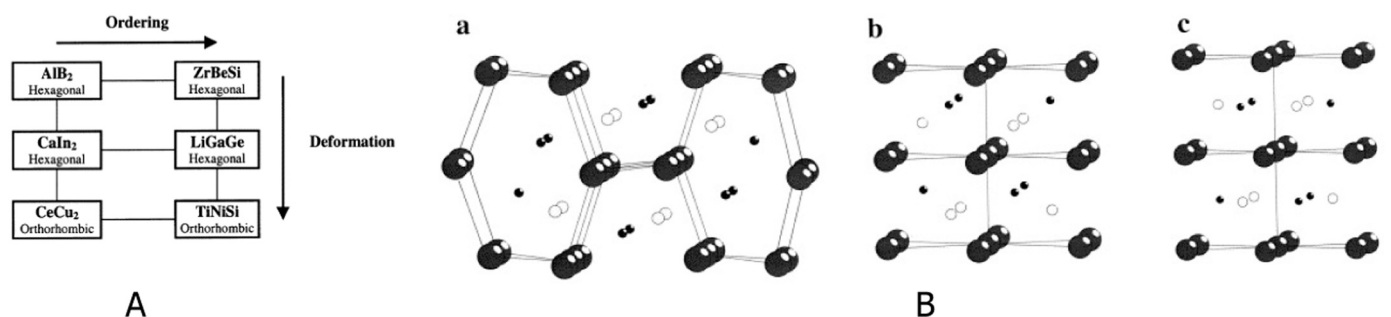
Magnetic ordering changes on the hydrogenation (see Table 1 for the details). Most frequently magnetic moment carried by the rare earth decreases while the temperature of magnetic ordering decreases as well. For the Kondo semimetal  $\text{CeNiSn}$  the hydrogenation induces antiferromagnetism in  $\text{CeNiSnD}$ .

**4.3.2.1. Magnetic ordering in  $\text{TbNiSnD}$  [71].** The neutron powder diffraction study of the crystal and magnetic structure of  $\text{TbNiSnD}$  shows that the deuteride crystallizes with the filled  $\text{TiNiSi}$  type structure (S.G.  $Pnma$ ;  $a = 7.0354(1)$ ;  $b = 4.20662(9)$ ;  $c = 8.1808(2)$  Å at 293 K;  $a = 7.008(1)$ ;  $b = 4.1954(6)$ ;  $c = 8.153(1)$  Å at 2 K) and is canted ferromagnetic below 10 K. The deuterium atoms occupy the tetrahedral  $\text{Tb}_3\text{Ni}$  sites. The Tb atoms carry a magnetic moment of



**Fig. 13.** (a) Selected possible hydrogen sites in  $\text{RTX}$  hydride with  $\text{TiNiSi}$  type structure.  $R$ : large atoms.  $T$ : dark medium atoms.  $X$ : light medium atoms. Interstitial sites:  $R_3T$  site ( $\bullet$ ),  $R_3TX$  ( $\blacksquare$ ),  $R_2TX$  ( $\ast$ ) and  $R_2TX$  ( $\blacktriangle$ ). The coordinations polyhedra are outlined. (b) Structure of  $\text{TbNiSiD}_{1.78}$  illustrated as D ( $\text{Tb}_3\text{Ni}$ ) tetrahedra sharing  $\text{Tb-Tb}$  edges and  $\text{Ni}$  corners [70]. Si (shaded isolated atoms) does not form bonds with hydrogen. Tb atoms are white and Ni atoms shaded. (Reprinted with permission from [70]. Copyright (2001) Elsevier).

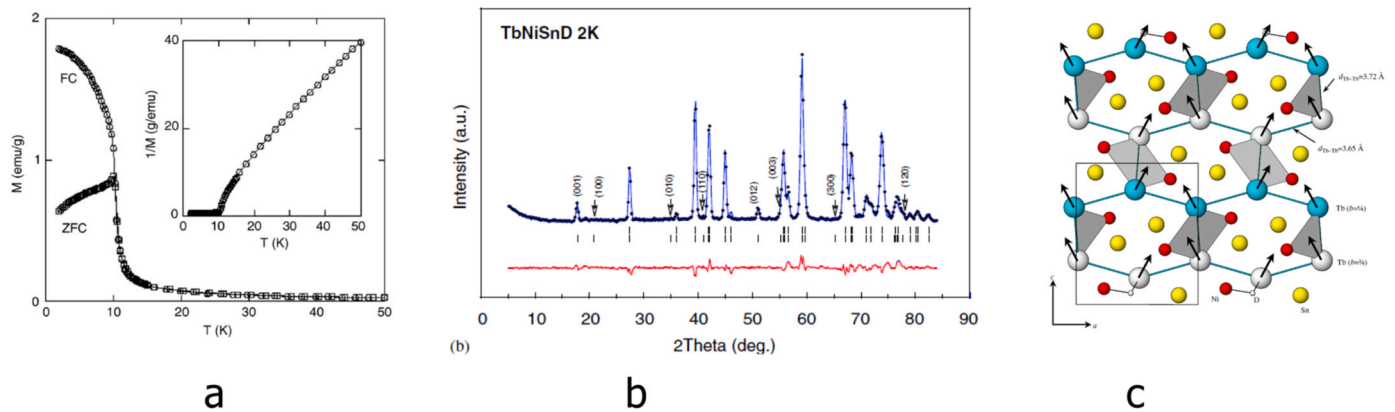
6.6(1)  $\mu_B$ /Tb. They order antiferromagnetically along the  $a$ -axis and ferromagnetically along the  $c$ -axis. No significant magnetic moment on the Ni atoms has been observed by neutron diffraction. Insertion of deuterium in  $\text{TbNiSn}$  causes dramatic changes of the original magnetic structure of  $\text{TbNiSn}$  and a reduction of the magnitude of the Tb magnetic moment (Fig. 14, Fig. 15 and Fig. 16).



**Fig. 12.** (A) An overview of the related structure types which for the ABC stoichiometries include  $\text{ZrBeSi}$ ,  $\text{LiGaGe}$  and  $\text{TiNiSi}$  types. (B) Comparison of the different structure types (a)  $\text{TiNiSi}$ , (b)  $\text{LiGaGe}$  and (c)  $\text{ZrBeSi}$ . The latter two have AA stacking of close-packed layers of  $R$  and differ by displacement of the smaller  $T$  and  $X$  (light) atoms, whereas the former is more distorted. (Reprinted with permission from [70]. Copyright (2001) Elsevier).

**Table 1**  
Crystallographic and magnetic properties of the ternary intermetallic compounds and their hydrides.

	Structure type and symmetry	Unit cell dimensions, Å	V, Å <sup>3</sup>	Type of magnetic ordering	Ordering Temperature, K	Size of magnetic moment on R atoms, μ
ErMn <sub>2</sub>	MgZn <sub>2</sub> , hexagonal, <i>P6<sub>3</sub>/mmc</i>	5.29416 8.64435 (350 K)		Ferromagnetic	15	7.72
ErMn <sub>2</sub> D <sub>2</sub>	MgZn <sub>2</sub> , hexagonal, <i>P6<sub>3</sub>/mmc</i> (350 K) Derivative from MgZn <sub>2</sub> , monoclinic, <i>C2/m</i> (150 K)	5.56787 9.07088 9.61247(9) 5.57558(4) 9.07102(8) β = 90.5451(5)	486.140(7)	Ferrimagnetic	205	9.6
CeNiSn	TiNiSi, orthorhombic, <i>Pnma</i>	7.54(1) 4.599(6) 7.639(9)	264.89	-	-	-
CeNiSnD	Deformed TiNiSi, orthorhombic, <i>Pna2<sub>1</sub></i>	7.238(4) 8.446(4) 4.398(2)	268.86	AF, with ferromagnetic sheets ordering + - -	5.1	1.37
TbNiSn	TiNiSi orthorhombic, <i>Pnma</i>	7.02204(8) 4.42515(5) 7.62601(9)	236.97	AF sine-modulated	18.5	8.9
TbNiSnD	TiNiSi orthorhombic, <i>Pnma</i>	7.0354(1) 4.20662(9) 8.1808(2)	242.112(9)	Canted ferromagnetic ordering	10	6.6
HoNiSn	TiNiSi orthorhombic, <i>Pnma</i>	7.0682(9) 4.4396(6) 7.6433(9)	239.85	AF sine-modulated	3	8.27
HoNiSnD <sub>0.67</sub>	ZrNiAl Hexagonal, <i>P6<sub>3</sub>2m</i>	7.24197(10) 3.93514(7)	238.31	AF sine-modulated	2.6	5.36
TbNiSi	TiNiSi, orthorhombic, <i>Pnma</i>	6.8715(23) 4.1439(17) 7.1508(35)	203.62(25)	AF sine-modulated	16	8.28
TbNiSiD <sub>1.78</sub>	ZrBeSi hexagonal, <i>P6<sub>3</sub>/mmc</i>	4.02895(7) 7.9585(2) (100 K)	223.76	AF	8	8.71



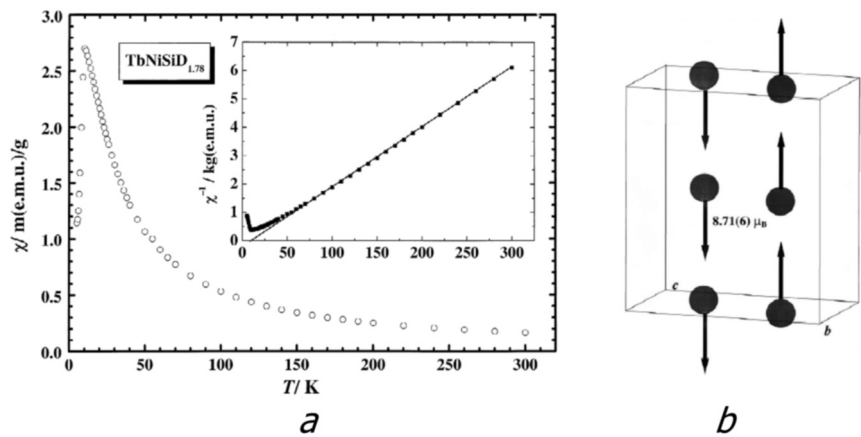
**Fig. 14.** (a) The zero-field-cooled (ZFC) and field-cooled (FC) magnetization curves show ordering of the magnetic structure proceeding below 10 K; (b) Neutron powder diffraction pattern measured below the ordering temperature at 2 K have a contribution from the nuclear and magnetic (shown by arrows) structures; (c) TbNiSnD has a canted ferromagnetic structure with a ferromagnetic component parallel to the *c*-axis and an antiferromagnetic component parallel to the *a*-axis. (Reprinted with permission from [71]. Copyright (2006) Elsevier).

**4.3.2.2. Magnetic ordering in TbNiSiD<sub>2-x</sub> [70,72].** The intermetallic compound TbNiSi adopts the orthorhombic TiNiSi-type structure (S.G. *Pnma*) [70]. Deuteration/hydrogenation) leads to increased symmetry from the orthorhombic to the hexagonal ZrBeSi-type structure (*P6<sub>3</sub>/mmc*) for TbNiSiD<sub>1.78</sub> [70]. By powder neutron diffraction, the magnetic ordering of TbNiSi was determined to be sinusoidal wave modulated below the Néel temperature at 16 K, and square-modulated below 9 K [72]. The magnetic moment of Tb (8.28 μ) is aligned along **b**.

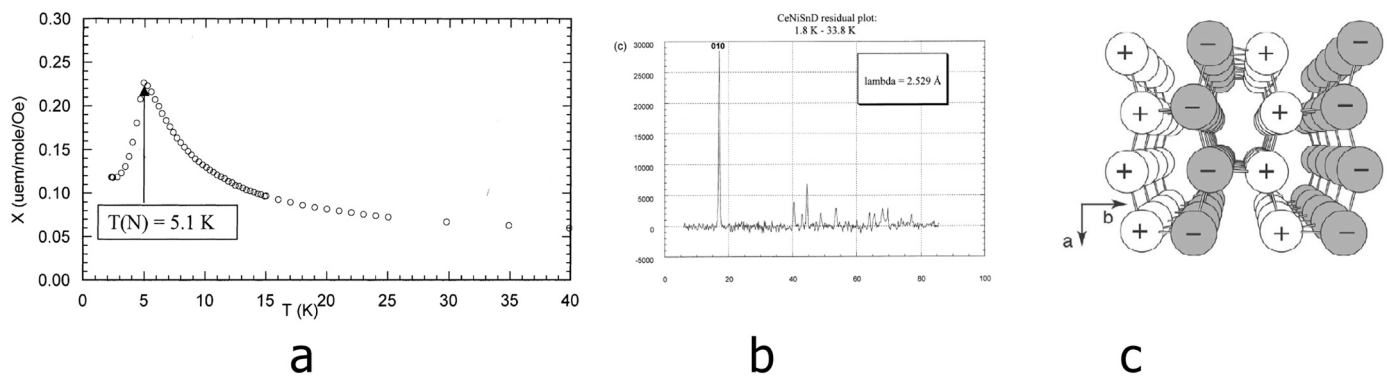
Below the magnetic ordering temperature of 10 K, at 2 K TbNiSiD<sub>1.78</sub> a small orthorhombic distortion of the lattice takes

place. The refined unit-cell dimensions at 2 K (S.G. *Pnma*) are *a* = 7.9505(2), *b* = 4.02502(14), *c* = 6.9823(2) Å. The magnetic moments of Tb are 8.71(6) μ, and are ordered antiferromagnetically along **a**.

**4.3.2.3. Magnetic ordering in CeNiSnD [73].** Neutron diffraction and magnetic measurements were used to investigate the changes in magnetic properties of the Kondo semimetal compound CeNiSn after charging with deuterium gas under moderate conditions. Charging with deuterium led to the formation of the monodeuteride CeNiSnD with filled TiNiSi structure. This compound gives rise to long-range antiferromagnetic ordering below *T<sub>N</sub>* = 5.1 K. The ordered magnetic



**Fig. 15.** (a) Magnetic susceptibility of TbNiSiD<sub>1.78</sub> in the applied field 1000 Oe as a function of temperature showing a magnetic ordering below 4 K. (Inset) Inverse magnetic susceptibility revealing a Curie-Weiss behaviour above 70 K. (b) The magnetic structure of TbNiSiD<sub>1.78</sub> with the magnetic moments of Tb atoms aligned along  $b$ . (Reprinted with permission from [72]. Copyright (2002) Elsevier).



**Fig. 16.** (a) Temperature dependence of the magnetic susceptibility for CeNiSnD peaking at  $T_N = 5.1$  K. (b) Difference plot of the neutron diffraction diagrams of the compound CeNiSnD at 33.8 K and at 2 K showing peaks of magnetic diffraction. (c) Stacking of the ferromagnetic layers along  $[010]$  in the magnetic structure of CeNiSnD. Ce moments are aligned along  $[001]$ . Opposite spins are indicated as + and - . (Reprinted with permission from [73]. Copyright (2006) Elsevier).

moment value equals  $1.37(3) \mu_B$  per Ce atom, the moments pointing along the  $[001]$  direction. No significant ordered magnetic moments on the Ni atoms have been observed by powder neutron diffraction.

**4.3.2.4. Magnetic ordering in HoNiSnD<sub>0.67</sub> [74,75].** The deuteration of HoNiSn intermetallic compound causes a change of the crystal structure from orthorhombic TiNiSi-type to the hexagonal ZrNiAl-type and a small volume contraction,  $-0.64\%$ . The atomic structure of HoNiSnD<sub>0.67</sub> is isostructural with chemically similar lower deuterides RNiInD<sub>0.5-0.6</sub> formed by La and Nd [76]. D atoms are partially filling the available Ho<sub>3</sub>Ni tetrahedra.

The magnetic ordering of HoNiSnD<sub>0.67</sub> is antiferromagnetic with the Néel temperature  $2.6(3)$  K. Initial intermetallic compound HoNiSn is antiferromagnet as well, and its ordering temperature,  $3.0$  K, remains practically unchanged on the hydrogenation.

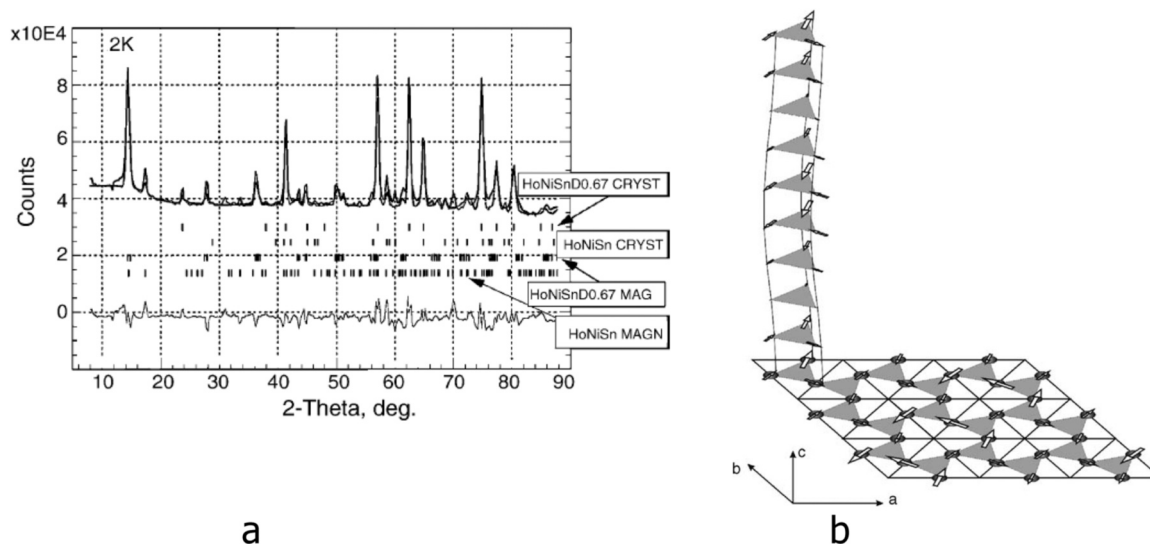
However, inserted hydrogen atoms cause change of the type of magnetic ordering. The reflections of magnetic origin, observed on the neutron diffraction pattern at 2 K (see Fig. 17a) are indexed by the propagation vector  $k_1 (= 0.34, 0.34, 0.117)$  for HoNiSnD<sub>0.67</sub> and  $k_2 (= 0.395, 0.330, 0)$  for HoNiSn. The Ho magnetic moment associated with  $k_1$  equal to  $5.36(5) \mu_B$ , lies in the basal plane  $(0\ 0\ 1)$  and forms a sine modulated structure (see Fig. 17b). The three moments in the crystal unit cell form noncollinear order. The magnetic moment on Ho atoms decreases as compared to the intermetallic HoNiSn ( $8.5 \mu_B$ ).

Other studied RENiSn based hydrides include LaNiSnH<sub>2</sub>, PrNiSnH, NdNiSnH and SnNiSnH<sub>0.6</sub> [77-79].

### 4.3.3. Multi-peak Thermal Desorption of deuterium in the YFe<sub>2</sub>-D<sub>2</sub> system [60]

The study of the YFe<sub>2</sub>-H<sub>2</sub> system attracted much attention regarding the influence of hydrogen absorption on the magnetic properties of RFe<sub>2</sub> compounds ( $R = \text{rare earth}$ ), its multi-plateau thermodynamic behaviour and the large variety of crystal structures for different hydrogen contents. The parent compound YFe<sub>2</sub> crystallizes in the cubic C15 Laves phase structure (S.G.  $Fd\bar{3}m$ ). On increasing hydrogen content, different single phase-hydrides such as YFe<sub>2</sub>H<sub>1.3</sub> (tetragonal  $I\bar{4}$ ), YFe<sub>2</sub>H<sub>1.75</sub>, (cubic  $I\bar{4}3m$ ), YFe<sub>2</sub>D<sub>1.9</sub> (tetragonal  $I\bar{4}$ ), YFe<sub>2</sub>H<sub>2.6</sub> (cubic  $Fd\bar{3}m$ ), YFe<sub>2</sub>H<sub>3.5</sub> (rhombohedral  $R\bar{3}m$ ), YFe<sub>2</sub>H<sub>4.2</sub> (monoclinic  $Pc$ ) and YFe<sub>2</sub>H<sub>5</sub> (orthorhombic  $Pmn2_1$ ) are formed [80-83]. In some cases, their neutron diffraction patterns display superstructure lines due to hydrogen ordering and the induced displacement of Y and Fe atoms to accommodate D-atoms [81]. The case of YFe<sub>2</sub>H<sub>4.2</sub> compound has received particular attention because a large isotopic effect is observed on the magnetic properties of the YFe<sub>2</sub>(H,D)<sub>4.2</sub> system [84,85]. Replacing deuterium with hydrogen leads to a sharp increase in the abrupt magnetostrictive transition, which has been attributed to the large cell volume difference between hydride and deuteride.

*In situ* neutron diffraction experiments in the YFe<sub>2</sub>-H<sub>2</sub> system have been focused on the formation and thermal ordering of YFe<sub>2</sub>HD<sub>1.3</sub>, YFe<sub>2</sub>D<sub>1.75</sub> and YFe<sub>2</sub>D<sub>1.9</sub> [81] as well as the thermal stability of YFe<sub>2</sub>D<sub>x</sub> ( $1.3 \leq x \leq 4.2$ ) deuterides [60]. For the latter studies, the aim was to shed new light on the multi-peak TD spectrum observed in this system and in general in AB<sub>2</sub> compounds [86,87]. Neutron



**Fig. 17.** (a) Observed and calculated neutron diffraction patterns of  $\text{HoNiSnD}_{0.67}$  at 2 K. Circles represent the experimental data, solid lines are for the calculated profile and at the bottom the difference between the observed and calculated data is shown. Vertical bars indicate the nuclear and magnetic peaks. 2 crystallographic structures and 2 magnetic structures were accounted during the refinements: (b) Magnetic structure of  $\text{HoNiSnD}_{0.67}$ . Only carrying magnetic moment Ho atoms are shown with a direction of the magnetic moments indicated by arrows.

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diffraction experiments were performed at D1B@ILL beamline during thermal desorption of  $\text{YFe}_2\text{D}_{4.2}$  compound. The sample was put in a stainless steel holder that was placed in a furnace with vanadium heating elements. The sample holder was heated from RT to 650 K with a ramp of 0.5 K/min and connected through a vacuum line equipped with a Pirani gauge and a rotary pump [60].

Fig. 18 shows the evolution of neutron diffraction patterns (Fig. 18c) and TD data (Fig. 18a) during thermal desorption of  $\text{YFe}_2\text{D}_{4.2}$ . On heating up to 360 K, the primitive monoclinic structure (S.G.  $Pc$ ) of this compound transforms, before the onset of deuterium desorption, to the high-symmetry cubic C15 structure (S.G.  $Fd\bar{3}m$ ) as result of an order/disorder transition (Fig. 18c). With further heating, C15 diffraction peaks shift to higher angles indicating the shrinking of C15 unit-cell volume due to deuterium desorption (Fig. 18b). The C15 structure is retained throughout the desorption process, but in the temperature range of 430–480 K, where superstructure lines belonging to  $\text{YFe}_2\text{D}_{1.9}$  and  $\text{YFe}_2\text{D}_{1.75}$  are observed. In addition, the line widths of the diffraction peaks undergo significant broadening from 350 to 445 K and from 495 to 560 K. Indeed, it has been shown that the temperatures at which these structural discontinuities occur are strongly correlated to the singularities of the TD spectrum (Fig. 18a). The TD spectrum could be deconvoluted in seven desorption events, each one related to deuterium release as a consequence of two-phase transformation between deuterides. This *in situ* study proved that the multi-peak TD spectra observed in  $\text{YFe}_2\text{D}_x$  deuterides are due to the many different phases that occur in this system as a function of deuterium content. In the past, multi-peak spectra in  $\text{AB}_2$  compounds were attributed to different local H environments in the C15 lattice, a fact that is questioned by the present study.

#### 4.3.4. Ternary hydrides of $\text{Nd}_6\text{Fe}_{13}\text{Si}$ and $\text{Mo}_2\text{FeB}_2$ types

NPD studies of  $\text{Nd}_6\text{Fe}_{13}\text{GaD}_{12.3}$  showed that deuterium atoms nearly completely occupy four different types of interstices, comprising  $\text{Nd}_2\text{Fe}_4$  octahedra and three types of tetrahedra,  $\text{Nd}_4$ ,  $\text{Nd}_3\text{Fe}$  and  $\text{Nd}_2\text{Fe}_2$  [88].  $\text{Pr}_6\text{Fe}_{13}\text{AuD}_{13}$  has a similar nuclear structure with deuterium atoms occupying various types of interstitial sites bonded with Pr and Fe atoms. A collinear antiferromagnetic ordering of the four Fe and the two Pr sublattices for the uncharged

$\text{Pr}_6\text{Fe}_{13}\text{Au}$  compound transforms into a collinear ferromagnetic structure of the deuteride [89].

An antiferromagnet  $\text{Nd}_2\text{Ni}_2\text{Mg}$  with the Neel temperature  $T_N \sim 19$  K forms a  $\text{Nd}_2\text{Ni}_2\text{MgH}_8$  hydride where the  $T_N$  becomes significantly decreased, to just 1 K [90]. The Hydrogenation of  $\text{Nd}_2\text{Ni}_2\text{Mg}$  leads to a pronounced deformation of the constituting it  $\text{AlB}_2$  and CsCl type blocks. A dynamic distribution of hydrogen atoms and no formation of covalent Ni-H complex ions are concluded for the structure of  $\text{Nd}_2\text{Ni}_2\text{MgH}_8$  (the results will be published separately).

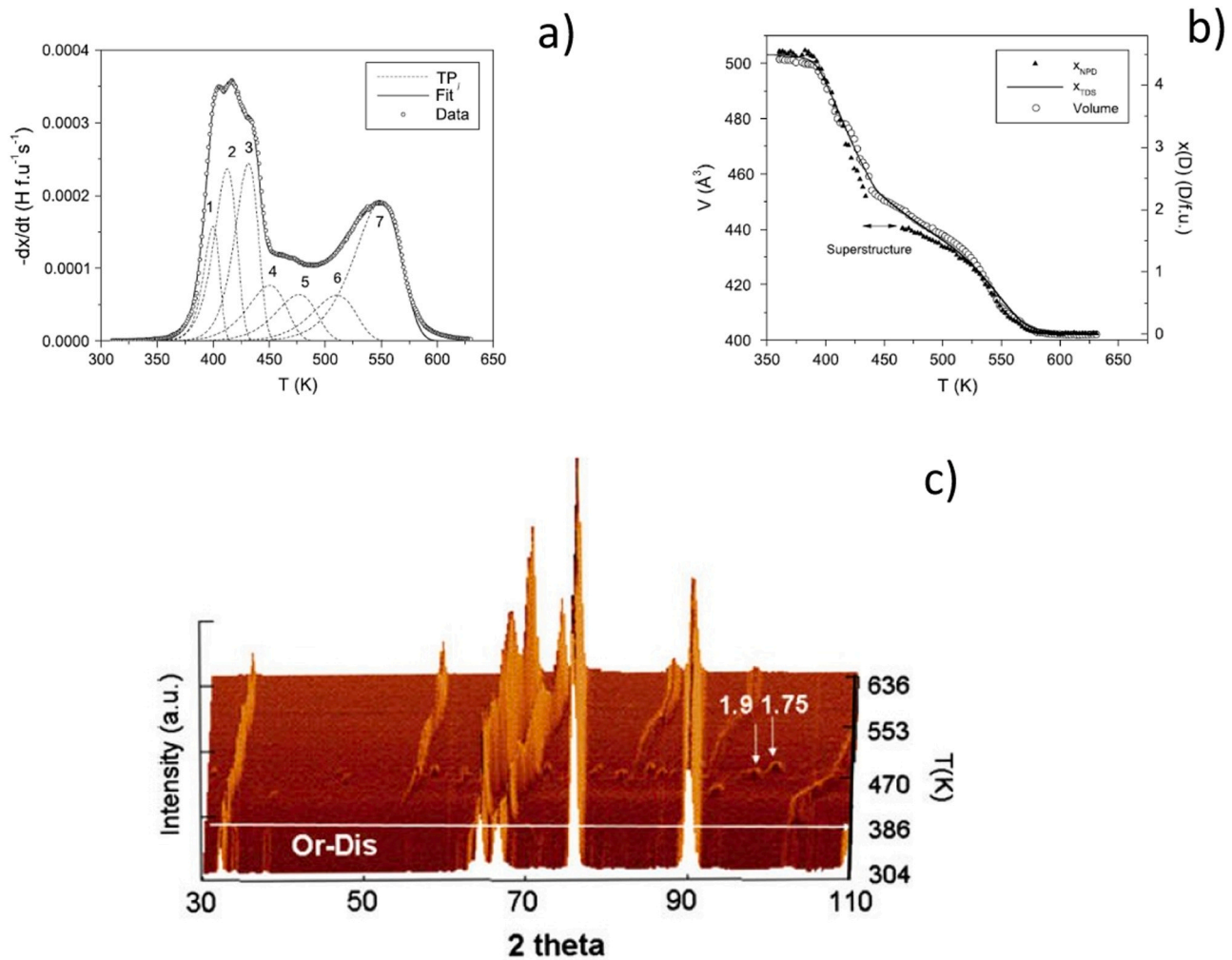
## 5. Hydrogen pressure-dependent transformations

### 5.1. Studies at pressures 1–100 bar $D_2$

#### 5.1.1. $\text{AB}_5\text{-H}_2$ systems

5.1.1.1.  $\text{LaNi}_{4.70-4.78}\text{Sn}_{0.22-0.30}\text{-H}_2$  system [91,92]. A specially designed cell for *in situ* studies (the studied sample is kept in a quartz glass capillary) by using synchrotron X-ray powder diffraction (SR XRD) in  $\text{H}_2$ -atmosphere has been used to monitor phase-structural transformations during hydride formation and decomposition in the advanced hydrogen storage alloy  $\text{LaNi}_{4.70-4.78}\text{Sn}_{0.22-0.30}$ . The cell was attached to a metal hydride hydrogen storage unit developed at IFE providing hydrogen gas at convenient constant temperature-dependent pressures. Temperature cycling between RT and 100 °C at  $\text{H}_2$  pressures of 2–3 bar was employed to  $\text{LaNi}_{4.70}\text{Sn}_{0.30}\text{H}_{\sim 6}$  which resulted in a reversible hydrogen absorption–desorption process (see Fig. 19). A large temperature-dependent homogeneity range of the  $\beta$ -hydride and a small hysteresis in the formation–decomposition of the  $\beta$ -hydride were revealed (Fig. 20).

5.1.1.2.  $\text{La}_{0.5}\text{Ce}_{0.5}\text{Ni}_4\text{Co-H}_2$  system [93]. A hexagonal  $\text{CaCu}_5$ -type  $\text{La}_{0.5}\text{Ce}_{0.5}\text{Ni}_4\text{Co}$  intermetallic alloy was studied as H storage material in as-cast condition using PCT and *in situ* SR XRD studies performed between 263 and 353 K. A significant shrinking of the unit cells takes place on a substitution of La with Ce. SR XRD showed the formation of two hexagonal hydrides, a  $\gamma$ -trihydride  $\text{La}_{0.5}\text{Ce}_{0.5}\text{Ni}_4\text{CoH}_3$ , and a saturated  $\beta$ -hexahydride  $\text{La}_{0.5}\text{Ce}_{0.5}\text{Ni}_4\text{CoH}_{\sim 6}$ . In addition, an  $\alpha$ -H solid solution was experimentally observed. Hysteresis of hydrogen



**Fig. 18.** Simultaneous acquisition of thermal desorption (TD) spectrum and neutron-diffraction (ND) patterns of  $\text{YFe}_2\text{D}_{4.2}$ . Heating rate 0.5 K/min. a) TD spectrum of  $\text{YFe}_2\text{D}_{4.2}$  deconvoluted in seven desorption events assigned to seven phase transformations. b) Evolution of C15 unit cell-volume (left-axis) and deuterium content in  $\text{YFe}_2$  (right-axis) with temperature. Deuterium contents are evaluated from Rietveld refinement of ND data and integration of TD spectrum. c) 3D representation of ND patterns with temperature. The arrows indicate superstructure lines denoting the formation of intermediate deuterides  $\text{YFe}_2\text{D}_{1.9}$  and  $\text{YFe}_2\text{D}_{1.75}$ . (Reprinted with permission from [60] Copyright (2008) Elsevier).

absorption and desorption between the flat single plateau isotherms measured during the PCT measurements was rather small, with no indication of the formation of the thermodynamically stable  $\gamma$ -(La,Ce)(Ni,Co) $_5\text{H}_3$ . The SR XRD revealed (see Fig. 20) that (a) the relative abundance of the  $\gamma$ -hydride is higher during the desorption compared to the absorption; its relative amount reaches 50 wt% in maximum; (b) anisotropic strains in all constituent phases during both hydrogen absorption and desorption are not developed; (c) the pressure of the formation of the  $\gamma$ -phase was found to be strongly temperature-dependent. The formation of large amounts of metastable  $\gamma$ -(La,Ce)(Ni,Co) $_5\text{H}_3$  hydride in quasiequilibrium experimental conditions applied during the SR XRD experiments and during reaching the thermodynamic equilibrium between the constituent phases can be responsible for the reduced hysteresis of the PCT isotherms and absence of noticeable anisotropic strains in the constituent phases.

### 5.1.2. Over-stoichiometric $\text{LaNi}_{5+x}$ - $\text{LaNi}_6$ - $\text{H}_2$ system

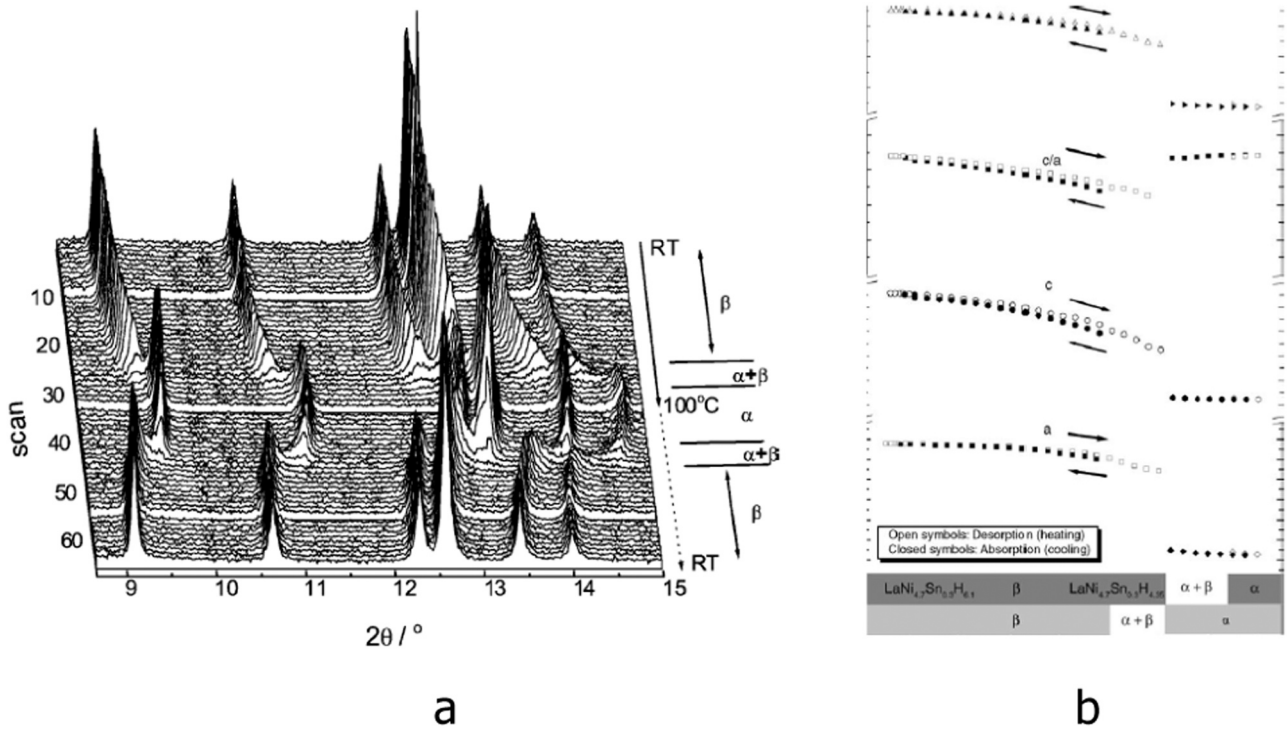
*In situ* neutron diffraction studies were performed on highly over-stoichiometric  $\text{LaNi}_{5+x}$  with  $x = 1.0$  -  $\text{LaNi}_6$  alloy - during the monitoring of desorption PCT curve from its deuteride at room temperature.

Experiments were performed at D1B@ILL at deuterium pressures between 0.01 and 10 MPa using a stainless steel sample holder and

the Sieverts' apparatus shown in Fig. 6 (Section 3). Contrary to the stoichiometric composition  $\text{LaNi}_5$  which crystallizes in  $\text{CaCu}_5$  structure-type, the crystal structure of over-stoichiometric  $\text{LaNi}_{5+x}$  alloys is better described in the  $\text{TbCu}_7$ -type structure in which La atoms are partially substituted by  $c$ -axis oriented Ni-dumbbells which can be thus described by the formula  $\text{La}_{1-y}\text{Ni}_{5+2y}$ . A highly-overstoichiometric  $\text{La}_{0.88}\text{Ni}_{5.24}$  alloy (roughly corresponding to the  $\text{LaNi}_6$  composition) was synthesized by using mechanochemical treatment and exists as a metastable phase [94]. The PCT desorption curve is displayed in Fig. 21 and is compared to that of the hydride of stoichiometric  $\text{LaNi}_5$  intermetallic alloy. In contrast to the  $\text{LaNi}_5$ - $\text{H}_2$  system, the highly overstoichiometric compound shows no clear plateau pressure and a maximum H storage capacity of 4.9 at.D/unit cell at 10 MPa  $\text{D}_2$ . However, as displayed in Fig. 21, a phase transition from  $\beta$ -deuteride to  $\alpha$ -solid solution shows a reduced  $\alpha + \beta$  coexistence region. This transition exhibits variable  $\alpha$ - and  $\beta$ -lattice volumes and a minor value of a discrete lattice expansion ( $\Delta V/V = 3.5 \text{ vol.}\%$ ) as compared to the stoichiometric  $\text{LaNi}_5$  ( $\Delta V/V = 21.5 \text{ vol.}\%$ ).

### 5.1.3. Laves type $\text{AB}_2$ deuterides with C14, C15 and C36 structures

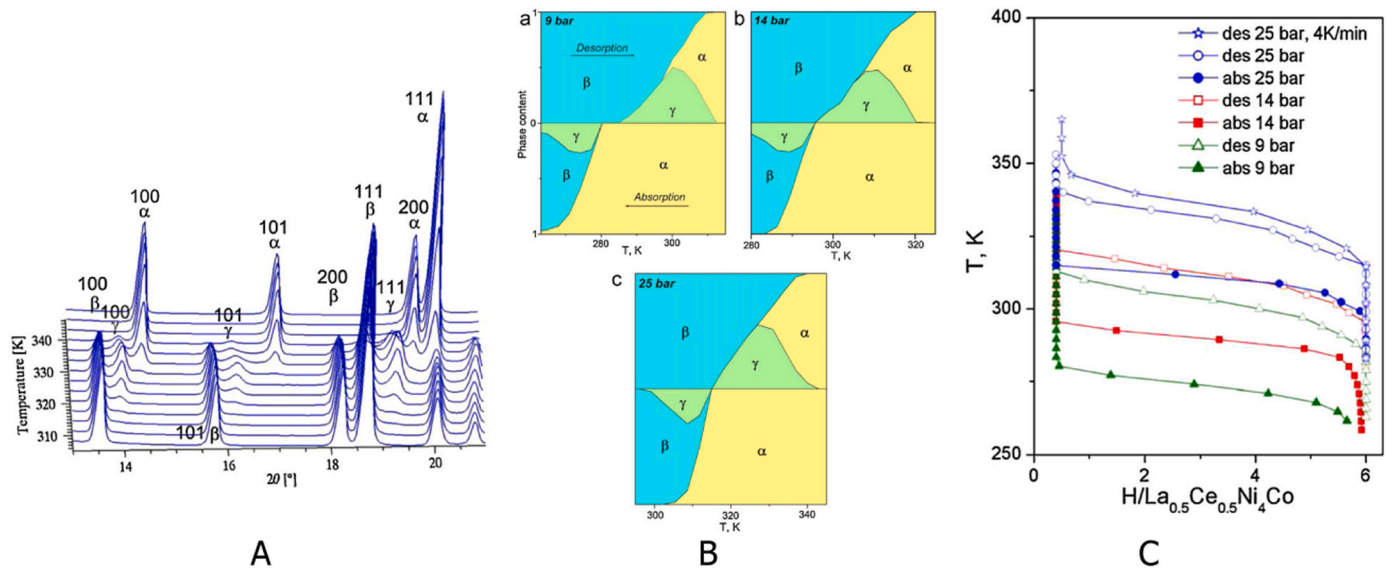
Hydride structures based on the Laves phases of composition  $(\text{Sc,Ti,Zr})\text{T}_2$  ( $T$  = transition metal of the first row) with C14  $\text{MgZn}_2$ , C15  $\text{MgCu}_2$  and C36  $\text{MgNi}_2$  structure types have been studied in



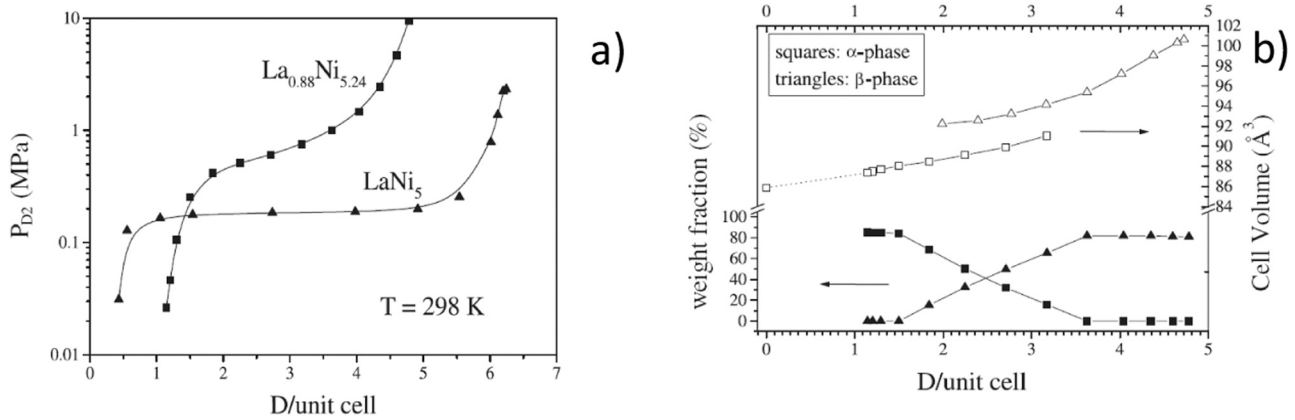
**Fig. 19.** (a) Powder diffraction profiles during transformation of  $\beta$ -phase to  $\alpha$ -phase on heating the saturated with hydrogen sample from 25° to 100°C in H<sub>2</sub> gas and subsequent cooling to 25 °C to form back the  $\beta$ -phase.(b) Variations of the unit-cell dimensions of  $\beta$ -hydride and  $\alpha$ -phase as a function of temperature. (Reprinted with permission from [92]. Copyright (2005) Elsevier).

detail [95–102]. Laves phase structures contain three different kinds of tetrahedral interstices:  $R_2T_2$ ,  $RT_3$  and  $T_4$ . When hydrogen content in the hydrides does not exceed  $H/RT_2 = 2.8$ –3.0, the  $R_2T_2$  sites are the only ones being filled (TiFe<sub>1.16</sub>V<sub>0.84</sub>D<sub>2.1</sub> [95], ZrMoFeD<sub>2.6</sub> [96],  $\lambda_1$ -ScFe<sub>2</sub>D<sub>2.5</sub> [97],  $\lambda_1$ -ScFe<sub>1.5</sub>Si<sub>0.5</sub>D<sub>2.1</sub> [97],  $\lambda_3$ -ScFe<sub>2</sub>D<sub>2.9</sub> (the only studied MgNi<sub>2</sub> C36 type hydride, D atoms fill Sc<sub>2</sub>Fe<sub>2</sub> sites of 6 types [98]), ZrV<sub>2</sub>D<sub>1.5</sub> [99]). When the hydrogen content increases up to about the saturation limit, the hydrogen atoms become accommodated in the

$RT_3$  and  $T_4$  interstices, which have less space and are less attractive in their surrounding for H atoms. Thus, in the ZrV<sub>2</sub>D<sub>2.8</sub> deuteride [99], a partial filling of the ZrV<sub>3</sub> interstices (5% of the total number of D atoms) was observed. In the ZrV<sub>2</sub>D<sub>4.5</sub> deuteride the main part of deuterium atoms in excess of  $D/ZrV_2 = 2.8$  is located in the ZrV<sub>3</sub> interstices, and this raises the fraction of the filled ZrV<sub>3</sub> interstices to  $\sim 1/3$ . This can be related to the progressing H content expansion of the unit cells, which causes the calculated values of the radii of the



**Fig. 20.** (A) Temperature evolution of the *in situ* SR XRD pattern during hydrogen desorption from the  $\beta$ -La<sub>0.5</sub>Ce<sub>0.5</sub>Ni<sub>4</sub>CoH<sub>6</sub> hydride. Miller indices of the diffraction peaks for three constituent phases are shown. (B) Phase-structural composition of the La<sub>0.5</sub>Ce<sub>0.5</sub>Ni<sub>4</sub>Co–H<sub>2</sub> system from the refinements of the SR XRD data during isobaric desorption (upper part; heating) and absorption (lower part, cooling) with a heating/cooling rate of 2 K/min under isobaric conditions, 9 bar H<sub>2</sub> (a), 14 bar H<sub>2</sub> (b) or 25 bar H<sub>2</sub> (c). (C) Calculated isobars of hydrogen desorption and absorption built on the basis of the refined SR XRD data. (Reprinted with permission from [93]. Copyright (2010) Elsevier).



**Fig. 21.** a) PCT deuterium desorption isotherms for stoichiometric  $\text{LaNi}_5$  and over-stoichiometric  $\text{LaNi}_6$  ( $\text{La}_{0.88}\text{Ni}_{5.24}$ ) alloys measured at 298 K. b) Evolution of weight fraction (full symbols) and unit-cell volumes (empty symbols) of  $\alpha$ - and  $\beta$ -phases for overstoichiometric  $\text{LaNi}_6$  during *in situ* monitoring by neutron diffraction when performing PCT desorption studies.

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$\text{ZrV}_3$  interstices to exceed  $0.4 \text{ \AA}$ . This creates chemical conditions suitable for stable metal hydride formation:  $d_{M-1} > r_M + 0.4 \text{ \AA}$  [100].

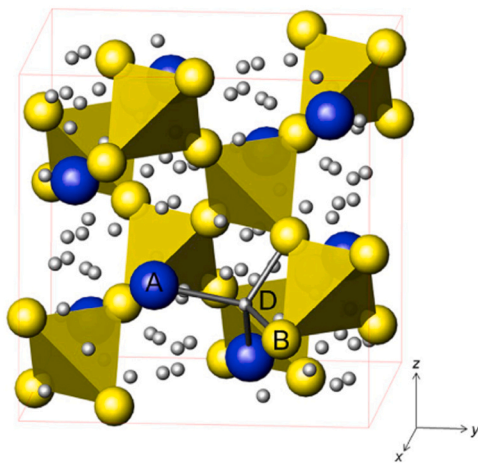
In the C14 type  $\text{ZrVFeD}_{3,6}$ ,  $\text{ZrVCuD}_{4,0}$  and  $\text{ZrVNbD}_{5,4}$  deuterides [101] in addition to the preferentially occupied  $\text{Zr}_2(\text{V},\text{T})_2$  interstices, deuterium atoms also partially fill  $\text{Zr}(\text{V},\text{T})_3$  and  $(\text{V},\text{T})_4$  interstices: 3% and 0%, 4% and 7%, 3% and 0%, respectively, of the total number of hydrogen atoms. All filled interstices have radii exceeding  $0.4 \text{ \AA}$ .

One typical example of the C15 type hydride is a multi-component  $(\text{Ti},\text{Zr},\text{V})(\text{Ni},\text{Mn},\text{Fe},\text{V})_2\text{D}_{3 \pm x}$  which for the alloy's composition  $\text{Ti}_{0.2}\text{Zr}_{0.8}\text{La}_{0-0.05}\text{Ni}_{1.2}\text{Mn}_{0.7}\text{V}_{0.12}\text{Fe}_{0.12}$  showed occupancy of the A sites not only by Ti and Zr but also by V and exclusive filling by D atoms of the  $\text{A}_2\text{B}_2$  tetrahedral sites (see Fig. 22) [102].

A number of ordered superstructures of D atoms in the metal lattice are formed in the C15 types hydrides at low temperatures [93]. When  $\text{D}/\text{AB}_2$  does not exceed 4 at.H/f.u., various types of ordering in the  $\text{A}_2\text{B}_2$  tetrahedra are realised. Interestingly, for the highest H/D content reached for  $\text{ZrV}_2\text{D}_6$ , the  $3 \text{ Zr}_2\text{V}_2 + 3 \text{ ZrV}_3$  ordering is realised [103], as in the expanded crystal lattice  $\text{ZrV}_3$  tetrahedra become equally attractive for the accommodation of D atoms as the  $\text{Zr}_2\text{V}_2$  tetrahedral sites.

#### 5.1.4. Ti(Fe,Mn)- $\text{H}_2$ system [55]

TiFe-based alloys have recently regained considerable interest for reversible hydrogen storage near room temperature due to their



**Fig. 22.** The crystal structure of the cubic  $\text{AB}_2\text{D}_{2.9}$  deuteride [102]. D atoms partially, 24%, fill 96 g sites shown as small spheres. A, B and D represent Zr/Ti/V (8a), Ni/Mn/Fe/V (16d), and D atoms, respectively. Spatial framework of  $\text{B}_4$  tetrahedra is shown. (Reprinted with permission from [102]. Copyright (2019) Elsevier).

relatively high hydrogen content (1.87 wt% and  $105 \text{ kg}_\text{H}/\text{m}^3$  for  $\text{TiFeH}_2$ ) and their low cost as compared to other intermetallics [104]. The equilibrium hydrogen pressure can be finely tuned by a partial chemical substitution to adapt the operating conditions to the specific applications [105]. The crystal structure of  $\text{TiFe}_{0.90-x}\text{Mn}_x$  alloys ( $x = 0, 0.05$  and  $0.10$ ) and their deuterides has been determined by *in situ* neutron diffraction, while recording Pressure-Composition Isotherms at room temperature [55]. Experiments were performed at D1B@ILL at deuterium pressures from 0.02 to 9 MPa with the help of a Sieverts' apparatus (see Fig. 6). Fig. 23a shows the evolution of 2D diffraction patterns for  $x = 0.05$  during a whole desorption/absorption PCI measurement. Three phases were detected, and their amount determined (Fig. 23b):  $\alpha$ -cubic solid solution (S.G.  $\text{Pm}\bar{3}m$ ),  $\beta$ -orthorhombic mono-deuteride (S.G.  $\text{P}222_1$ ) and  $\gamma$ -orthorhombic dideuteride (S.G.  $\text{Cmmm}$ ). Mn substitution does not significantly change the volumetric expansion during the phase transitions, affecting only the deuterium content for the  $\gamma$  phase and the variation of cell volume for the  $\beta$  phase. Volume expansions are 11–15% for the  $\beta$  phase and 17% for the  $\gamma$  phase, with respect to the dehydrogenated  $\alpha$  phase (Fig. 23c). This parameter is of high industrial relevance and should be properly considered for the integration of TiFe-type alloys in hydrogen tanks.

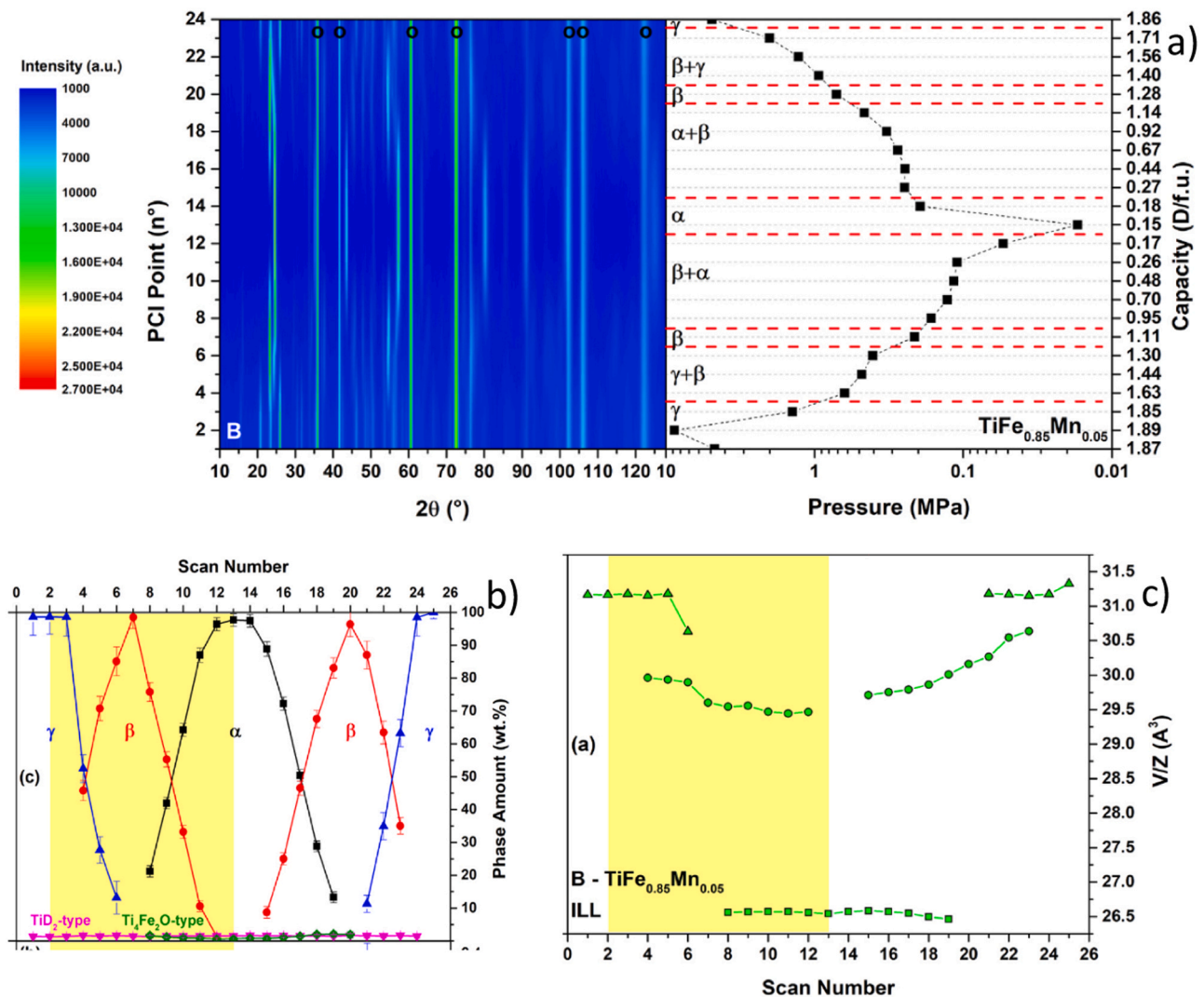
## 5.2. Studies of metal-H systems at high pressures up to 1000 bar $\text{D}_2$

### 5.2.1. $\text{CeNi}_5\text{-D}_2$ [39,40]

The initial deuteration of a single phase  $\text{CeNi}_5$  intermetallic with  $\text{CaCu}_5$  type structure ( $a = 4.88415(7)$ ;  $c = 4.00016(7) \text{ \AA}$ ) was performed on D1B at ILL without a preliminary activation of the alloy by applying high deuterium pressures. A solid solution of deuterium in  $\text{CeNi}_5$ ,  $\alpha\text{-CeNi}_5\text{D}_{\sim 0.5}$ , was formed first which extended until 600 bar  $\text{D}_2$ . At higher deuterium pressures the  $\alpha\text{-CeNi}_5\text{D}_{\sim 0.5} \rightarrow \beta\text{-CeNi}_5\text{D}_{\sim 6.3}$  transformation took place. During the first deuteration, the midpoint of the  $\alpha \rightarrow \beta$  plateau at  $20^\circ \text{C}$  was observed at 632 bar  $\text{D}_2$ ; maximum D content in the saturated beta deuteride, according to the volumetric measurements, equals 6.5 at.D/f.u. at  $\sim 800$  bar  $\text{D}_2$ . During the deuterium desorption, the midpoint of the plateau pressure for the  $\beta \rightarrow \alpha$  transformation was observed at a significantly lower pressure, 117 bar  $\text{D}_2$  at  $60^\circ \text{C}$  (see Fig. 24a). Variation of the NPD pattern as related to the pressure of  $\text{D}_2$  desorption is shown in Fig. 24b. During the second deuteration, a mid-plateau pressure for the  $\alpha \rightarrow \beta$  transformation at  $30^\circ \text{C}$  was significantly decreased, to just 237 bar  $\text{D}_2$ . (Fig. 25).

A small difference between the data of the volumetric measurements and the results of the crystal structure refinements can be





**Fig. 23.** a) *In situ* neutron diffraction 2D patterns (counter plot,  $\lambda = 1.286 \text{ \AA}$ ) and PCI curve at RT during deuterium desorption/absorption at D1B@ILL between 0.02 and 9 MPa for  $\text{TiFe}_{0.85}\text{Mn}_{0.05}$  alloy. Diffraction lines from the stainless steel sample holder are marked by a circle (o). b) amount of  $\alpha$ ,  $\beta$  and  $\gamma$  phases, c) Normalized cell volume for phases  $\alpha$  ( $\square$ ),  $\beta$  (O) and  $\gamma$  ( $\Delta$ ).

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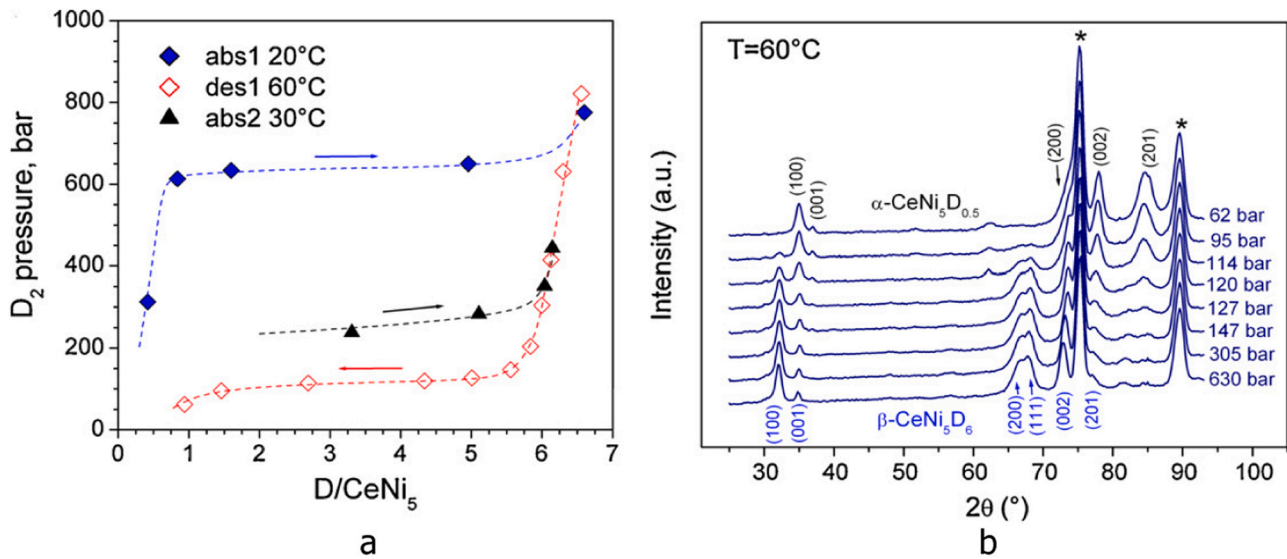
caused by the presence of undistinguishable (by neutron diffraction) amounts of the  $\beta$ -deuteride (ca. 5%) in the sample already at 613 bar  $\text{D}_2$ . Such small quantities are difficult to confidently identify using the NPD data with a lower resolution, as at the D1B instrument. However, the high flux beam allowed fast collection of a single data set, on a time scale of approximately 2 min, allowing kinetic measurements of the processes of phase-structural transformations caused by absorption or desorption of deuterium gas.

NPD data showed a complete transformation into the  $\beta$ -deuteride at  $\text{D}_2$  pressure of 776 bar. During the deuteration, the unit cell expands by 30.1%, with  $\Delta a/a = 10.0\%$  and  $\Delta c/c = 7.5\%$ . Volume expansion during the formation of the  $\beta$ - $\text{CeNi}_5\text{D}_{6.3}$  is significantly larger as compared to  $\text{LaNi}_5\text{D}_{6.5}$  (~24%) [106]. No superstructure peaks were observed in the NPD pattern; thus, the structure of the  $\text{CeNi}_5\text{D}_{6.3}$  deuteride was described in the space group  $P6/mmm$ . However, a possible hydrogen ordering in the structure resulting in the doubling of the unit cell and lowering of the symmetry cannot be excluded, similar as in the  $\beta$ - $\text{LaNi}_5\text{D}_{6+x}$  [107].

Three different interstitial sites were found to be partially filled by D atoms, including  $\text{Ce}_2\text{Ni}_2$  tetrahedra;  $\text{Ce}_2\text{Ni}_1\text{Ni}_2$  sites (half of the octahedron  $\text{Ce}_2\text{Ni}_1\text{Ni}_2$ ) and tetrahedra  $\text{Ni}_4$ . Refined deuterium content of 6.3(2) at.D/f.u. agrees well with the volumetric

measurements, 6.5 D/f.u. The values of the unit cell parameters for the saturated  $\text{CeNi}_5\text{D}_{6.3}$  deuteride,  $a = 5.374(2)$ ;  $c = 4.3012(7) \text{ \AA}$ , are very close to the corresponding values for  $\text{LaNi}_5\text{D}_6$ . This is in contrast to a significant difference between the sizes of the unit cells for  $\text{LaNi}_5$  and  $\text{CeNi}_5$ , where a significant contraction of the unit cells takes place on a replacement of La by Ce. Thus, a decrease of the Ce valence from a partial  $4^+$  state in  $\text{CeNi}_5$  to a trivalent state in  $\text{CeNi}_5\text{D}_{6.3}$  is likely to take place during the hydrogenation process.

Finally, a pronounced anisotropic diffraction peak broadening upon the first hydrogenation of the  $\text{CeNi}_5$  was observed, during the  $\alpha$ - $\text{CeNi}_5\text{D}_{\sim 0.5} \rightarrow \beta$ - $\text{CeNi}_5\text{D}_{\sim 6}$  transformation. The anisotropic strain broadening is not observed in the initial period of the hydrogenation, until the  $\beta$ -deuteride appears and, after that, remains a component of the metal-hydrogen system during further desorption-absorption cycling. Such an anisotropic broadening is attributed to the misfit dislocations created during the first hydrogenation. During the refinements, the anisotropic peak broadening was modelled using the function developed by Latroche et al. [108], where it is described as originating from the local variations of the unit cell parameters  $a$  and  $c$  of the hexagonal cells, using three microstrain coefficients. The strain coefficients obtained for the  $\text{CeNi}_5\text{D}_{6.3}$  are comparable to those reported for the  $\text{LaNi}_5\text{D}_{6.5}$  deuteride [109]. As suggested in



**Fig. 24.** (a) P–C diagrams for  $\text{CeNi}_5\text{-D}_2$  system obtained by volumetric measurements during *in situ* NPD experiments; (b) evolution of NPD pattern during deuterium desorption (des1 60 °C) in  $\text{CeNi}_5\text{-D}_2$  system. The asterisks show reflections from a stainless steel liner of the high-pressure sample cell. (Reprinted with permission from [39]. Copyright (2010) Elsevier).

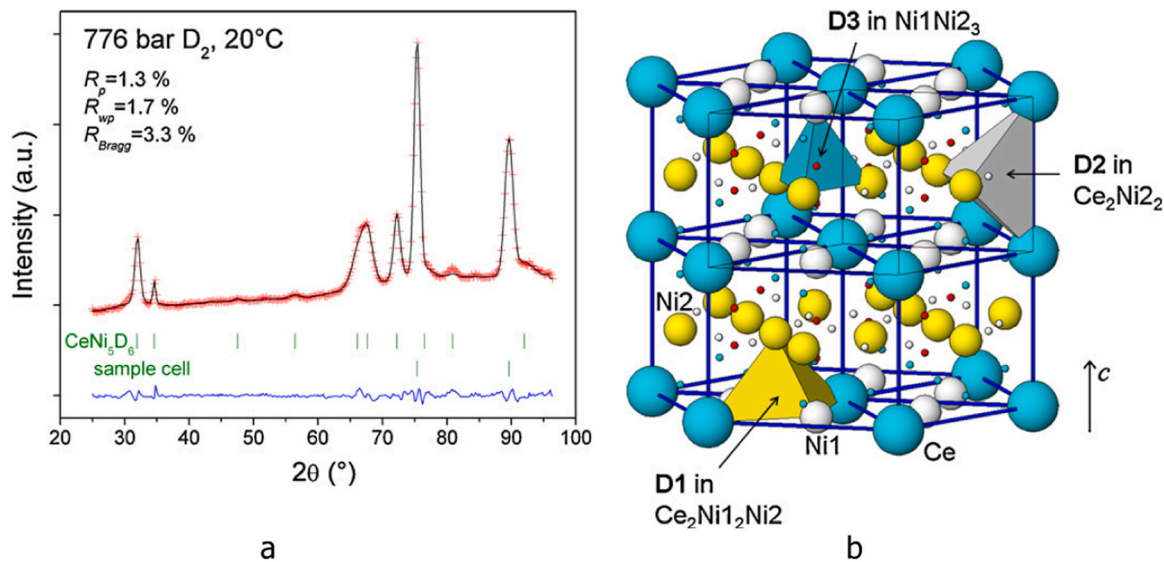
[108], a decrease in the discrete unit cell volume expansion values during the  $\alpha \rightarrow \beta$  transition in the substituted  $\text{LaNi}_{5-x}\text{M}_x$  ( $M = \text{Co}, \text{Mn}, \text{Al}, \text{Fe}, \text{Cu}, \text{Sn}$ ) or overstoichiometric  $\text{LaNi}_{5+x}$  ( $x = 0\text{-}0.5$ ) compounds leads to a reduction of the strain parameters. Thus, it is natural to propose a similar explanation for the  $\text{CeNi}_5 \rightarrow \text{CeNi}_5\text{D}_{6.3}$  transformation; here the larger, as compared to the  $\text{LaNi}_5\text{D}_{6.5}$  [108], values of the microstrain parameters developed during the deuteration of the Ce-based intermetallic are caused by significantly larger volume expansion during the deuterium absorption in  $\text{CeNi}_5$ .

### 5.2.2. $\text{ZrFe}_{1.98}\text{Al}_{0.02}\text{-D}_2$ [39]

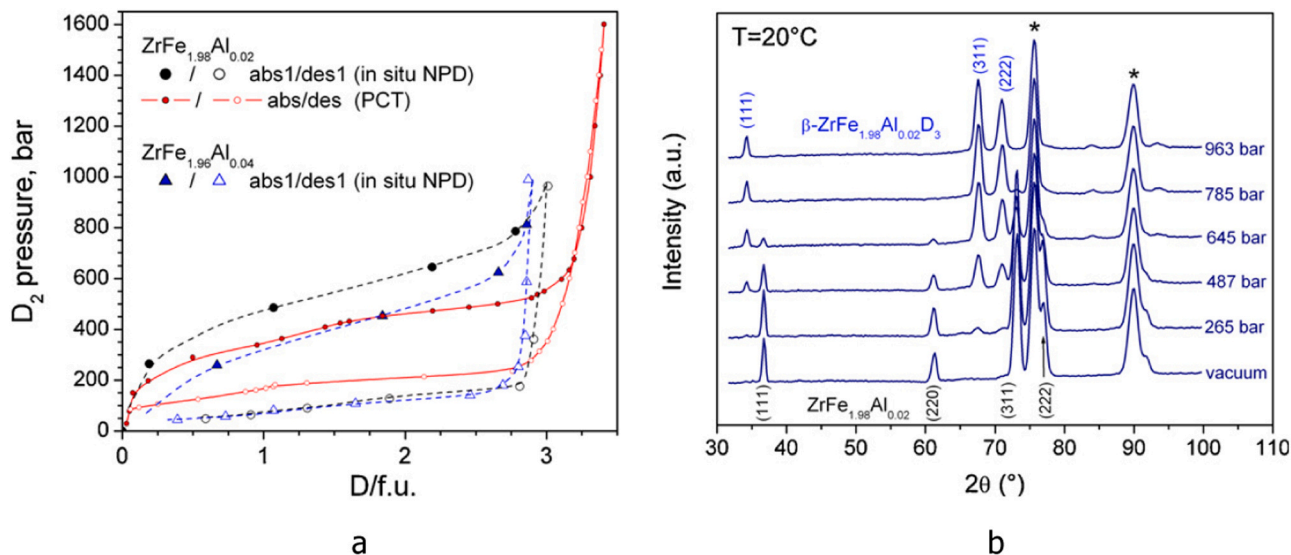
$\text{ZrFe}_2$  intermetallic compound crystallizes with C15 ( $\text{MgCu}_2$ ) Laves-type structure. SR XRD studies of the  $\text{ZrFe}_{2-x}\text{Al}_x$  alloys with  $x = 0.02$  and  $0.04$  showed that Fe substitution by Al increases the unit cell parameters from  $a = 7.072 \text{ \AA}$  for  $\text{ZrFe}_2$  [39] to  $a = 7.08492(8) \text{ \AA}$  for

$\text{ZrFe}_{1.98}\text{Al}_{0.02}$ ,  $a = 7.08658(7) \text{ \AA}$  for  $\text{ZrFe}_{1.96}\text{Al}_{0.04}$ , while the FCC crystal structure remains unchanged.

*In situ* hydrogenation-dehydrogenation studies demonstrated a very fast kinetics of H/D exchange in the  $\text{ZrFe}_{2-x}\text{Al}_x$  hydrides caused by the introduction of aluminium instead of iron, in contrast with the  $\text{ZrFe}_2$  binary Laves phase compound. This allowed a quick synthesis of the hydrogenated samples. For  $\text{ZrFe}_{1.98}\text{Al}_{0.02}$  after activation by several hydrogen absorption–desorption cycles, at 1600 bar its maximum hydrogenation capacity reached 3.4 at.D/f.u. (see Fig. 26a), close to the capacity of the  $\text{ZrFe}_2$  intermetallic alloy of 3.5 at.H/f.u. at 1800 bar  $\text{H}_2$  [39]. For the NPD studies (see Fig. 26b), the absorption was performed at room temperature without preliminary activation of the as-cast samples (crushed into the particles 0.5–1 mm in size). Because of the absence of activation and cycling, and as a result of lower synthesis pressure applied, maximum reached D storage capacity was slightly lower than in the Sieverts



**Fig. 25.** (a) Refined neutron diffraction pattern for  $\text{CeNi}_5\text{D}_{6.3}$  (observed (+), calculated (—) and difference (bottom line) profiles). (b) Crystal structure of the  $\text{CeNi}_5\text{D}_{6.3}$  deuteride. Three types of interstices occupied by D atoms including tetrahedron  $\text{Ce}_2\text{Ni}_2$ , trigonal bipyramid  $\text{Ce}_2\text{Ni}_3$  and tetrahedron  $\text{Ni}_4$  are shown. (Reprinted with permission from [39]. Copyright (2010) Elsevier).



**Fig. 26.** (a) P–C diagrams measured at room temperature for the first D<sub>2</sub> absorption–desorption cycle in the ZrFe<sub>2-x</sub>Al<sub>x</sub>-D<sub>2</sub> (x = 0.02 and 0.04) systems. The data are collected by the volumetric measurements during *in situ* NPD experiments. Absorption–desorption isotherms for the activated ZrFe<sub>1.98</sub>Al<sub>0.02</sub> were obtained during the conventional PCT measurements and are given for comparison. (b) Evolution of the NPD pattern during deuterium absorption by ZrFe<sub>1.98</sub>Al<sub>0.02</sub>. The asterisks show the peaks from the high-pressure sample cell.

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1600 bar D<sub>2</sub> experiment, reaching 3.0 at.D/f.u. for ZrFe<sub>1.98</sub>Al<sub>0.02</sub> and 2.9 at.D/f.u. for ZrFe<sub>1.96</sub>Al<sub>0.04</sub>.

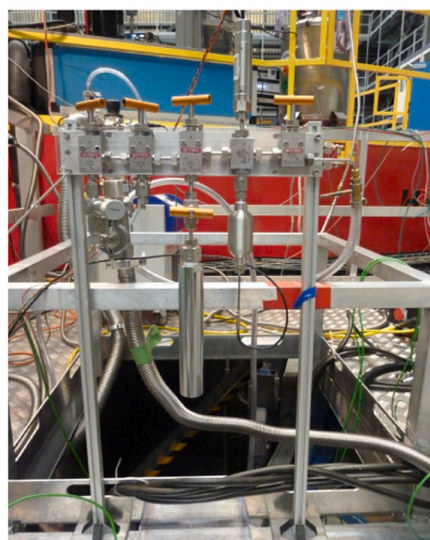
### 5.2.3. La<sub>1.09</sub>Mg<sub>1.91</sub>Ni<sub>9</sub>D<sub>9.5</sub> and La<sub>0.91</sub>Mg<sub>2.09</sub>Ni<sub>9</sub>D<sub>9.4</sub> [110]

Magnesium can replace 2/3 of La in LaNi<sub>3</sub> intermetallic compound to reach a limiting composition LaMg<sub>2</sub>Ni<sub>9</sub>. Two deuterated La<sub>3-x</sub>Mg<sub>x</sub>Ni<sub>9</sub> samples with slightly different La/Mg ratios both only marginally deviating in a composition from LaMg<sub>2</sub>Ni<sub>9</sub>, La<sub>1.09</sub>Mg<sub>1.91</sub>Ni<sub>9</sub> (sample 1) and La<sub>0.91</sub>Mg<sub>2.09</sub>Ni<sub>9</sub> (sample 2), were studied using *in situ* neutron powder diffraction.

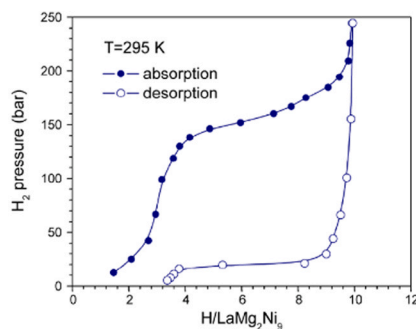
NPD characterisation of the La<sub>1.09</sub>Mg<sub>1.91</sub>Ni<sub>9</sub>D<sub>9.5</sub> (1) and La<sub>0.91</sub>Mg<sub>2.09</sub>Ni<sub>9</sub>D<sub>9.4</sub> (2) deuterides was performed at 25 bar D<sub>2</sub> (1)

and 918 bar D<sub>2</sub> (2) using the setup shown in Fig. 27a. The hydrogenation properties of the (1) and (2) are dramatically different from those for LaNi<sub>3</sub>. The Mg-containing intermetallics reversibly form hydrides with  $\Delta H_{\text{des}} = 24.0 \text{ kJ/molH}_2$  (Fig. 27c) and an equilibrium pressure of H<sub>2</sub> desorption of 18 bar at 20 °C (La<sub>1.09</sub>Mg<sub>1.91</sub>Ni<sub>9</sub>). A pronounced hysteresis of H<sub>2</sub> absorption and desorption, ~100 bar, is observed (Fig. 27b).

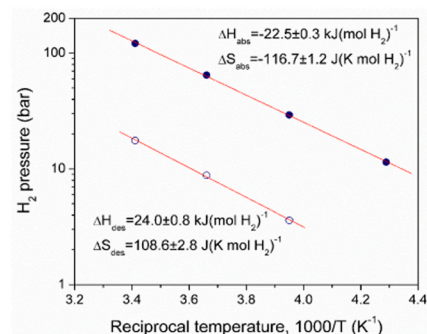
LaNi<sub>5</sub>-assisted hydrogenation of MgNi<sub>2</sub> in the LaMg<sub>2</sub>Ni<sub>9</sub> hybrid structure takes place. In the La<sub>1.09</sub>Mg<sub>1.91</sub>Ni<sub>9</sub>D<sub>9.5</sub> (1) and La<sub>0.91</sub>Mg<sub>2.09</sub>Ni<sub>9</sub>D<sub>9.4</sub> (2) ( $a = 5.263/5.212$ ;  $c = 25.803/25.71 \text{ \AA}$ ) D atoms are accommodated in both Laves and CaCu<sub>5</sub>-type slabs (see Fig. 28a). In the LaNi<sub>5</sub> CaCu<sub>5</sub>-type layer, D atoms fill three types of interstices;



a



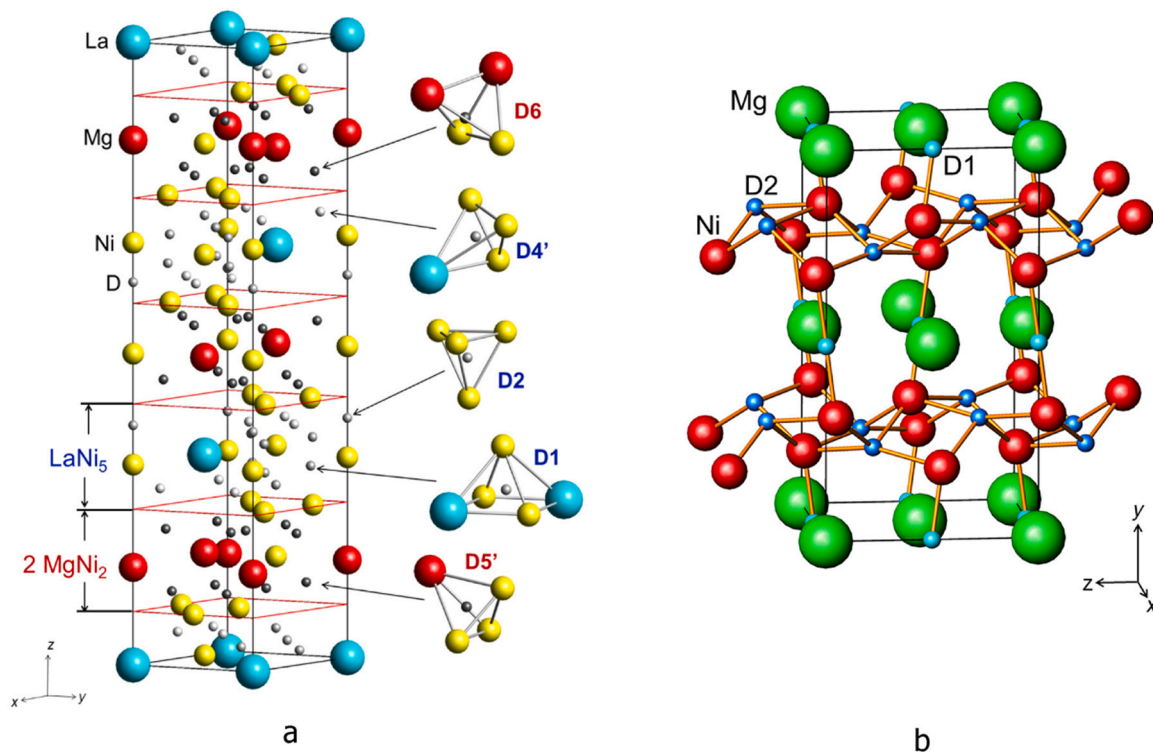
b



c

**Fig. 27.** (a) High pressure synthesis setup for the *in situ* NPD measurements at pressures up to 1000 bar D<sub>2</sub>; (b) Room temperature isotherms of hydrogen absorption and desorption; van't Hoff plots (c) for La<sub>1.09</sub>Mg<sub>1.91</sub>Ni<sub>9</sub>-based hydride. At room temperature equilibrium pressure of hydrogen absorption is ~120 bar D<sub>2</sub>, while for the desorption  $P_{\text{eq}}$  equals to ~20 bar D<sub>2</sub>.

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**Fig. 28.** (a) Crystal structure of  $\text{La}_{1 \pm 0.1}\text{Mg}_{2 \pm 0.1}\text{Ni}_9\text{D}_{9.4-9.5}$  and types of the filled interstices [110]; (b) The orthorhombic  $Cmca$  crystal structure of  $\text{MgNi}_2\text{D}_3$  [111]. Two types of sites occupied by deuterium include a  $\text{Mg}_4\text{Ni}_2$  octahedron for H1(D1) and a chair  $\text{Ni}_4$  configuration for H2(D2) located within the buckled Ni-H nets containing bended spirals  $-\text{Ni}-\text{H}_2-\text{Ni}-\text{H}_2-$ . (Reprinted from an open access article [110] distributed under the Creative Commons Attribution license (<http://creativecommons.org/licenses/by/4.0/>)). (Reprinted with permission from [111]. Copyright (2015) Elsevier).

a deformed octahedron  $[\text{La}_2\text{Ni}_4]$ ,  $[(\text{La},\text{Mg})_2\text{Ni}_2]$  and  $[\text{Ni}_4]$  tetrahedra. The overall chemical compositions can be presented as  $\text{LaNi}_5\text{H}_{5.6/5.0} + 2 * \text{MgNi}_2\text{H}_{1.95/2.2}$  showing that the hydrogenation of the  $\text{MgNi}_2$  slab proceeds at mild  $\text{H}_2/\text{D}_2$  pressure of just 20 bar. In the  $\text{MgNi}_2$  slab a partial filling by D atoms of four types of the tetrahedral interstices takes place, which include  $[\text{MgNi}_3]$  and  $[\text{Mg}_2\text{Ni}_2]$  tetrahedra.

$\text{LaNi}_5$ -assisted hydrogenation of  $\text{MgNi}_2$  is observed in the  $\text{LaMg}_2\text{Ni}_9$  hybrid structure. Formation of  $\text{LaMg}_2\text{Ni}_9\text{D}_{9.5}$  proceeds via an isotropic expansion of the trigonal unit cell. D atoms are accommodated in both Laves and  $\text{CaCu}_5$ -type slabs H atoms filling interstitial sites in both  $\text{LaNi}_5$  and  $\text{MgNi}_2$  structural fragments. Limits of Mg solubility in  $\text{LaNi}_3$  are not confined to  $\text{LaMg}_2\text{Ni}_9$  and extend to the composition  $\text{La}_{0.91}\text{Mg}_{2.09}\text{Ni}_9$  with a refined composition of the  $\text{CaCu}_5$ -type block of  $\text{La}_{0.95}\text{Mg}_{0.05}\text{Ni}_5$ . Within the  $\text{LaNi}_5$   $\text{CaCu}_5$ -type layer, D atoms fill three types of interstices; a deformed octahedron  $[\text{La}_2\text{Ni}_4]$ , and two types of tetrahedra,  $[\text{LaNi}_3]$  and  $[\text{Ni}_4]$ , to yield  $\text{LaNi}_5\text{D}_{5-5.6}$  composition. D distribution is very similar to that in the individual  $\beta$ - $\text{LaNi}_5\text{D}_7$  deuteride.

In the  $\text{MgNi}_2$  slab hydrogen atoms fill two types of tetrahedra,  $[\text{Mg}_2\text{Ni}_2]$  and  $[\text{MgNi}_3]$ . The hydrogen sublattice formed is unique and was not observed earlier in the studied structures of the Laves-type intermetallic hydrides. These include the crystal structure of  $\text{MgNi}_2\text{D}_3$  [111], [112], [113] featuring accommodation of hydrogen atoms into the  $\text{Mg}_2\text{Ni}_4$  octahedra and nickel nets, see Fig. 28b, during the formation of which metal sublattice become rebuilt.

A significant mutual influence of the  $\text{LaNi}_5$  and  $\text{MgNi}_2$  slabs causes a dramatic altering of their hydrogenation behaviours leading to:

- (a) significant decrease of the stability of the  $\text{LaNi}_5$ -type hydride;.
- (b) much easier hydrogenation of the  $\text{MgNi}_2$  slabs compared to the parent intermetallic compound;.
- (c) increased hysteresis.

## 6. Dynamic transformations in the Laves type metal-hydrogen systems

### 6.1. $\text{AB}_2\text{-D}_2$ system studied at 5/10 bar $\text{D}_2$ in the temperature window

#### 6.1.1. 300–473 K during the heating and cooling

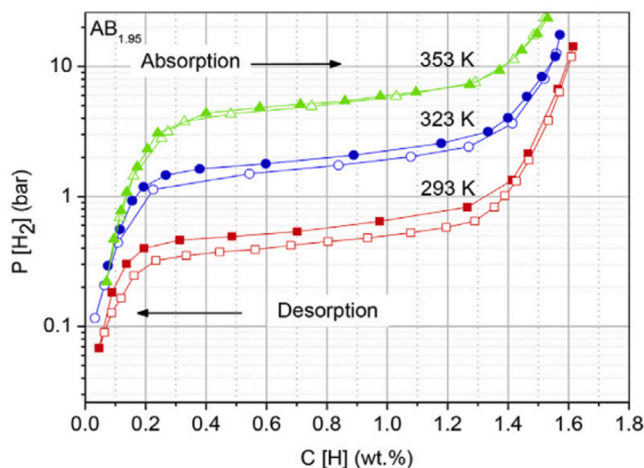
$\text{AB}_2$  Laves type intermetallics form the most abundant group of the hydride-forming alloys which contains more than 900 binary and ternary compounds. The chemistry of these compounds is very rich as they are composed from a variety of the very different in chemical behaviours elements at the A (Mg, Ca, Ti, Zr, Rare Earth Metals) and B (V, Cr, Mn, Fe, Co, Ni, Al) sites. Recently hydrides of the Laves type intermetallics as hydrogen storage materials were reviewed in [16,114].

When the ratio between the atomic radii of the A and B components has an ideal value of  $r_A/r_B = 1.225$ , these structures can be presented as a close packing of different in size spheres A and B, while the structures contain  $\text{A}_2\text{B}_2$ ,  $\text{AB}_3$  and  $\text{B}_4$  interstices present in stoichiometric ratio 12:4:1.

Broad solid solubility ranges on both A and B sites allow to optimize the performance of the hydrides – reversible storage capacity, thermodynamics and kinetics of their interaction with hydrogen and thermal stability.

One family which has been recently in focus of the research efforts can be described by a general formula  $(\text{Zr}_{1-x}\text{Ti}_x)(\text{Ni}, \text{Mn}, \text{V}, \text{Fe})_{2 \pm y}$ .

When the content of the components at the B site is fixed as  $\text{Ni}_{1.2}\text{Mn}_{0.7}\text{V}_{0.12}\text{Fe}_{0.12}$ , this allows to utilise a modification of the ratio between Zr and Ti (in a range from 88:12–78:12) and, furthermore, a ratio between B and A (between 1.90 and 2.07) to optimize thermodynamic properties of the hydrides – equilibrium pressures of hydrogen desorption and absorption. One attractive feature of the



**Fig. 29.** Isotherms of hydrogen absorption and desorption in the  $Zr_{0.80}Ti_{0.20}Ni_{1.2}Mn_{0.7}V_{0.12}Fe_{0.12}-H_2$  system measured at 293, 323 and 353 K. Reprinted from an open access article [115] distributed under the Creative Commons Attribution license (<http://creativecommons.org/licenses/by/4.0/>).

formed hydrides with H storage capacity close to 3 at.H/f.u. $AB_{2\pm y}$  (2.9–3.2) is in a very small hysteresis between the equilibrium pressures of hydrogen absorption and desorption and presence of the flat plateaux at the isotherms of  $H_2$  absorption and desorption, see Fig. 29. Please, note that with increasing temperature from 293 K to 353 K, the difference between the absorption and desorption branches of the isotherms decreases because the measured T become close to the critical temperature in the studied system (the differences with disappear at  $T = T_{critical}$ ).

The high rates of hydrogen absorption and desorption steps (See Fig. 30, a and b) are jointly governed by hydrogen diffusion and grain boundary nucleation of alpha-solid solution and beta-hydride identified as a rate limiting steps by fitting the kinetic equation  $X = 1 - \exp(-Kt^n)$  with the values of  $n = 1-1.25$ .

*In situ* neutron powder diffraction patterns were measured for the deuteride sample placed into a 6 mm SS autoclave and equilibrated at 5 bar  $D_2$  and at 10 bar  $D_2$  in separate experiments. Then the samples were slowly heated and, later, cooled down while keeping a slow linear heating and cooling rate, so a heating experiment from

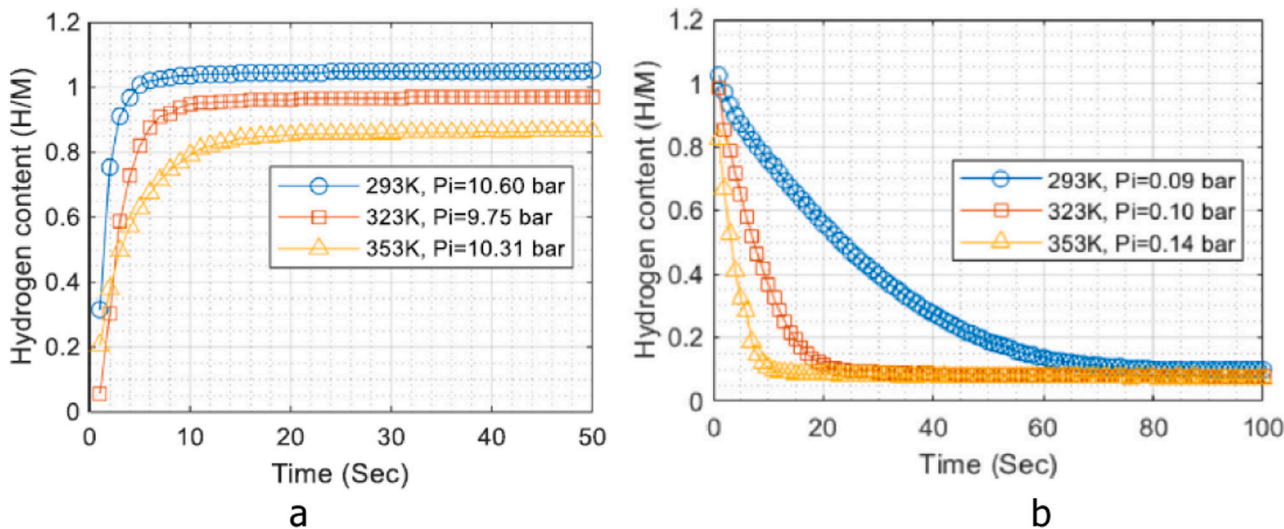
300 to 473 K lasted 4 h and 20 min. The data collection has been performed by recording individual data sets recorded every 7 min. These individual data sets were all successfully fitted by Rietveld profile refinements and gave nice fits of the data. A typical example is presented in the Figure Fig. 31a, while the Fig. 31b shows a 2-D plot of the development of the temperature-dependent development NPD pattern manifesting a reversible decomposition of the deuteride to form the initial alloy and of the rehydrogenation process.

The fitting of the NPD data allowed to reach the following conclusions. The maximum hydrogen storage capacity reached was 3.22 at.D/f.u.  $AB_{1.95}$ . H/D capacity is pressure-dependent and grows up when the pressure increases from 5 to 10 bar  $D_2$ .  $\beta$ -deuteride has a broad homogeneity range, extending between 2.45 and 3.22 at.D/ $AB_{1.95}$ . D content in  $\alpha$ -solid solution within a two-phase region is appr. 0.40–0.45 at.D/f.u.  $AB_{1.95}$ .

The collected data allowed to build the isobar diagrams of the dependence between the deuterium content D/ $AB_{1.95}$  and temperature for the system  $AB_{1.95}-D_2$  studied at two setpoint pressures of deuterium, 5 and 10 bar  $D_2$  and showed an excellent agreement (see Fig. 32) between the experimental data (red spheres – 10 bar  $H_2$ ; black rectangles – 5 bar  $H_2$ ) and modelled dependences when using thermodynamic constants for the  $Zr_{0.80}Ti_{0.20}Ni_{1.2}Mn_{0.7}V_{0.12}Fe_{0.12}H_3$  hydride,  $\Delta H_{des} = 35.25$  kJ/mol  $H_2$ ;  $\Delta S_{des} = 113.33$  J/mol  $H_2$  K.

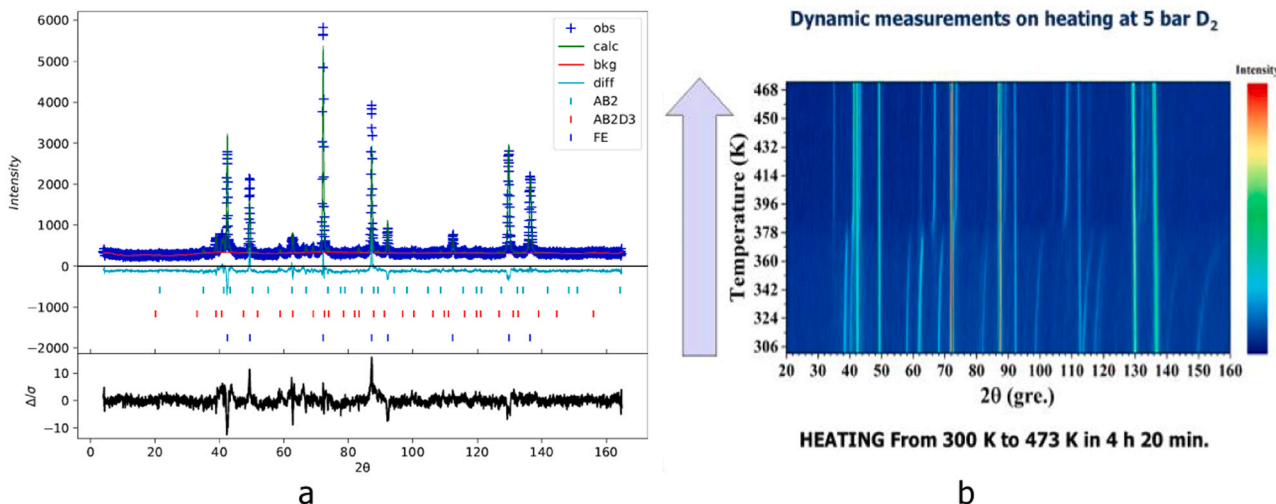
## 6.2. Reversible deuterium loading and kinetics in $MgD_2-TiD_2$ nanocomposites [36]

$MgD_2-TiD_2$  nanocomposites are outstanding materials for hydrogen storage because of their high reversible capacity (> 4 wt%), fast reaction kinetics and good cycle life at atmospheric pressure and moderate temperature, typically 573 K [118,119]. They can be efficiently prepared by mechanochemistry of Mg and Ti powders under deuterium pressure [120]. For a better understanding of these properties, *in situ* neutron diffraction experiments were carried out at D1B@ILL on two representative  $(1-x)MgD_2-xTiD_2$  samples with  $x = 0$  and 0.3 [36]. Simultaneous to the acquisition of neutron diffraction data, first thermodesorption (TD) from RT to 600 K and then PCI measurements at 548 K were performed. For the former, the deuterium outlet of the stainless steel sample holder was connected to a Baratron device and rotatory pump, while for the latter, a

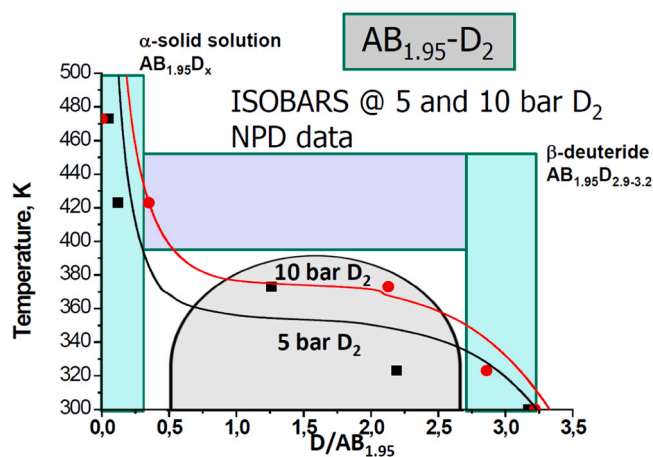


**Fig. 30.** Hydrogen absorption (a) and desorption (b) kinetic curves showing that at specified P/T conditions the absorption process is very fast and can be completed in just 9 s, while the desorption process is clearly slower.

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**Fig. 31.** (a) Refinements of an individual data set measured after equilibrating the sample at 323 K / 10 bar  $D_2$  during the cooling experiment. Collection time of the data pattern was 7 min [117]. (b) 2-D plot of the temperature-dependent variation of the diffraction pattern measured in the T range between 300 and 473 K for the sample accommodated in a SS autoclave at 5 bar  $D_2$  showing a transformation from the deuteride to the alloy.

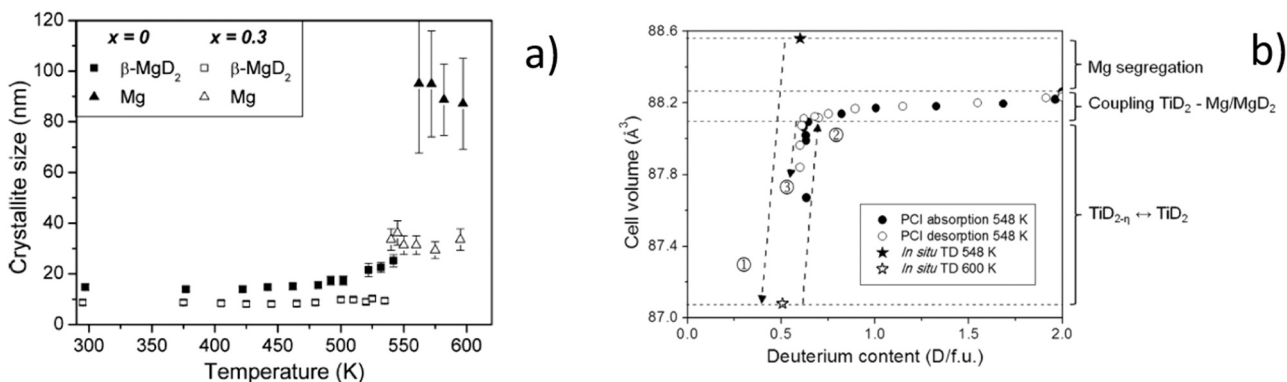


**Fig. 32.** Experimental data from the NPD studies of the phase-structural equilibria in the  $Zr_{0.80}Ti_{0.20}Ni_{1.2}Mn_{0.7}V_{0.12}Fe_{0.12} - H_2$  system forming a saturated hydride  $AB_{1.95}H(D)_{3.2}$  [117].

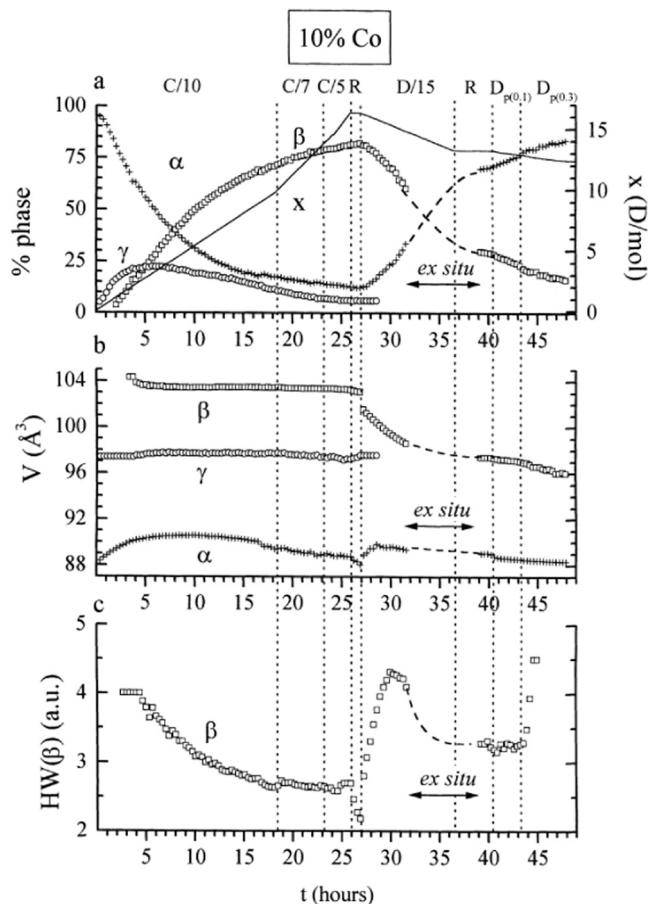
Sieverts' apparatus (see Fig. 6) was connected allowing controlled deuterium uptake/release within the pressure range 0.01–1 MPa.

Fig. 33a shows the evolution the crystallite size of tetragonal  $\beta$ - $MgD_2$  (S.G.  $P4_2/mnm$ ) and hexagonal Mg (S.G.  $P6_3/mmc$ ) phases during TD experiments on  $x=0$  and 0.3 samples. For the Ti-free sample ( $x=0$ ), the Mg phase exhibits severe grain coarsening with a crystal size of  $90 \pm 20$  nm after desorption of  $MgD_2$  above 550 K. In contrast, in the presence of  $TiD_2$  phase, grain coarsening of Mg crystallites is limited to  $30 \pm 5$  nm at the same temperature. It is concluded that, following the Zener effect,  $TiD_2$  inclusions limit the grain growth of Mg and  $MgD_2$  phases allowing for short D-diffusion paths and, therefore, fast reaction kinetics associated to the  $MgD_2$ /Mg phase transformation and long cycle-life.

Fig. 33b shows the evolution the  $TiD_2$  cell volume during the TD and PCI experiments for  $0.7MgD_2-0.3TiD_2$  sample ( $x=0.3$ ). During TD, the  $TiD_2$  cell volume first decreases from  $88.6$  to  $88.2 \text{ \AA}^3$  due to segregation of Mg out of the  $TiD_2$  structure and then to  $87.1 \text{ \AA}^3$  due to the formation of sub-stoichiometric  $TiD_{2-\eta}$ . Next, during PCI experiments, variations in the range  $88.1 - 88.3 \text{ \AA}^3$  are concomitant to the Mg/ $MgD_2$  transformation and are therefore attributed to the coherent coupling between the  $TiD_2$  phase and Mg/ $MgD_2$  interfaces.



**Fig. 33.** (a) Evolution of the crystallite size of  $\beta$ - $MgD_2$  and Mg phases during thermal desorption (TD) of  $(1-x)MgD_2-xTiD_2$  samples with  $x=0$  and 0.3. (b) Variation of  $TiD_2$  cell volume during *in situ* TD and PCI measurements. Event ① corresponds to cell volume decrease during *in situ* TD. Events ② and ③ correspond to the cell volume variations during partial absorption and desorption of deuterium in  $TiD_2$  phase at 548 K. (Reprinted with permission from [36]. Copyright (2013) American Chemical Society).



**Fig. 34.** (a) Electrochemical cycling of the 10 wt% Co electrode from a discharged initial state: (a) intensities of the  $\alpha$  (+),  $\gamma$  (O) and  $\beta$  ( $\square$ ) phases (left axis) and expected capacity  $x$  (full line) determined from coulometric measurements (right axis). (b) Cell volume  $V$  of the  $\alpha$  (+),  $\gamma$  (O) and  $\beta$  ( $\square$ ) phases. (c) Evolution of the half width of the  $\beta$  ( $\square$ ) phase. (Reprinted with permission from [63]. Copyright (1999) Elsevier).

These results suggest that the  $\text{TiD}_2$  phase facilitates H-mobility through coherent coupling between  $\text{TiD}_2$  and  $\text{Mg}/\text{MgD}_2$  phases and the presence of sub-stoichiometric  $\text{TiD}_{2-\eta}$ .

Hydrogen storage properties of Mg-based alloys and Mg-based composites depend upon the chemical nature of the additives. Various mechanisms of the favourable catalytic influence of the additives, structure and metallurgical state of the alloys have been uncovered during their phase-structural characterisation [66,113,121].

## 7. Electrochemical transformations in battery electrode materials

### 7.1. Metal hydride anodes

The use of metal hydrides as anode electrode materials of the Nickel-Metal Hydride batteries is the most developed application of the metal hydrides with the huge worldwide market of commercial Ni-MH batteries. Typical example is seen nowadays with the use of this technology as energy storage unit for a significant share of the hybrid electric vehicles (HEV) produced so far. Cheaper, safer and simpler than their competitor (Li ion batteries), the alkaline Ni-MH can be used to store electricity produced by renewable energies, intermittent in space and time. Several recent review papers present state-of-the-art status of the development of the electrode materials for the metal hydride anodes [122–124].

Layered intermetallics of  $\text{AB}_3$  and  $\text{A}_2\text{B}_7$  stoichiometries which structures can be formed via the stacking of the  $\text{AB}_5$  ( $\text{CaCu}_5$  type) and  $\text{AB}_2$  (Laves type) layers when  $\text{A}=\text{La}$  and  $\text{Ce}$ ;  $\text{B}=\text{Ni}$ , show a very unusual, strongly anisotropic expansion of the hexagonal/trigonal unit cells exclusively limited to the Laves type slabs associated with the formation of the special networks of Ni-H covalent bonds [74, 125–129]. Anisotropic expansion of the unit cells was also observed for the  $\text{RENiIn}$ -based hydrides ( $\text{RE}=\text{La}, \text{Ce}, \text{Nd}$ ) and has been associated with a different reason – formation of strong Ni-H bonds and very short H...H separations, as short as 1.56 Å for Nd compound, where the “rule of 2 Å” is violated [76, 130–133].

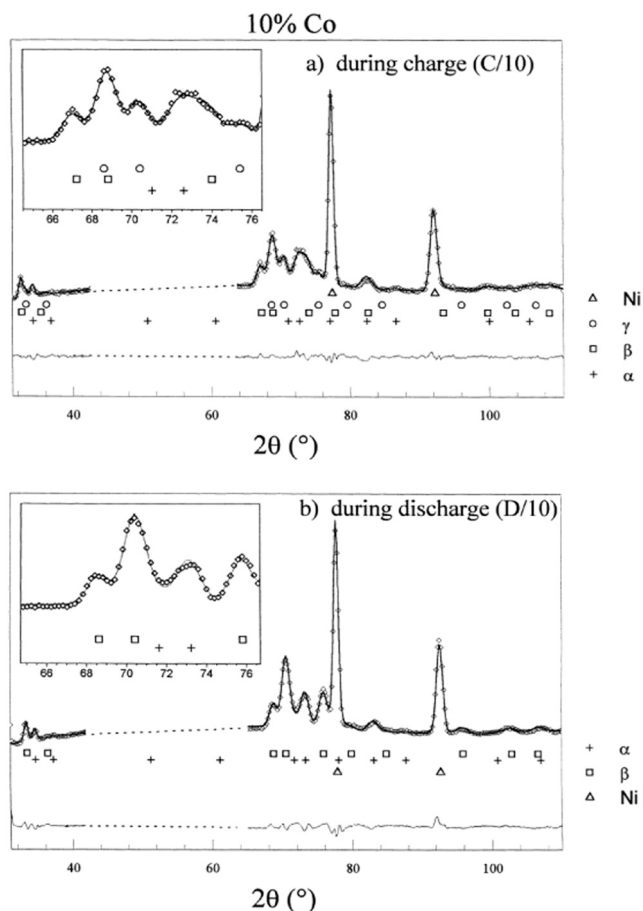
Unfortunately, hydrides of layered RE-Ni intermetallics are thermodynamically unstable and are prone to the disproportionation during their cycling in hydrogen gas. Among the materials with the highest electrochemical storage capacities are intermetallic alloys of the  $\text{AB}_3$  and  $\text{A}_2\text{B}_7$  types where Mg partially substituted lanthanum forming a solid solution without a change of the type of crystal structure. It appears that a partial replacement of RE metals by magnesium causes stabilisation of the metal sublattice. These materials offer 30% improvement of the electrochemical performance as compared to the initially developed  $\text{AB}_5$  type metal hydride anodes and have been extensively studied in the recent years [52, 134–140]. At the same time, Co-containing  $\text{AB}_5$  alloys still hold a significant share of the market.

#### 7.1.1. $\text{LaNi}_5$ -derived electrodes

$\text{LaNi}_5$ -type alloys ( $\text{CaCu}_5$ -structure type,  $P6/mmm$  space group) are the most common active material used in the negative electrode of commercial Ni-MH batteries. The proof-of-concept of Ni-MH batteries dates back to the early 1970's [141,142] but their commercialization was only possible from the 1990's through the adjustment of  $\text{LaNi}_5$  chemistry to the requirements of battery operation. Through suitable chemical substitutions, multi-elementary  $\text{LaNi}_5$ -derived alloys were developed with the aim of increasing the thermodynamic stability of hydride  $\text{LaNi}_5\text{H}_6$  while mitigating alloy deprecipitation and corrosion during electrochemical cycling [143].

*In situ* neutron diffraction experiments played a major role to characterize the electrochemical properties of the  $\text{LaNi}_5$ -type alloys [47,48,63]. Latroche et al. showed that single-substitution of Ni by Al, Mn or Co decreases the equilibrium plateau pressure and, moreover, that the cycle-life improves when the width of the  $\beta$ -hydride solution branch increases at the expense of the  $\alpha + \beta$  plateau region [47]. The latter effect was assigned to a reduction of deprecipitation phenomena, which was greatly observed in multi-substituted compounds such as  $\text{LaNi}_{3.55}\text{Mn}_{0.4}\text{Al}_{0.3}\text{Co}_{0.75}$  alloy.

This concept was latter better understood in a subsequent manuscript devoted to the influence of cobalt content in  $\text{MmNi}_{4.3-x}\text{Mn}_{0.3}\text{Al}_{0.4}\text{Co}_x$  alloy. Two alloys with cobalt content of 5 and 10 wt% ( $x = 0.36$  and  $0.69$ , respectively) were analyzed [63]. Alloy pulverization significantly decreased with cobalt content. *In situ* neutron diffraction patterns were monitored for the 10 wt% Co sample  $\text{La}_{0.51}\text{Ce}_{0.24}\text{Nd}_{0.23}\text{Pr}_{0.02}\text{Ni}_{3.61}\text{Mn}_{0.30}\text{Al}_{0.39}\text{Co}_{0.69}$  during a charging /discharging cycle with different regimes (Fig. 34a). The refinement of diffraction patterns required to consider a novel intermediate phase, denoted as  $\gamma$ , in addition to classical  $\alpha$ -intermetallic and  $\beta$ -hydride. This phase has a cell volume between that of  $\alpha$  and  $\beta$  phases while keeping same space group (Fig. 34b), while the halfwidth parameters of the  $\beta$  phase changed (Fig. 34c). The  $\gamma$  phase is metastable and is only observed during the charge step, as shown in Fig. 35. The occurrence of this intermediate deuteride implies a two-step process in the  $\alpha$  to  $\beta$  transformation and it allows the system to reduce the strains induced by the large volume expansion during phase transition. As a consequence, the deprecipitation process and subsequent corrosion effects were mitigated.



**Fig. 35.** Refined neutron diffraction pattern of  $MmNi_{3.61}Mn_{0.3}Al_{0.4}Co_{0.69}$  (10 wt% Co) alloy. (a) After 9 h of charge at  $C/10$ , (b) after 5 h of discharge at  $D/15$ , showing the experimental, calculated and difference curves. The inset shows the detailed contribution between  $64 < 2\theta < 78$  involving three phases during the charge and only two phases for the discharge. Ni lines come from both the collector grid and the counter electrode.

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Further studies by *in situ* neutron diffraction demonstrated that the occurrence of the intermediate  $\gamma$  phase depends not only on the Co content but also on the mischmetal composition and the alloy stoichiometry, *i.e.* the atomic ratio between the transition metals  $B$  and the rare earths  $A$  [49]. Thus, the probability of the appearance of the  $\gamma$  phase decreases with cobalt content but increases with  $AB_{5+y}$  over-stoichiometry ( $y$ ). Owing to the high cost of cobalt as compared to other transition metals, intensive research was conducted to reduce its content without impacting the cycle-life. This was achieved by optimizing the alloy stoichiometry, which results in the shrinking of the  $\alpha + \beta$  plateau region [35]. Moreover, by replacing Co by Fe, the formation of intermediate  $\gamma$ -phase is favoured and long cycle-life achieved. Indeed, the cycle-life of the Co-free and overstoichiometric  $AB_{5.3}$  compound  $La_{0.70}Ce_{0.22}Nd_{0.06}Pr_{0.01}Ni_{4.07}Mn_{0.64}Al_{0.19}Fe_{0.42}$  is reported to be comparable to that of equivalent Co-containing compound [50].

*In situ* neutron diffraction studies have been conducted in  $LaNi_5$ -type electrodes for determining also kinetic mechanisms. Reaction kinetics of an  $AB_5$ -type electrode of composition  $La_{0.50}Ce_{0.35}Nd_{0.12}Pr_{0.03}Ni_{3.60}Al_{0.30}Mn_{0.40}Co_{0.70}$  at high charge / discharge rates (regime over 2 C, 700 mA/g) have been analysed by Latroche et al. [45]. The active material was pulverised by solid-gas hydrogenation cycling and was afterwards surface-treated by hot-alkali etching to ensure high catalytic activity. By using high neutron flux at D20/ILL beamline and the large position sensitive detector, they succeeded to

analyse its reaction kinetics during full absorption/desorption in 1 h (Fig. 36). At such high reaction rates, out of equilibrium states were observed simultaneously in the  $\alpha$ -intermetallic and  $\beta$ -hydride phases showing that the rate limiting factor of electrode charging/discharging is the  $\alpha/\beta$  phase transformation not the deuterium diffusion coefficient within any of these phases.

### 7.1.2. Mg-containing $RE_{3-x}Mg_xNi_9 - D_2$ systems

$LaNi_5$ -type anode materials are effective, but their intrinsic capacity is restricted to  $300 \text{ mAh} \cdot \text{g}^{-1}$ . To improve the energy density, it is necessary to introduce new materials with larger H storage and electrochemical capacity. One approach is to use layered structure phases. These structures are built by stacking along the  $c$ -axis different subunits with compositions  $AB_2$  and  $AB_5$  (A: rare earths and B: transition metals). By adding different slabs as  $mAB_5 + 2AB_2$ , various stoichiometries of the general A-B phase diagram can be described such as  $AB_3$  ( $m=1$ ),  $A_2B_7$  ( $m=2$ ), and  $A_5B_{19}$  ( $m=3$ ). Costly rare-earth can be partially replaced by additions of magnesium leading to pseudo-binary or ternary systems A-Mg-B. Mg substitution has been demonstrated to be effective as it enables a lower molar mass, enhances the structural stability, and stabilizes the equilibrium pressure as well as reducing the overall cost.

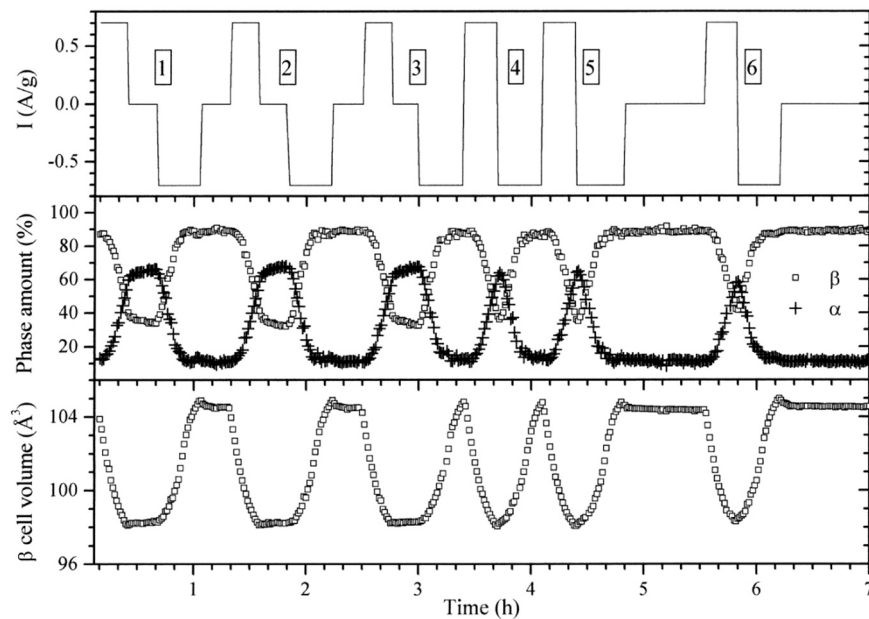
**7.1.2.1.  $La_2MgNi_9$  anode.** The electrochemical behaviour of a composite electrode made of  $La_2MgNi_9$  has been thoroughly investigated by *in situ* neutron powder diffraction using deuterated samples at different charge/discharge rates. From the data analysis, combining diffraction and PCT measurements (see Fig. 37, a, b, c), the mechanisms and the kinetic of the electrode performance have been determined. A fairly good agreement is observed between the solid-gas and electrochemical capacities (see Fig. 37a). The electrode material works by following a hydride to intermetallic transformation through a *beta* solid solution domain and a *beta* to *alpha* transformation with low capacity attributed to the *alpha* solid solution domain. During charge, heavy line broadening is however observed and might be related to the formation of an intermediate *gamma* phase as previously observed in the  $LaNi_5$  systems. The electrochemical reaction easily follows the current variations at all rates indicating absence of kinetic limitations for the hydrogen exchange in the materials when the charge-discharge rates  $C/10$  and  $C/5$  are applied.

However, for the studied bulky electrode, part of the reversible capacity is lost because of (a) formation of electrochemically stable  $\alpha$  H solid solution which contains up to 1 at. H/f.u. and (b) incompleteness of the conversion of the metal hydride anode electrode alloy into the *beta* hydride phase during the charging process. Thus, care should be taken to achieve efficient performance of the metal hydride electrodes on their scaling up [52].

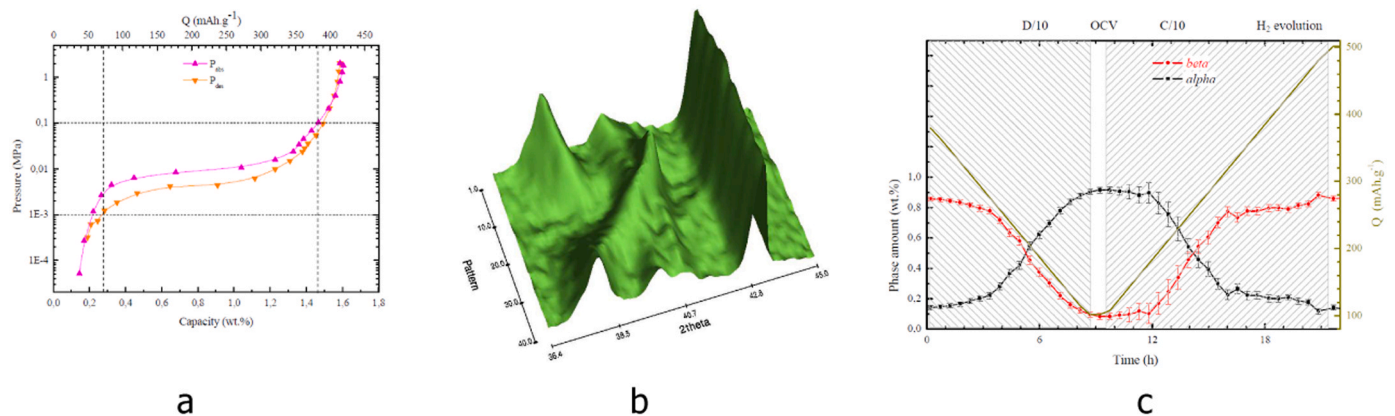
**7.1.2.2.  $LaNdMgNi_9$  anode.**  $LaNdMgNi_9$  intermetallic alloy crystallizes with a rhombohedral  $PuNi_3$  type structure and has slightly smaller unit cell parameters as compared to  $La_2MgNi_9$ . The partial replacement of La by Nd induces a volume contraction of 1.2% for the metallic lattice.  $LaNdMgNi_9$  displays a single pressure plateau in the PCT isotherm measured at 293 K, with a desorption pressure around 0.03 MPa  $H_2$  and a maximum storage capacity close to 12.9 at. H/f.u. (see Fig. 38a),(Fig. 39).

The presence of Nd in the  $La_{2-y}Nd_yMgNi_9$  alloys results in the appearance of more flat and longer plateaux, both in the case of hydrogen absorption-desorption and under electrochemical charge-discharge processes. Furthermore, increased Nd substitution for La in  $La_2MgNi_9$  promotes the formation of a more homogeneous alloy and increases the high rate dischargeability (HRD) of the electrode. The electrochemically charged alloy showed a hydrogen storage capacity of 9.8 D/f.u., 23% lower as compared to the deuteride obtained by solid-gas route at 1.6 MPa of deuterium pressure (12.9 D/f.u.). NPD studies of the saturated deuteride revealed a nonuniform





**Fig. 36.** Current density (top), phase amount (middle) and  $\beta$ -phase cell volume during galvanostatic cycling at high rate (700 mA/g) of surface-treated  $\text{La}_{0.50}\text{Ce}_{0.35}\text{Nd}_{0.12}\text{Pr}_{0.03}\text{Ni}_{3.60}\text{Al}_{0.30}\text{Mn}_{0.40}\text{Co}_{0.70}$  electrode. (Reprinted with permission from [45]. Copyright (2002) Elsevier).



**Fig. 37.** (a) Pressure-Composition-Temperature (PCT) isotherm curve for  $\text{La}_2\text{MgNi}_9$  during a full absorption/desorption cycle at 20 °C. The dotted lines stand for the electrochemical window (between 0.001 and 0.1 MPa) and define about  $310 \text{ mAh}\cdot\text{g}^{-1}$  of reversible capacity (data from [52]) whereas the total capacity at 0.1 MPa reaches  $380 \text{ mAh}\cdot\text{g}^{-1}$ ; (b) 3D view of the NPD pattern evolution as function of time during the first cycle (D/10 + C/10) of the working electrode at a current rate of  $33 \text{ mA}\cdot\text{g}^{-1}$ ; (c) Comparison between the  $\alpha$  and  $\beta$  phase amounts and the capacity  $Q$  during the in-beam charge/discharge cycle at D/10 + C/10 for the working electrode. The initial charge state  $Q$  at  $t=0$  is set at  $380 \text{ mAh}\cdot\text{g}^{-1}$ , a value determined from the PCT curve (Fig. 39a) at a pressure of 0.1 MPa. (Reprinted with permission from [52]. Copyright (2014) American Chemical Society).

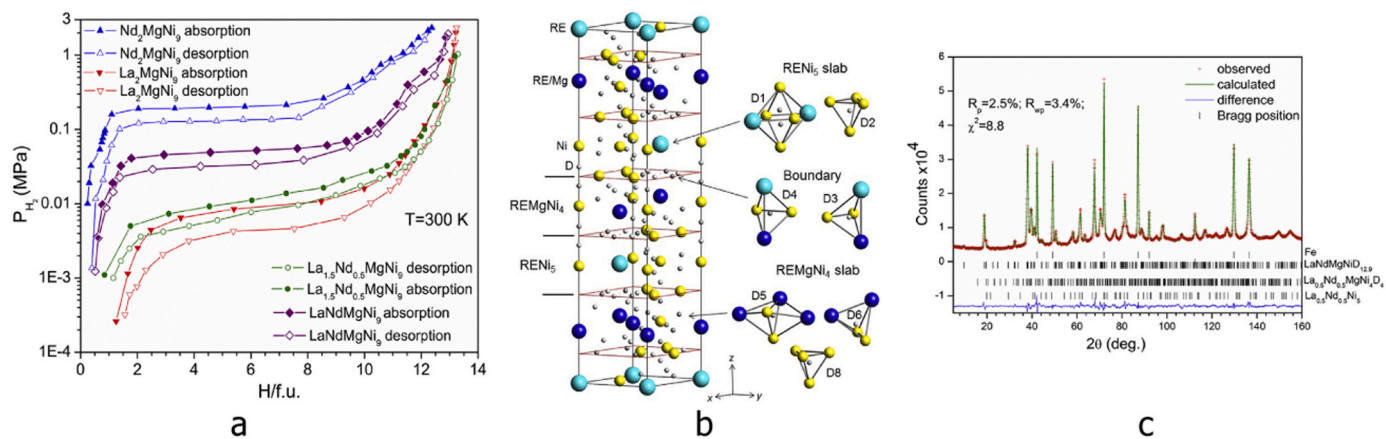
distribution of D atoms within the  $\text{REMgNi}_4$  and  $\text{RENi}_5$  layers and filling by D atoms of 7 types of interstitial sites (Fig. 38b,c). During the electrochemical cycling (Fig. 39a), formation of an extended  $\alpha$ -solid solution spanning from  $\text{LaNdMgNi}_9\text{D}_{0.8}$  to  $\text{LaNdMgNi}_9\text{D}_{1.6}$  was observed, whereas the homogeneity range for the  $\beta$ -deuteride occurred between  $\text{LaNdMgNi}_9\text{D}_{8.2}$  and  $\text{LaNdMgNi}_9\text{D}_{9.8}$ . The reversible cycling was accompanied by the discrete transformation between  $\alpha$  and  $\beta$  phases (Fig. 39b). The same reaction pathway was observed during the charge and discharge yielding a completely reversible cycling of the metal hydride electrode. A fully discharged state (i.e., 100% of the hydrogen-free intermetallic alloy) was not reached and the amount of electrochemically released hydrogen was significantly smaller than for the metal-gas PCT measurement.

Comparison of the properties of the hydride materials containing a variable amount of Nd and increasing in a row  $\text{La}_2\text{MgNi}_9$ - $\text{La}_{1.5}\text{Nd}_{0.5}\text{MgNi}_9$ - $\text{LaNdMgNi}_9$ - $\text{Nd}_2\text{MgNi}_9$ , showed that presence of Nd is important to improve the anode electrodes of the  $\text{RE}_2\text{MgNi}_9$  type

and shows their clear advantages (broader and flatter plateaus, faster H exchange rates, better cycle stability and high-rate discharge performance) over the Nd-free compositions. Optimization of the Nd content in the  $\text{La}_{2-y}\text{Nd}_y\text{MgNi}_9$  alloys will help in achieving the most efficient performance of the electrodes working in a high power regime [144].

### 7.1.3. TiNi-D<sub>2</sub> based systems

As rare earths are critical raw materials [145], their replacement in the negative electrodes of the commercial Ni-MH batteries would be a great success. Among all rare-earth free hydride forming alloys, AB-type compounds are of particular interest as they can be formed by light and relatively cheap metals. TiFe, one of the most popular hydrogen stores, forms  $\text{TiFeH}_2$  with a hydrogen capacity of 1.87 wt% near standard conditions of pressure and temperature [104]. It would be a good candidate as negative electrode. But, unfortunately, it exhibits no electrochemical activity due to passivation issues. In



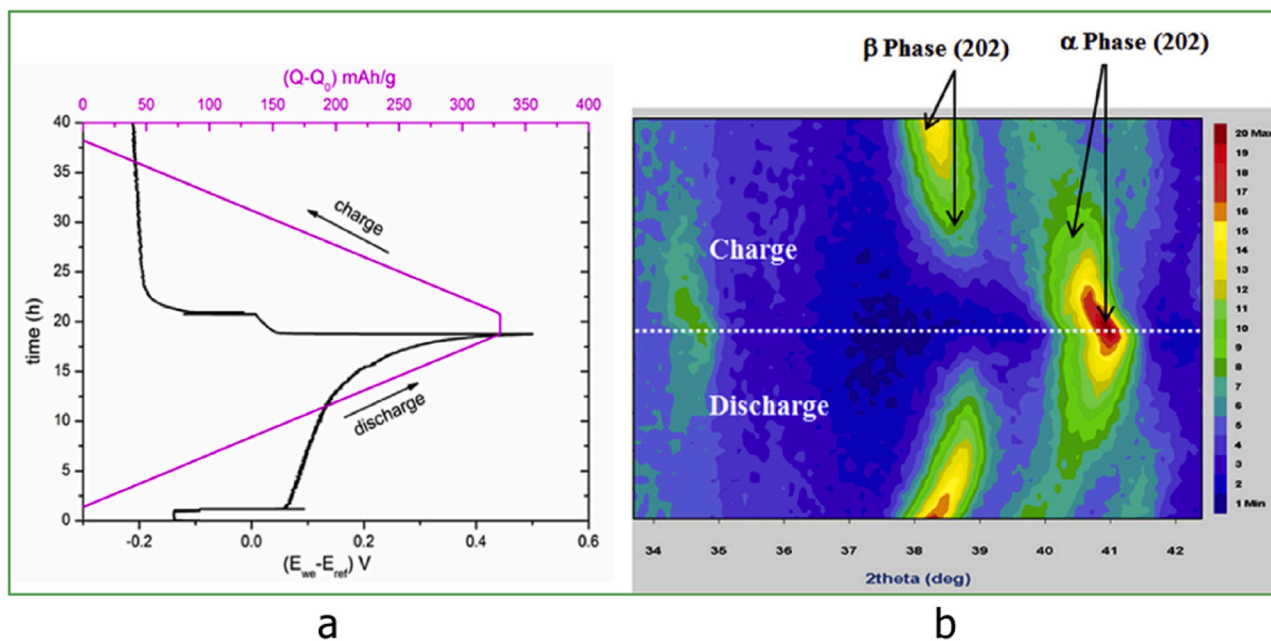
**Fig. 38.** (a) PCT isotherms measured for the  $\text{LaNdMgNi}_9\text{-H}_2$  system at 300 K [144] as compared to the reference data for  $\text{La}_2\text{MgNi}_9\text{-H}_2$  [134],  $\text{Nd}_2\text{MgNi}_9\text{-H}_2$  [135] and  $\text{La}_{1.5}\text{Nd}_{0.5}\text{MgNi}_9\text{-H}_2$  [138] systems; (b) Crystal structure of  $\text{LaNdMgNi}_9\text{D}_{12.9}$  showing the stacking of the (La,Nd)Ni<sub>5</sub> and (La,Nd,Mg)Ni<sub>4</sub> slabs. La and Nd are labelled as RE (rare earth metals). From 7 sites filled by D, two are located inside the (La/Nd)Ni<sub>5</sub> layer (D1, D2), three inside the (La/Nd)MgNi<sub>4</sub> layer (D5, D6, D8), and two at the boundary between the two slabs (D3, D4). The vacant D7 site is not shown. (c) Rietveld refinements of the NPD pattern of  $\text{LaNdMgNi}_9$  alloy saturated with deuterium. Wavelength used for NPD -  $\lambda = 1.494\text{ \AA}$ . The pattern show presence of three identified phase constituents, including  $\text{LaNdMgNi}_9\text{D}_{12.9} - 91.3(2)\text{ wt\%}$ ,  $\text{La}_{0.5}\text{Nd}_{0.5}\text{MgNi}_4\text{D}_4 - 6.0(2)\text{ wt\%}$ ,  $\text{La}_{0.5}\text{Nd}_{0.5}\text{Ni}_5 - 2.7(2)\text{ wt\%}$  and a stainless steel sample holder (Fe). (Reprinted with permission from [144]. Copyright (2017) Elsevier).

contrast, TiNi has an excellent electrocatalytic activity and limited corrosion in KOH-concentrated alkali media [146]. Indeed, TiNi-based electrodes were first proposed as active materials for Ni-MH batteries. The major handicap of TiNi is however its low reversible capacity, 150 mAh/g, which corresponds to an exchange of 0.6 H atoms per TiNi formula unit [147]. This amount is much lower than for the fully hydrogenated hydride  $\text{TiNiH}_{1.4}$ , which suggests that some hydrogen is irreversibly trapped in the alloy during electrochemical cycling.

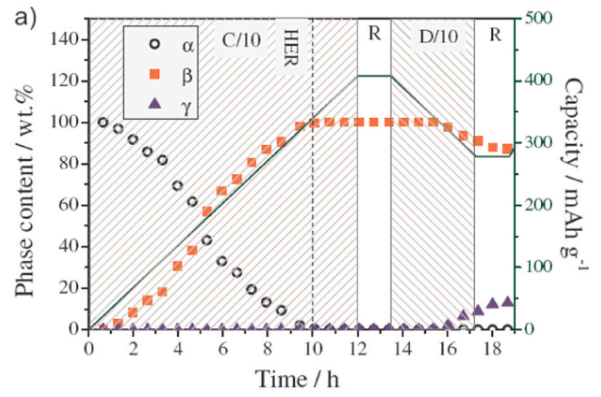
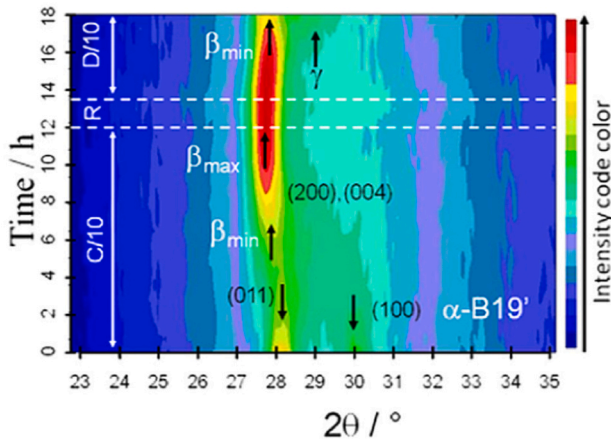
To improve the properties of TiNi electrodes, different chemical substitutions have been studied: Zr for Ti as well as Pd, Co and Cu for Ni [51, 147–149]. The best electrochemical properties have been obtained so far with Cu substitution for which the reversible capacity is doubled for  $\text{TiNi}_{0.8}\text{Cu}_{0.2}$  alloy compared to binary TiNi [51]. As now detailed, *in situ* neutron diffraction studies coupled to

galvanostatic electrochemical cycling for TiNi and  $\text{TiNi}_{0.8}\text{Cu}_{0.2}$  shed light on the origin of this remarkable result [51].

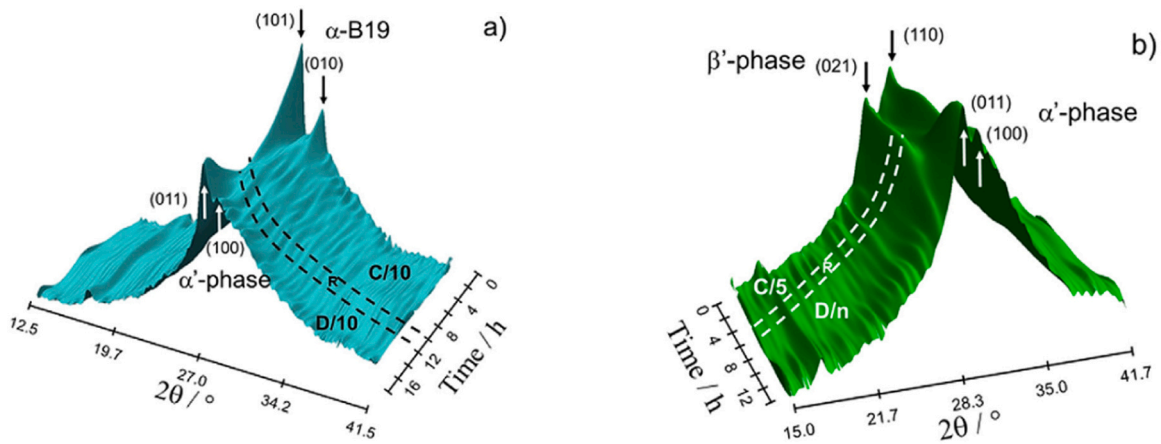
For binary TiNi, the electrode was charged at C/10 rate for 12 h, relaxed for 1.5 h in open circuit, and discharged at D/10 rate up to the cut-off potential of 0.5 V versus  $\text{Cd}/\text{Cd}(\text{OH})_2$ . The 2D diffraction patterns displayed in Fig. 40a show that the diffraction peaks referring to the pristine  $\alpha\text{-B19'}$  martensitic phase of TiNi (S.G.  $P2_1/m$ ) gradually disappear on charge while those belonging to  $\beta$ -hydride (S.G.  $I4/mmm$ ) emerge and steadily shift to lower angles. Thus, diffraction patterns reflect first the  $\alpha$  to  $\beta$  phase transformation followed by a solid solution of hydrogen from  $\beta_{\min}$  ( $\text{TiNiH}_{1.0\pm 0.1}$ ) to  $\beta_{\max}$  ( $\text{TiNiH}_{1.3\pm 0.1}$ ) in the hydride phase. On discharge, diffraction peaks from  $\beta$ -hydride shift back to higher angles evidencing hydrogen release within the hydride phase. In addition, approaching to the cut-off potential, a novel diffraction peak starts appearing attributed to



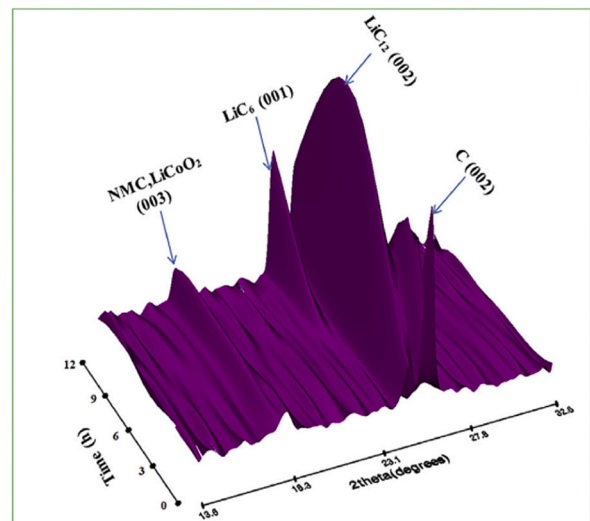
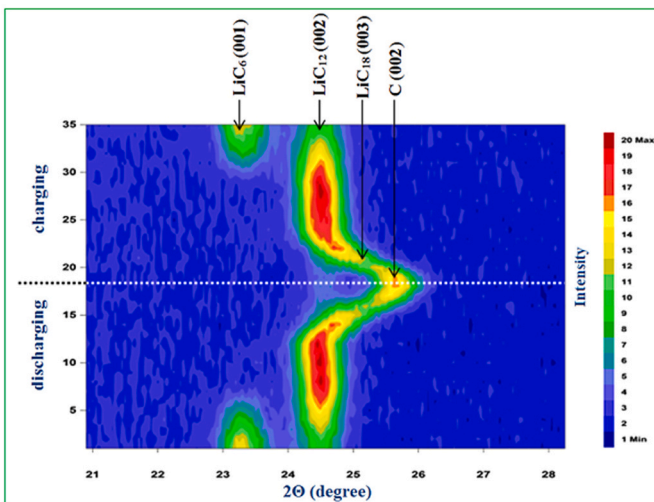
**Fig. 39.** (a) Discharge/charge potential profile and electrochemical capacity for the  $\text{LaNdMgNi}_9$  electrode. (b) 2D contour plot of NPD patterns within the angular domain  $34^\circ < 2\theta < 42^\circ$  showing the evolution of (202) diffraction lines of  $\alpha$ -metal and  $\beta$ -deuteride phases during the galvanostatic cycling. (Reprinted with permission from [144]. Copyright (2017) Elsevier).



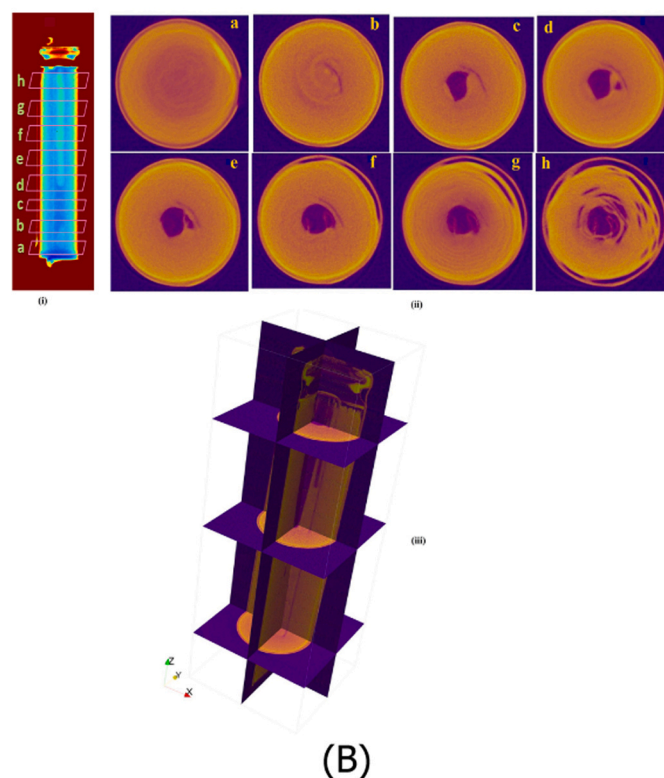
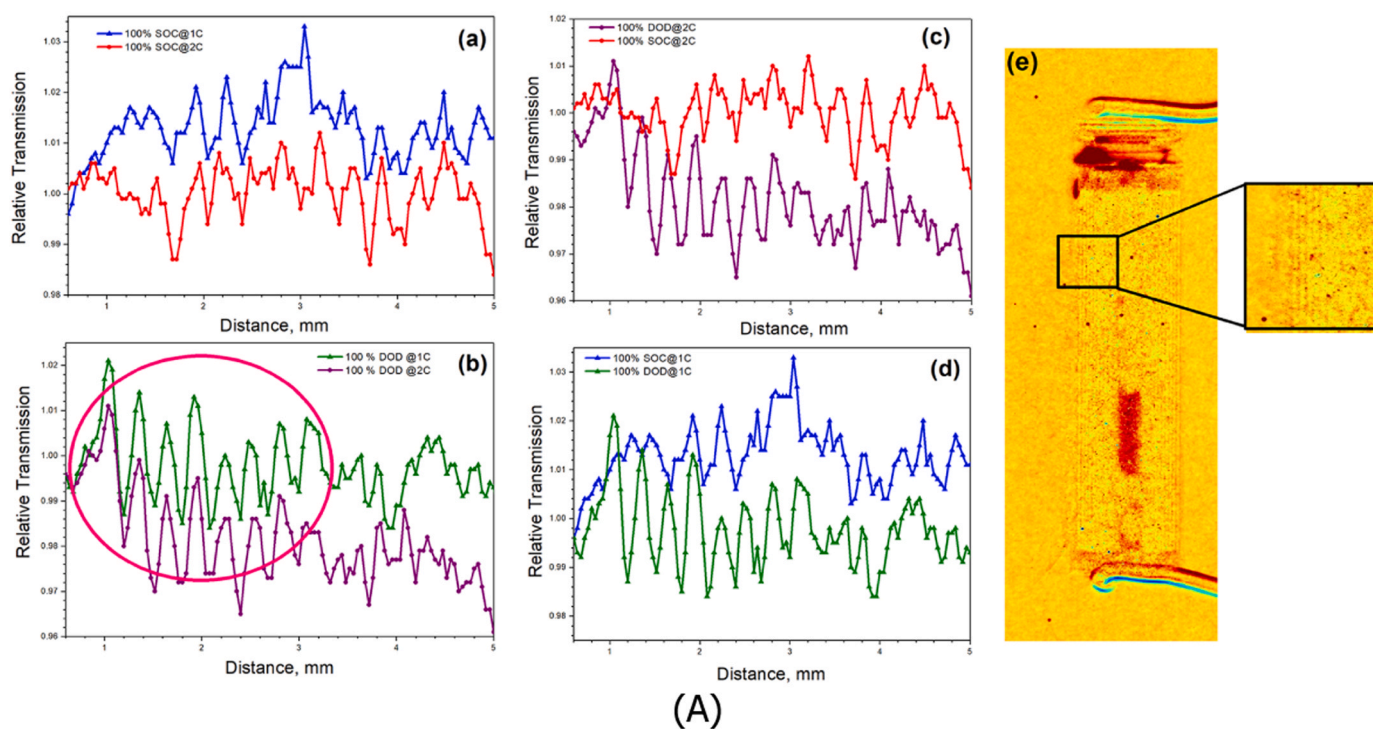
**Fig. 40.** (a) First charge/discharge cycle of TiNi electrode 3D. 2D view of the neutron diffraction pattern evolution as function of time; (b) phase amount as a function of time. (Reprinted with permission from [51]. Copyright (2014) Elsevier).



**Fig. 41.** 3D *in situ* neutron diffraction patterns of TiNi<sub>0.8</sub>Cu<sub>0.2</sub> electrode. (a) First charge/discharge cycle; (b) Fourth galvanostatic cycle. Phase identification and indexation of main peaks are given. (Reprinted with permission from [51]. Copyright (2014) Elsevier).



**Fig. 42.** (a) 2D-plot of a selected area of the *in situ* ND pattern of anode during a discharge-charge cycle. The color scale (right axis) represents the diffraction peaks intensities (red: the highest intensity; blue - the lowest intensity). (b) 3D plot of the *in situ* NPD pattern during the charging of the battery. A parallel development of the (003) Bragg peak of the LiCoO<sub>2</sub> and Li(Ni,Mn,Co)O<sub>2</sub>, together with graphite (002), LiC<sub>12</sub> (002) and LiC<sub>6</sub> (001) peaks is shown. (Reprinted with permission from [150]. Copyright (2016) Elsevier).



**Fig. 43.** (a) Comparison of fully charged states at 1 C and 2 C; (b) Comparison of fully discharged states at 1 C and 2 C; (c) Comparison of fully charged and fully discharged states at 2 C; (d) Comparison of fully charged and fully discharged states at 1 C; (e) Radiographic image of the discharged battery normalized to fully charged state and details of the area vertically averaged and depicted in (a–d). (b) Tomography slices from 3D reconstructions (ii) at different heights of the battery marked in the radiograph (i) and (iii) 3D visualization of the commercial ICR 10440 Li-ion cell (fully charged), reconstructed from the tomographic slices acquired with a total of 300 projections with 90 s exposure time. Reprinted from an open access article [151]. Distributed under the Creative Commons Attribution license (<http://creativecommons.org/licenses/by/4.0/>).

the formation of a minor amount of a novel  $\gamma$ -phase, the structure of which could not be determined. The evolution of phase contents during the whole cycle is shown in Fig. 40b. The reversible capacity for this cycle attains 130 mAh/g, being most of it provided by the

hydrogen solid solution in  $\beta$ -hydride and, to a much lower extent, by incomplete transformation ( $\sim 10\%$ ) between  $\beta$  and  $\gamma$ -phases.

For  $\text{TiNi}_{0.8}\text{Cu}_{0.2}$  electrode, structural transformations were observed to change on cycling (Fig. 41a,b). On the first charge at C/10

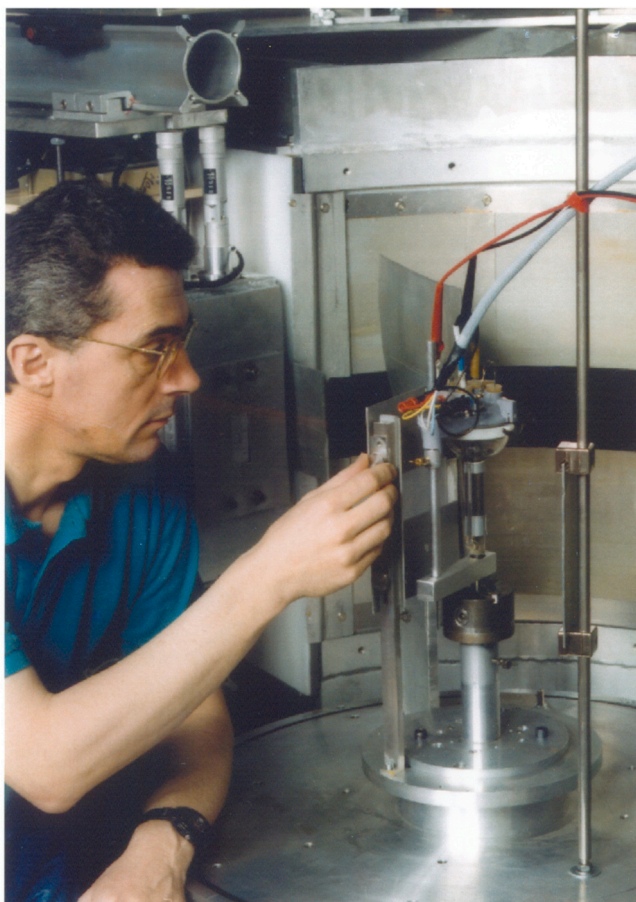


Fig. 44. Michel Latroche during the experiments with metal hydride anodes performed at D20, ILL in 1998.

regime,  $\alpha$ -B19 martensitic phase (S.G.  $Pmmb$ ) undergoes a phase transformation to a poorly crystallized phase. Then, on discharge at 10 D regime, a novel intermetallic phase  $\alpha'$  (S.G.  $P2_1/m$ ) with structure similar to B19' martensitic phase is formed. The reversible capacity for this first cycle attained 235 mAh/g. Later, the capacity of the electrode increased up to the fourth cycle due to an activation process. A 3D-diffraction plot of neutron diffraction pattern for the four cycle is shown in Fig. 41b. At this cycle, a reversible two-phase transformation occurs between the intermetallic phase  $\alpha'$  (S.G.  $P2_1/m$ ) and a hydride phase  $\beta'$  which structure could not be fully resolved. The reversible capacity for this two-phase transformation attained as much as 330 mAh/g.

To summarize, *in situ* neutron diffraction studies demonstrated that the low reversible electrochemical capacity in TiNi mainly originates from hydrogen stored in  $\beta$ -hydride solid-solution. In contrast, the high electrochemical capacity of  $\text{TiNi}_{0.8}\text{Cu}_{0.2}$  originates from a two-phase transformation between intermetallic and hydride phases. Complementary PCT thermodynamic studies on pseudobinary  $\text{TiNi}_{1-x}\text{Cu}_x$  ( $0 \leq x \leq 0.5$ ) alloys showed that the change of the reaction mechanism with Cu substitution results from Cu-induced destabilization of  $\beta$ -hydride [51].

### 7.2. Lithium ion batteries studied *in situ* by neutron diffraction, neutron radiography and tomography [150,151]

The structural evolution of the electrode materials in an ICR 10440 commercial cylindrical lithium-ion battery, which has a discharge capacity of 360 mAh and a nominal voltage of 3.7 V has been studied using *in situ* neutron diffraction. A three-phase mixture of Li (Ni,Mn,Co) $\text{O}_2$ , LiCo $\text{O}_2$  and LiMn $\text{O}_4$  was identified as the active

material of the cathode, with graphite acting as the anode material. The study revealed that the graphite anode underwent structural changes to form a series of insertion-type lithiated derivatives, with up to 12.7% volume expansion for the Li-saturated compound LiC $_6$  (see Fig. 42a). The charge-discharge behaviour was more complex for the cathode. Here, the charge process was associated with partial lithium depletion from the initially Li-saturated compounds, leading to volume shrinkage for Li(Ni,Mn,Co) $\text{O}_2$ , in contrast to (Ni,Mn)-free LiCo $\text{O}_2$  (see Fig. 42b). Electrochemical discharge experiments performed under a fast regime (2 C) at 5, 25 and 45 °C revealed that the discharge capacity followed the trend of an increased diffusion rate of Li $^+$  ions in the electrolyte and Li atoms in both electrodes, being highest for 45 °C. At the lowest tested temperature (5 °C), a rapid drop in the discharge capacity took place using the same kinetic regime.

An *operando* neutron imaging study of a commercial ICR 10440 Li ion battery during charge and discharge showed (Fig. 43a,b) that the cylindrical battery has a spiral configuration which is composed of a multiphase layered oxide cathode and graphite anode. In spite of a two-dimensional nature of the projection data of this time-resolved study, structural and functional details of the neutron radiography study were successfully uncovered and visualized. The spatially resolved measurements with a resolution of 40  $\mu\text{m}$  enabled to observe Li redistribution between the electrodes as well as a circulation of the electrolyte between the central column and the electrode layers at different states of charge (SoC) and at different current rates. Furthermore, *ex-situ* tomographic studies of the battery revealed the fine details of the structural inhomogeneity within the cell.

## 8. Concluding remarks

This review presents numerous examples of the studies of hydrogen and energy storage materials, where neutron diffraction, and particularly when combined with Synchrotron X-ray diffraction, becomes a valuable tool in *in situ* probing of hydrogen storage properties. Even when using small, 100–500 mg, samples, these studies allow establishing the mechanism of phase-structural transformations and their kinetics based on rapid data collection at various charge-discharge conditions. These conditions include a broad range of applied deuterium pressures, from vacuum to high pressures reaching 1000 bar H $_2$ /D $_2$ , and temperatures, from cryo-cooling (2 K) to as high as 1273 K (1000 °C), while various state-of-charge and discharge setpoints have been achieved by measuring the PCT diagrams of H $_2$ /D $_2$  absorption-desorption during the collection of the diffraction data or by applying variable current densities during the testing of the battery electrodes.

Collection of high quality *in situ* diffraction data benefits from a high flux beam and its combination with high resolution measurements, while the use of dedicated sample cells allows the application of appropriate testing conditions to reproduce the conditions of use of energy storage materials in the lab and in real applications. This is because the profile intensities of the diffraction pattern are utilized in the most efficient way as the data fitting covers a large  $d$ -spacing range and ensures improvements in the peak-to-background and signal-to-noise ratios thus mitigating the challenges caused by the small size of some studied samples, their low crystallinity, incoherent scattering of hydrogen atoms and diffraction signals from complex sample environments such as electrolytes. Most importantly, complex structures with large unit cells and low symmetry can be successfully determined, which is of particular importance in the studies of the phase-structural composition of single phase and multiphase hydride materials, even when in small amounts.

*In situ* experiments, performed under hydrogen/deuterium pressure at a range of temperatures, make demands on the type of sample cell used. It is necessary to simultaneously satisfy the

requirements of a material whose yield strength can support the experimental pressures used to the temperatures reached, does not deteriorate or embrittle with hydrogen and does not otherwise hinder the acquisition of quality diffraction data. A number of cells for specific solid-gas applications has been discussed. Moreover, for structural studies of metal hydride electrodes, amorphous silica cells which do not produce diffraction peaks have been used. Ancillary devices such as Sieverts' apparatuses, pressure gauges and galvanostats are used to control and monitor hydrogen/deuterium uptake and release during *in situ* neutron diffraction measurements. They are used in solid-gas and electrochemical studies.

The ability of neutrons to penetrate deeply into the matter makes neutron diffraction especially valuable for the study of bulk materials. The use of non-isotopically enriched samples makes many experiments far more accessible, and this has been established during *in situ* studies of protium-containing materials.

The use of multiple experimental diffraction data sets has become increasingly important. SR XRD data provide accurate information about the positions of heavy atoms (metal components) in the structure, which can then be fixed in the subsequent refinements of hydrogen positions using NPD data. Simultaneous refinement of the SR XRD and NPD data is the most preferable option which results in excellent resolution of hydrogen positions.

Major neutron instruments are open to scientists around the world. However, access to these instruments is competitive, based on the evaluation of the competing proposals, and is not normally straightforward even for accepted proposals as there is a time gap between the submission of the proposals and the scheduled experiments. The High Flux Isotope Reactor at ORNL, high-flux reactor at Institut Laue-Langevin (Grenoble, France), Neutron-Beam Split Core Reactor at NIST, FRM II (Munich, Germany), and OPAL reactor at the Australian Centre for Neutron Scattering (Lucas Heights, Australia) are on the list of the available nuclear research reactor institutions for the international community. This list is complemented by the neutron spallation sources which include the Spallation Neutron Source at ORNL, ISIS Neutron and Muon Source at Rutherford Appleton Laboratory (Chilton, UK), J-PARC (Japan), and the Swiss Spallation Neutron Source (SINQ), a continuous source at Paul Scherrer Institute (Villigen, Switzerland). Finally, the European Spallation Source is under construction in Lund, Sweden, and will be operating soon.

This paper is dedicated to the memory of Michel Latroche. In the context of this manuscript on *in situ* scattering analyses of structural determination and phase transformation in energy storage materials, the contribution of Michel Latroche is remarkable. Latroche's work has been mainly focused on the fundamental understanding of the relationships between structures, thermodynamic properties and physical properties of materials, without neglecting their applications in the field of chemical and electrochemical energy storage. In this perspective, his major achievements concern the fundamentals of metal hydrides and their application either as hydrogen gas storage for fuel cells or as negative electrodes for Ni-MH and Li-ion batteries [123, 143, 152–157].

After joining the Percherons-Guégan laboratory of CNRS France in 1989, Michel Latroche carried out pioneering work on *in situ* neutron diffraction studies of LaNi<sub>5</sub> based electrodes [45, 47–50, 63, 158]. In the 1990s, in collaboration with his colleague Yves Chabre, he developed the *in situ* neutron diffraction cell shown in Fig. 5 [45], which has been used during the experiments at D20, ILL in 1998 (see Fig. 44), and is still used today. In the present century, electrochemical, structural and mechanistic studies using *in situ* neutron diffraction have expanded from AB<sub>5</sub>-type electrodes to AB<sub>3</sub> [52,144,154], AB [51] and Mg-based electrodes [159]. In parallel with the structural studies concerning electrochemical storage, Latroche's work on *in situ* neutron diffraction has also provided a broad overview of materials for solid state hydrogen storage. In addition to the

classical AB<sub>5</sub> compounds, over-stoichiometric ones [35,109,160], AB<sub>2</sub> [60,81,161] and AB [55] intermetallics were considered. Last but not least, many other classes of materials have attracted his attention beyond intermetallics, such as Pd@C hybrids [162], Mg-based composites [36] and imide/amide systems [37]. Structural analysis of both battery and hydrogen gas storage materials has provided an in-depth understanding of the mechanisms underlying reversible hydrogen uptake, as well as their kinetic behaviour and cycle-life properties.

The recent progress in the development of advanced materials for hydrogen-based energy storage have been presented in recently published review papers [90,114,124,163]. We anticipate that *in situ* diffraction studies of hydrogen and energy storage materials, providing insights into understanding of their performance and allowing optimisation of the materials, will be growing further based on the use of advanced large-scale research facilities accessible to a broad international research community. This will facilitate progress in the socially important field of energy storage and conversion.

### CRedit authorship contribution statement

**V.A. Yartys:** Conceptualization, Writing – review & editing, Funding acquisition. **F. Cuevas:** Conceptualization, Writing – review & editing, Funding acquisition. **C.J. Webb:** Conceptualization, Writing – review & editing, Funding acquisition.

### Data Availability

Data will be made available on request.

### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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